

2-D problems

$$-\frac{\partial}{\partial x} \left( a_{11} \frac{\partial u}{\partial x} + a_{12} \frac{\partial u}{\partial y} \right) - \frac{\partial}{\partial y} \left( a_{21} \frac{\partial u}{\partial x} + a_{22} \frac{\partial u}{\partial y} \right) + a_0 u - f = 0$$

For Poisson's eqn  $a_{12} = a_{21} = a_0 = 0$

For Laplace eqn  $a_{12} = a_{21} = a_2 = f = 0$

$$\int_{\Omega} w \left[ \underbrace{-\frac{\partial}{\partial x} \left( a_{11} \frac{\partial u}{\partial x} + a_{12} \frac{\partial u}{\partial y} \right)}_{F_1} - \underbrace{\frac{\partial}{\partial y} \left( a_{21} \frac{\partial u}{\partial x} + a_{22} \frac{\partial u}{\partial y} \right)}_{F_2} + a_0 u - f \right] dx dy = 0$$

$$\int_{\Omega} w \left[ -\frac{\partial}{\partial x} F_1 - \frac{\partial}{\partial y} F_2 \right] dx dy + \int_{\Omega} w a_0 u dx dy - \int_{\Omega} w f dx dy = 0$$

$$w \frac{dk}{dx} = \frac{d(wk)}{dx} - k \frac{dw}{dx}$$

$$-w \frac{dk}{dx} = k \frac{dw}{dx} - \frac{d(wk)}{dx}$$

$$\int_{\Omega} \left( \frac{\partial w}{\partial x} F_1 + \frac{\partial w}{\partial y} F_2 \right) dx dy - \int_{\Omega} \frac{\partial (w F_1)}{\partial x} dx dy - \int_{\Omega} \frac{\partial (w F_2)}{\partial y} dx dy + A - B = 0$$

$$\int_{\Omega} \frac{\partial (w F_1)}{\partial x} dx dy = \int_S w F_1 n_x ds$$

$$\int_{\Omega} \frac{\partial (w F_2)}{\partial y} dx dy = \int_S w F_2 n_y ds$$

$$\int_{\Omega} \left[ \frac{\partial w}{\partial x} \left( a_{11} \frac{\partial u}{\partial x} + a_{12} \frac{\partial u}{\partial y} \right) + \frac{\partial w}{\partial y} \left( a_{21} \frac{\partial u}{\partial x} + a_{22} \frac{\partial u}{\partial y} \right) \right] dx dy + \int_{\Omega} w a_0 u dx dy = \int_{\Omega} w f dx dy + \int_S w \left( a_{11} \frac{\partial u}{\partial x} + a_{12} \frac{\partial u}{\partial y} \right) n_x ds + \int_S w \left( a_{21} \frac{\partial u}{\partial x} + a_{22} \frac{\partial u}{\partial y} \right) n_y ds$$

Simplified form,  $\int_{\Omega} \left( \frac{\partial w}{\partial x} a_{11} \frac{\partial u}{\partial x} + \frac{\partial w}{\partial y} a_{21} \frac{\partial u}{\partial x} \right) dx dy = \int_{\Omega} w f dx dy + \int_S w q_n ds$

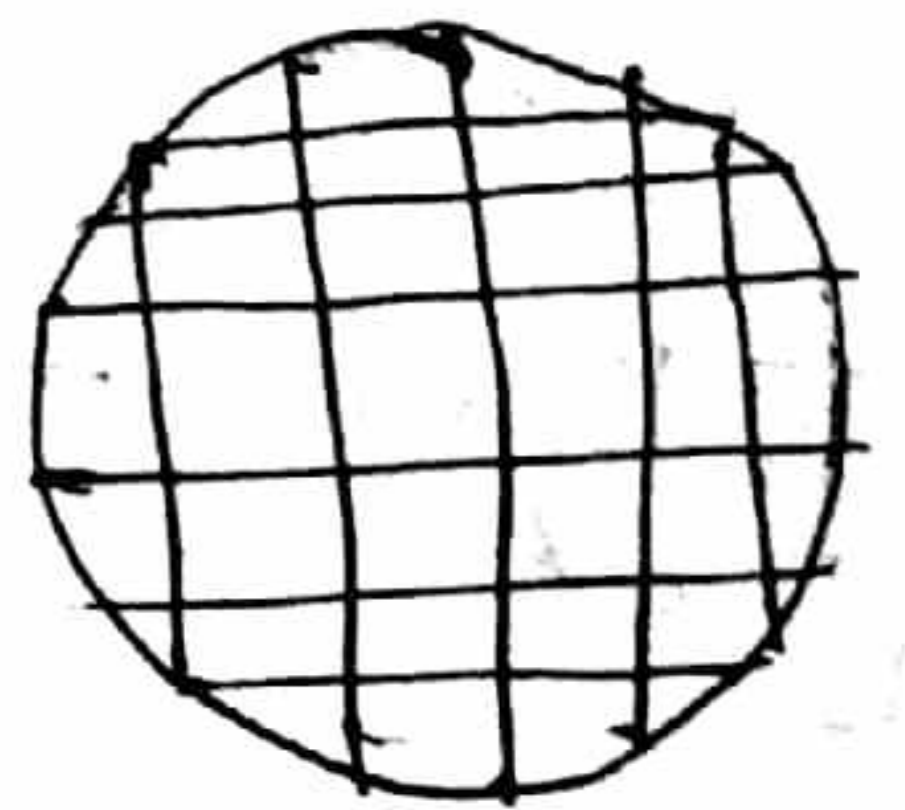
$$q_m = a \left( \frac{\partial \psi}{\partial x} n_x + \frac{\partial \psi}{\partial y} n_y \right)$$

R-R-F  $\Rightarrow \int_{\Omega} \left[ \frac{\partial \psi_i}{\partial x} a \sum_{j=1}^n \frac{\partial}{\partial x} (u_j \psi_j) + \frac{\partial \psi_i}{\partial y} a \sum_{j=1}^n \frac{\partial}{\partial y} (u_j \psi_j) \right] dx dy$

$$= \int_{\Omega} \psi_i A dx dy + \int_S \psi_i 2n ds$$

Matrix of  $n \times n$

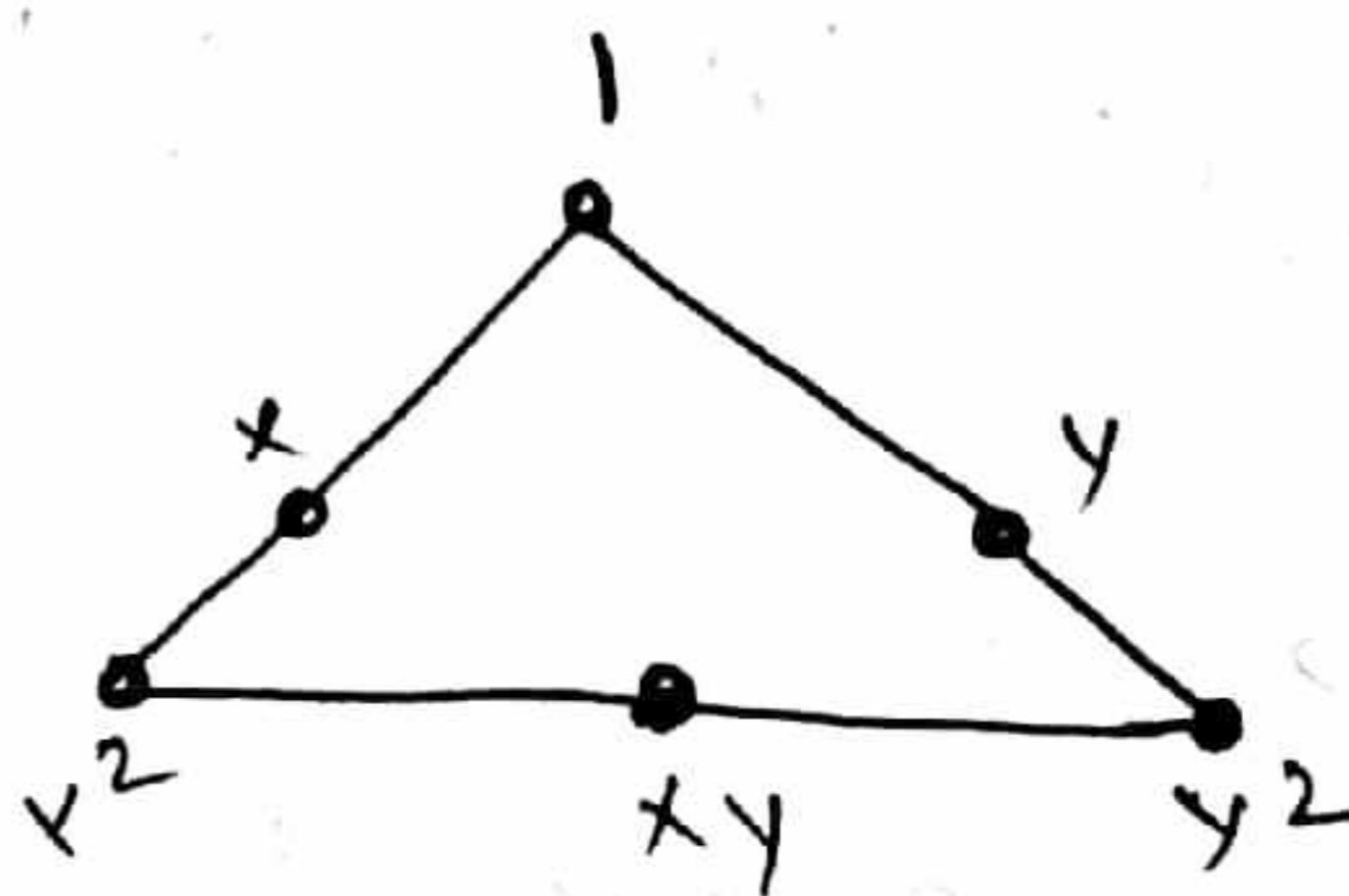
The elements are of two types ① Rectangular ② Triangular



- $x$
- $x^2$
- $x^3$
- $x^4$
- $xy$
- $x^2y$
- $x^2y^2$
- $xy^2$
- $xy^3$
- $y^2$
- $y^3$
- $y^4$

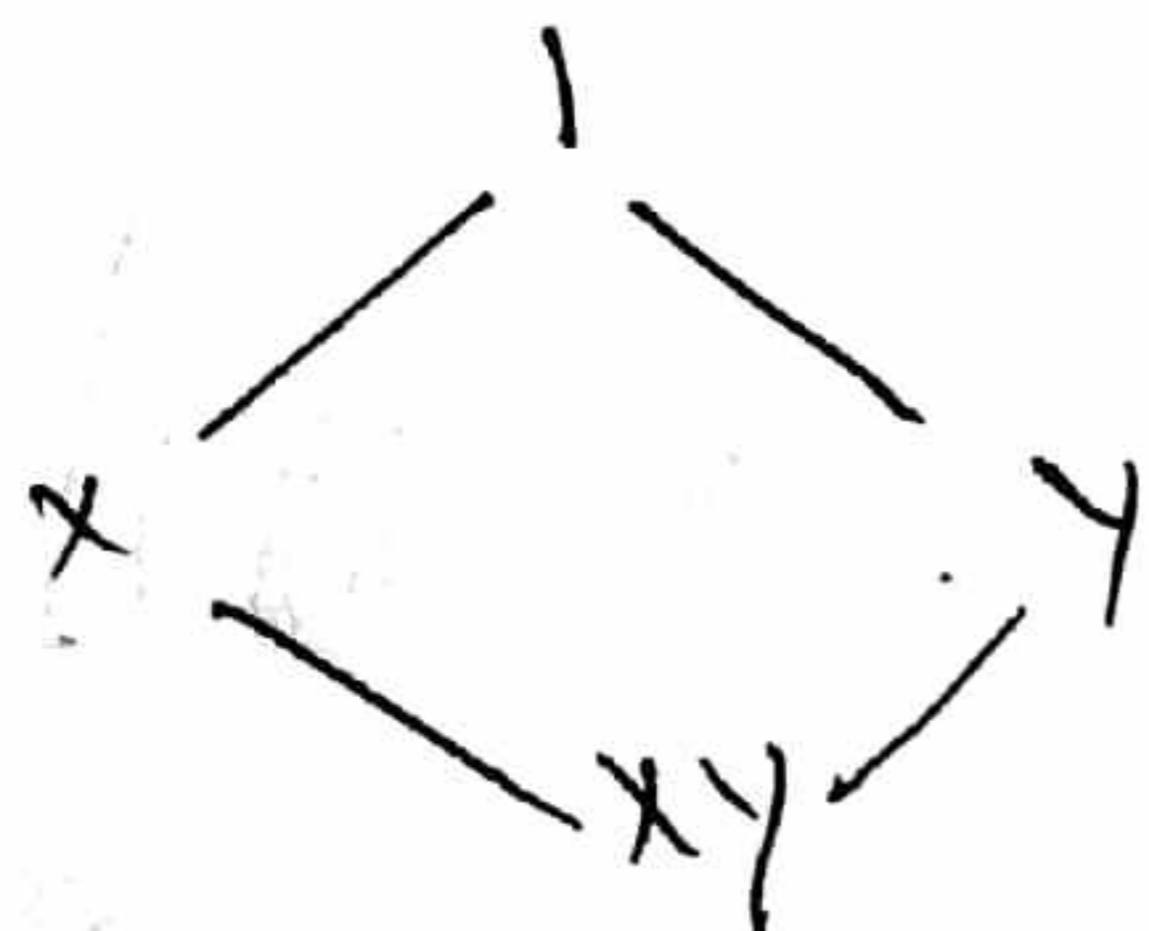
For quadratic triangular element.

$$1 + x + x^2 + xy + y^2 + y$$

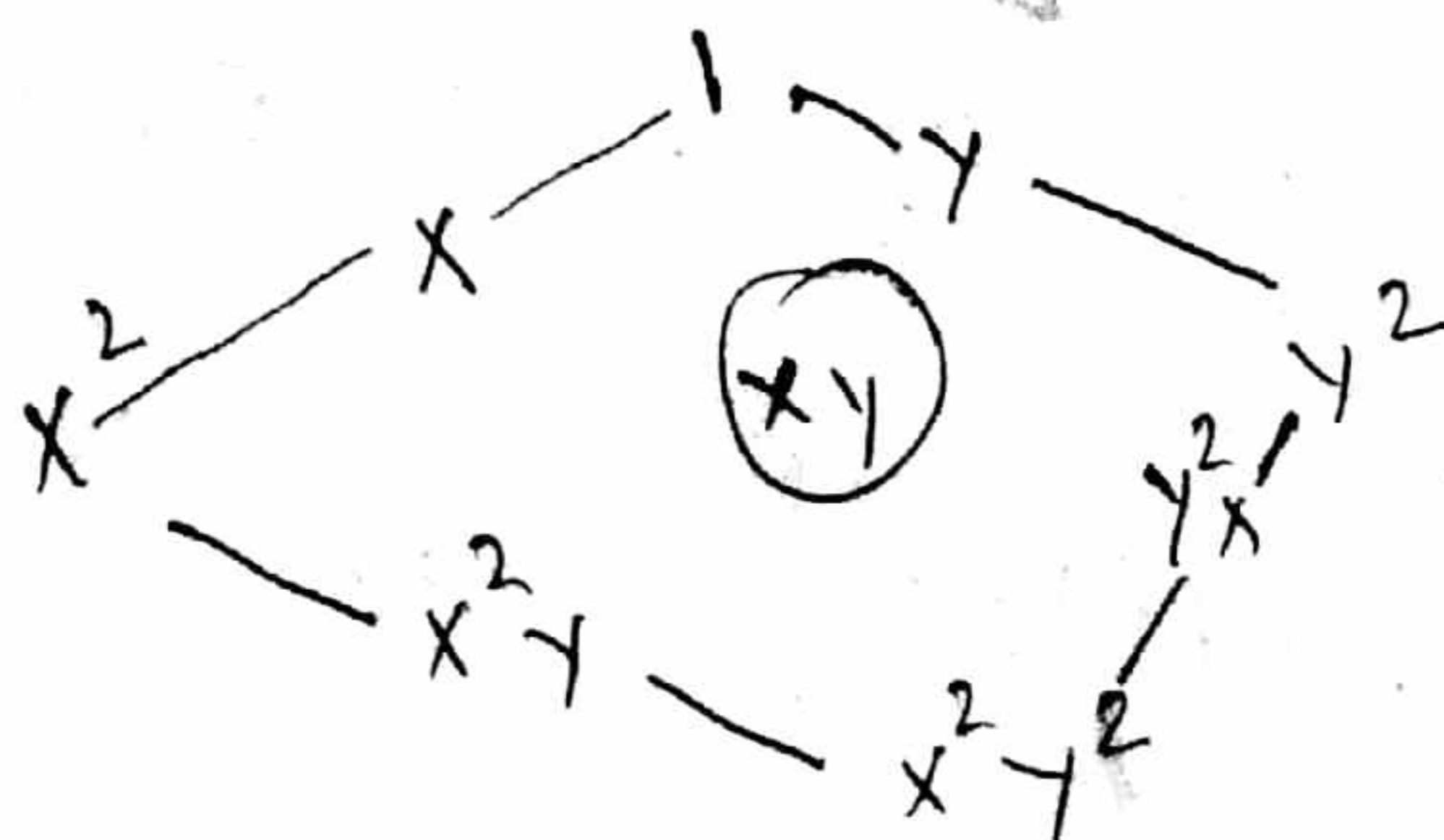


Hermitian interpolation function  $\rightarrow$  when the differential eq needs the continuity of derivative of field variables.

For Linear rectangular element.



For quadratic rectangular element



Total no. of nodal points = 9

$xy$  is at the center of the rectangle.

Boundary terms  $\rightarrow$  coefficient of weight functions / for  
 EBC  $\rightarrow$  ~~coefficient~~ The form of  $w$  in weak form / 1-D

In global coordinates -  

$$y_i = \frac{1}{2A} (\alpha_i + \beta_i x + \gamma_i y)$$

{ linear triangular element }

$A =$  area of the triangle.

$$\alpha_i = x_j y_k - x_k y_j$$

$$\beta_i = y_j - y_k$$

~~$\gamma_i = x_j - x_k$~~   $\gamma_i = -(x_j - x_k)$

$$\alpha_1 + \alpha_2 + \alpha_3 = 2A$$

$$1 \rightarrow 2 \rightarrow 3$$

$$2 \rightarrow 3 \rightarrow 1$$
  

$$i \quad j \quad k$$

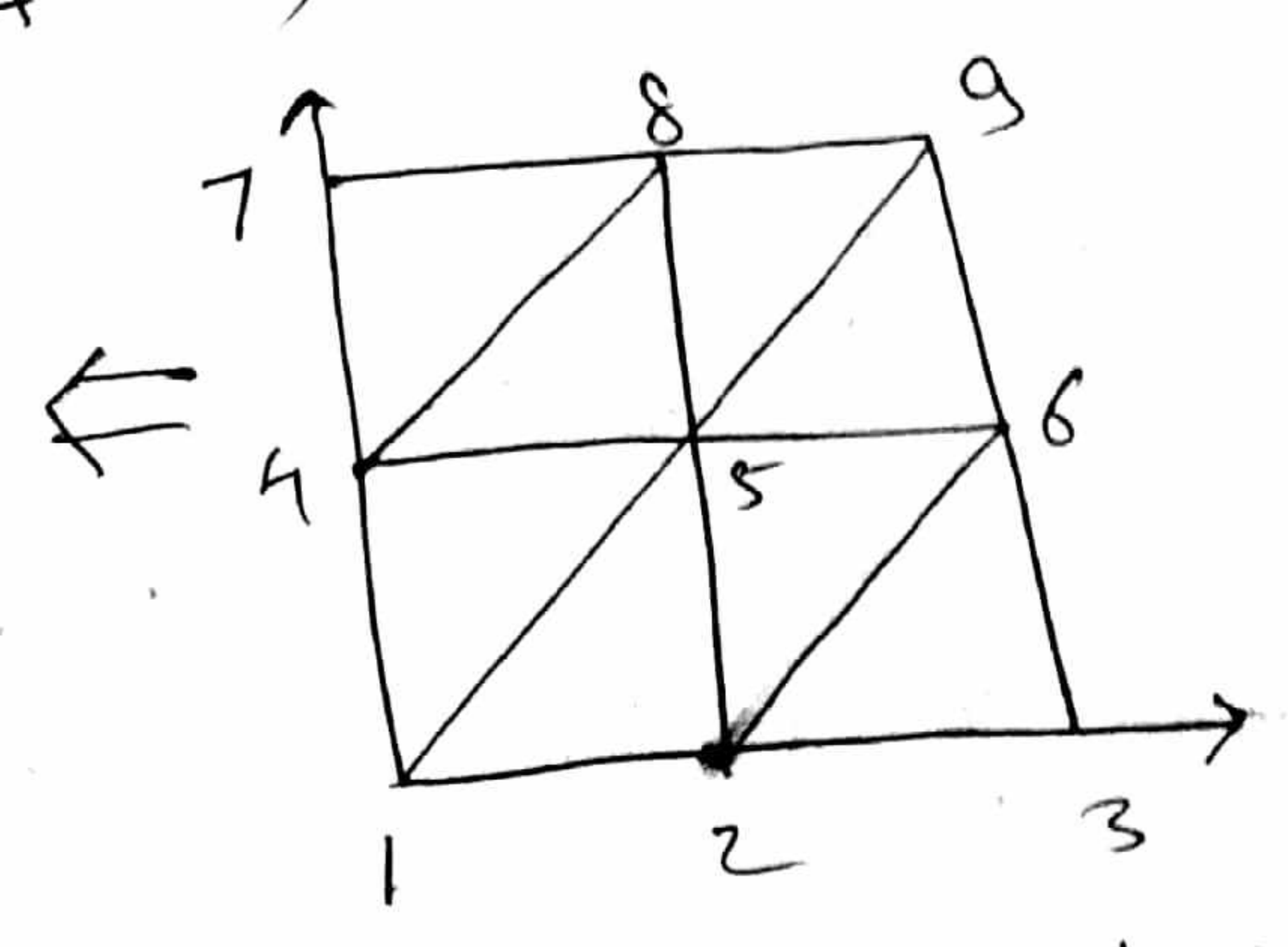
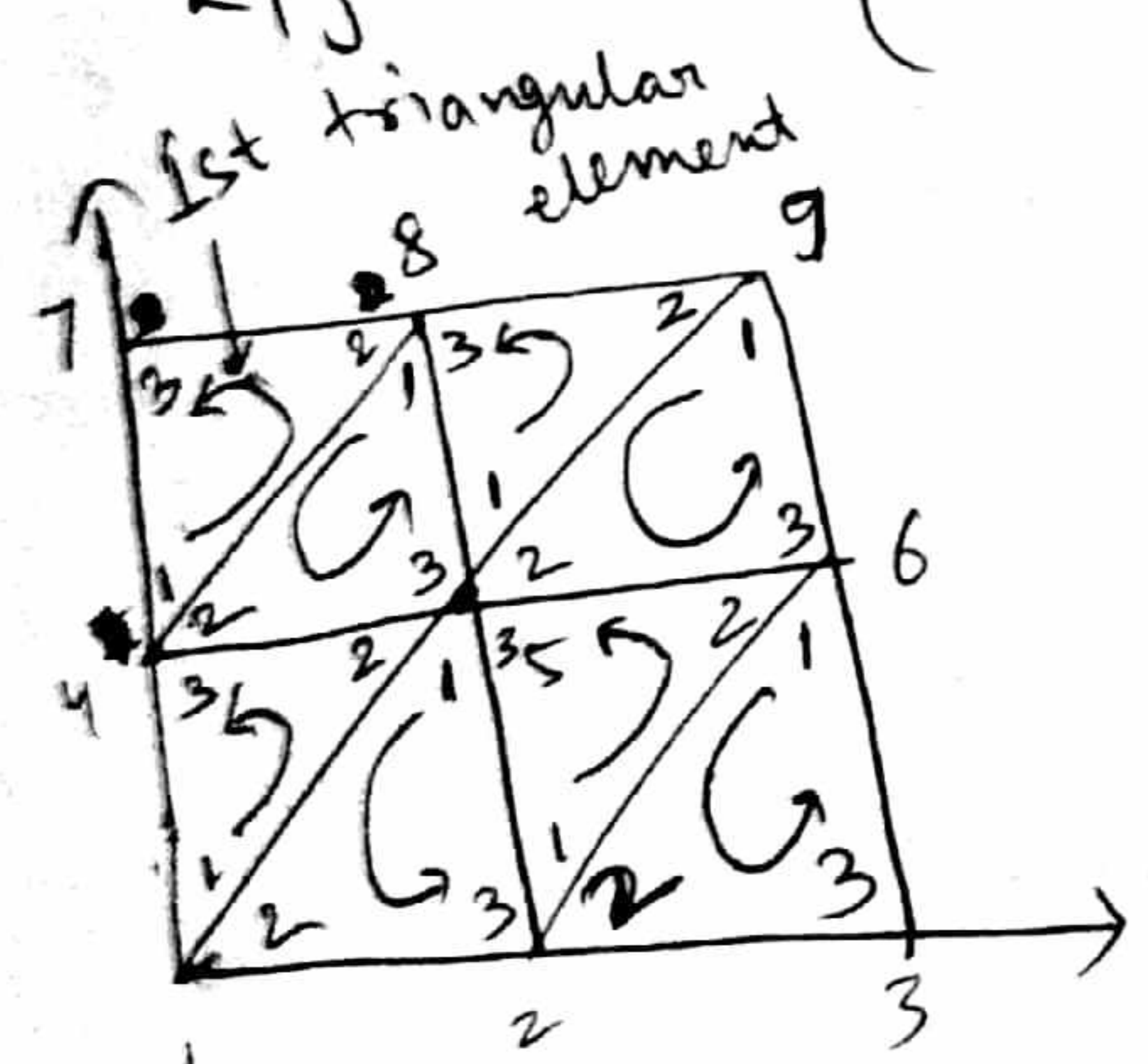
$$3 \rightarrow 1 \rightarrow 2$$
  

$$i \quad j \quad k$$

$$\alpha_1 + \alpha_2 + \alpha_3 = 2A$$

$$k_{ij} = a \int_{\Omega} \left( \frac{\beta_i}{2A} * \frac{\beta_j}{2A} + \frac{\gamma_i}{2A} \frac{\gamma_j}{2A} \right) dx dy$$

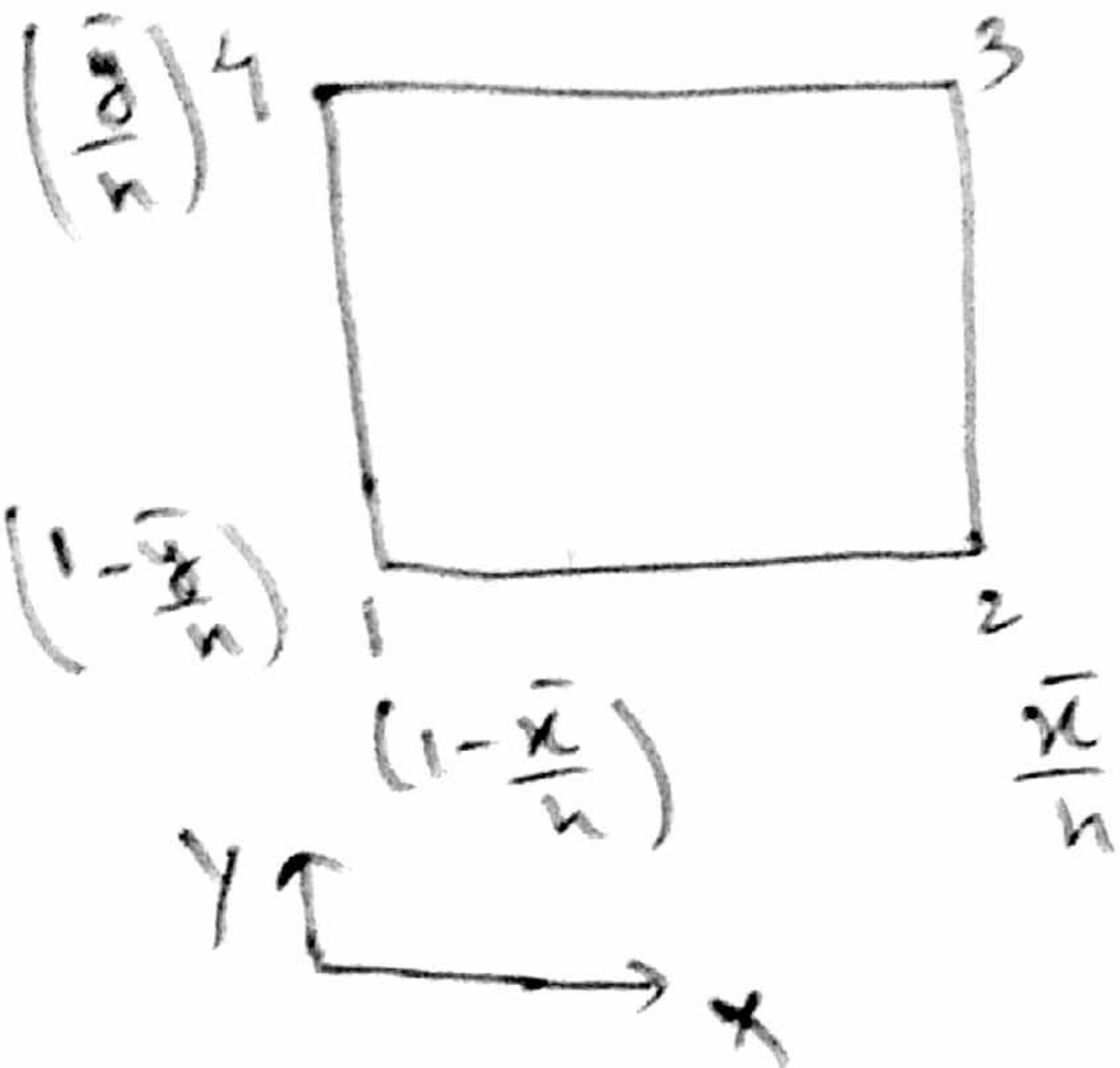
$$k_{ij} = a \left( \frac{\beta_i \beta_j + \gamma_i \gamma_j}{4A} \right)$$



1st element numbering in local coordinates was chosen arbitrarily but for other ~~some~~ <sup>successive</sup> elements we need to follow the 1st element numbering scheme otherwise the area computation will be wrong.

Linear rectangular element

Local coordinates



$$\Psi_4 = \left(1 - \frac{\bar{x}}{h}\right) \frac{\bar{y}}{h}$$

$$\Psi_1 = \left(1 - \frac{\bar{x}}{h}\right) \left(1 - \frac{\bar{y}}{h}\right)$$

$$\Psi_2 = \frac{\bar{x}}{h} \left(1 - \frac{\bar{y}}{h}\right)$$

$$\Psi_3 = \frac{\bar{x}}{h} \frac{\bar{y}}{h}$$

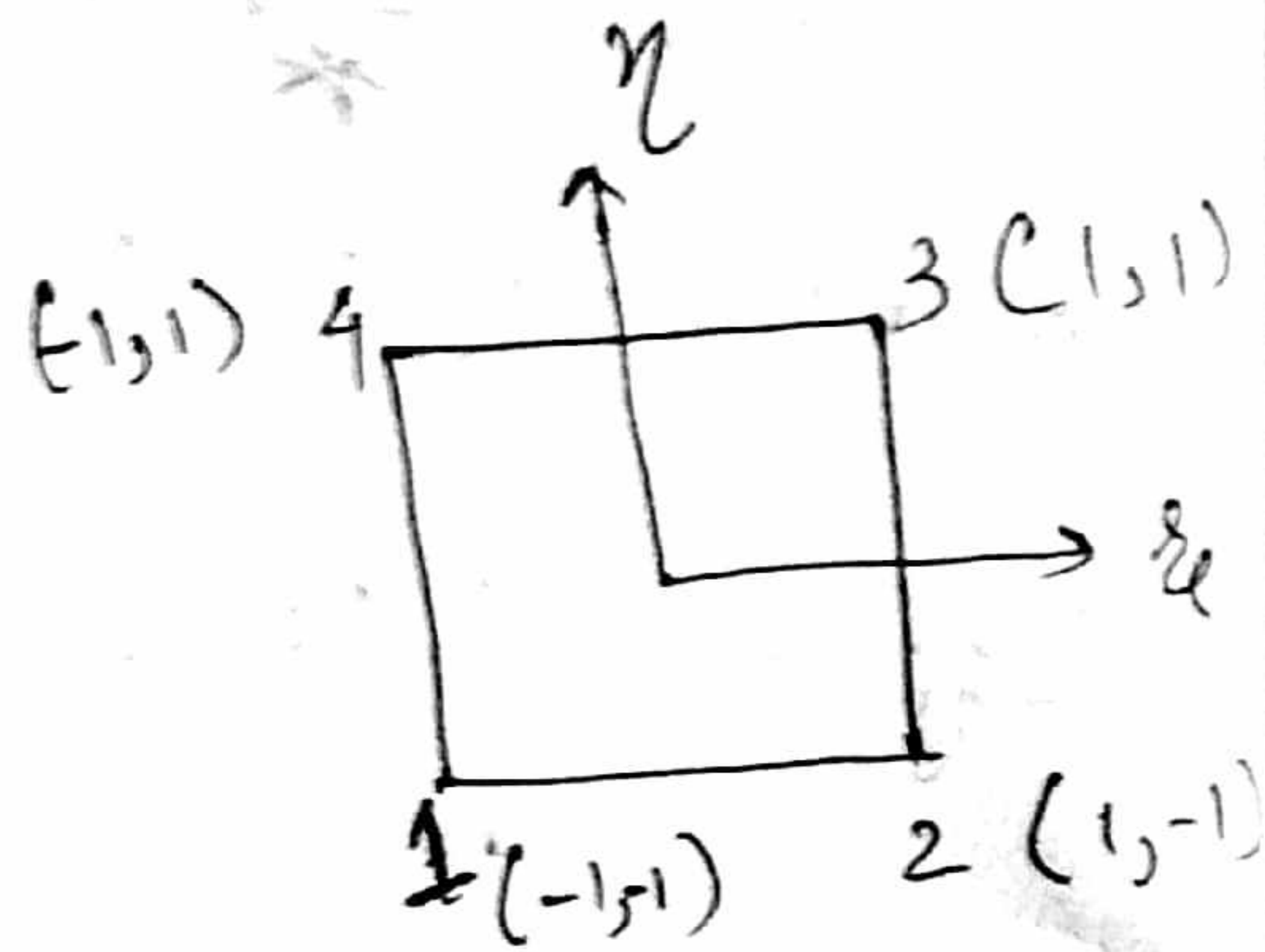
Natural coordinates

$$\Psi_4 = \left(\frac{1-\xi}{2}\right) \left(\frac{1+\eta}{2}\right)$$

$$\Psi_1 = \left(\frac{1-\xi}{2}\right) \left(\frac{1-\eta}{2}\right)$$

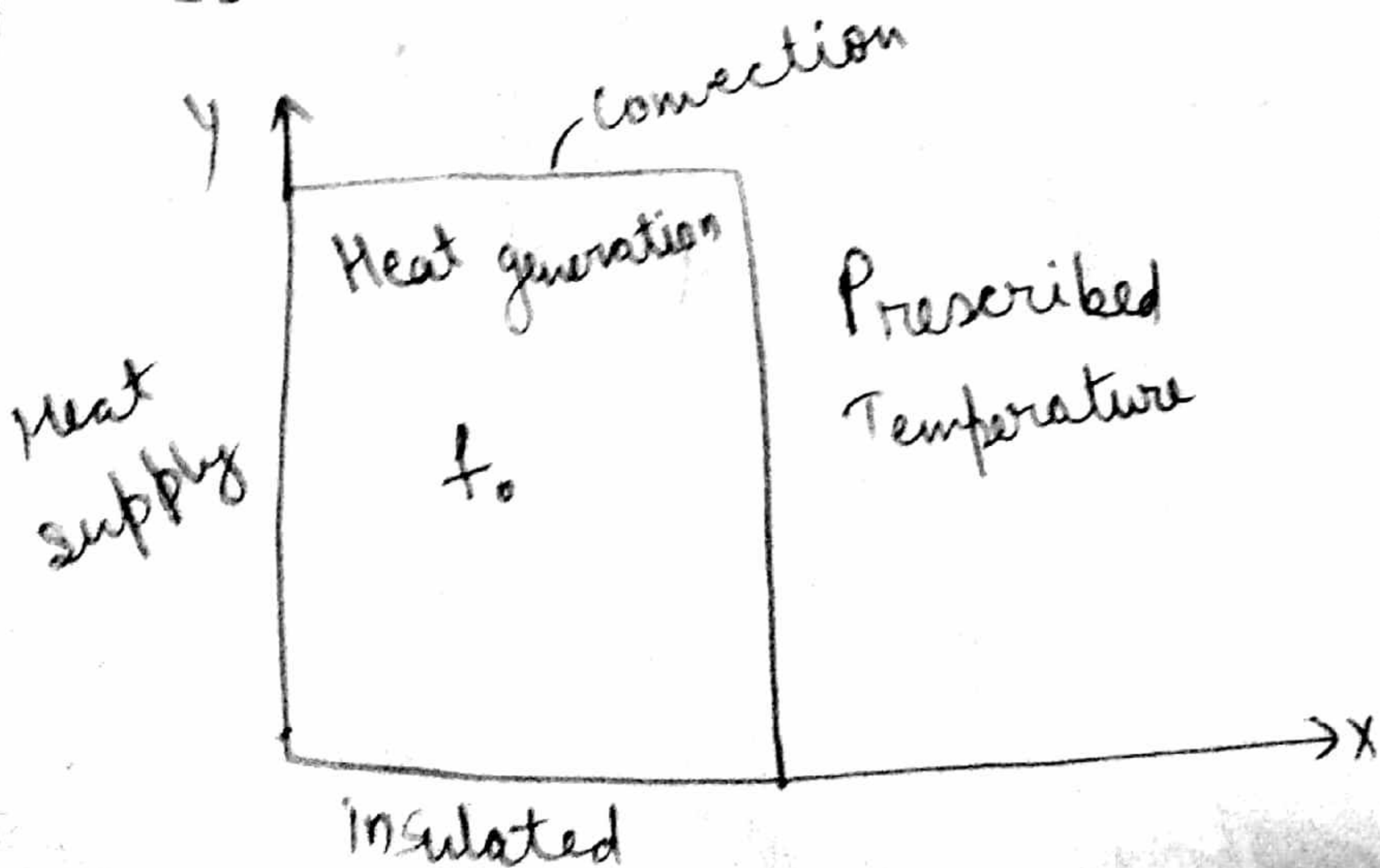
$$\Psi_2 = \left(\frac{1+\xi}{2}\right) \left(\frac{1-\eta}{2}\right)$$

$$\Psi_3 = \left(\frac{1+\xi}{2}\right) \left(\frac{1+\eta}{2}\right)$$



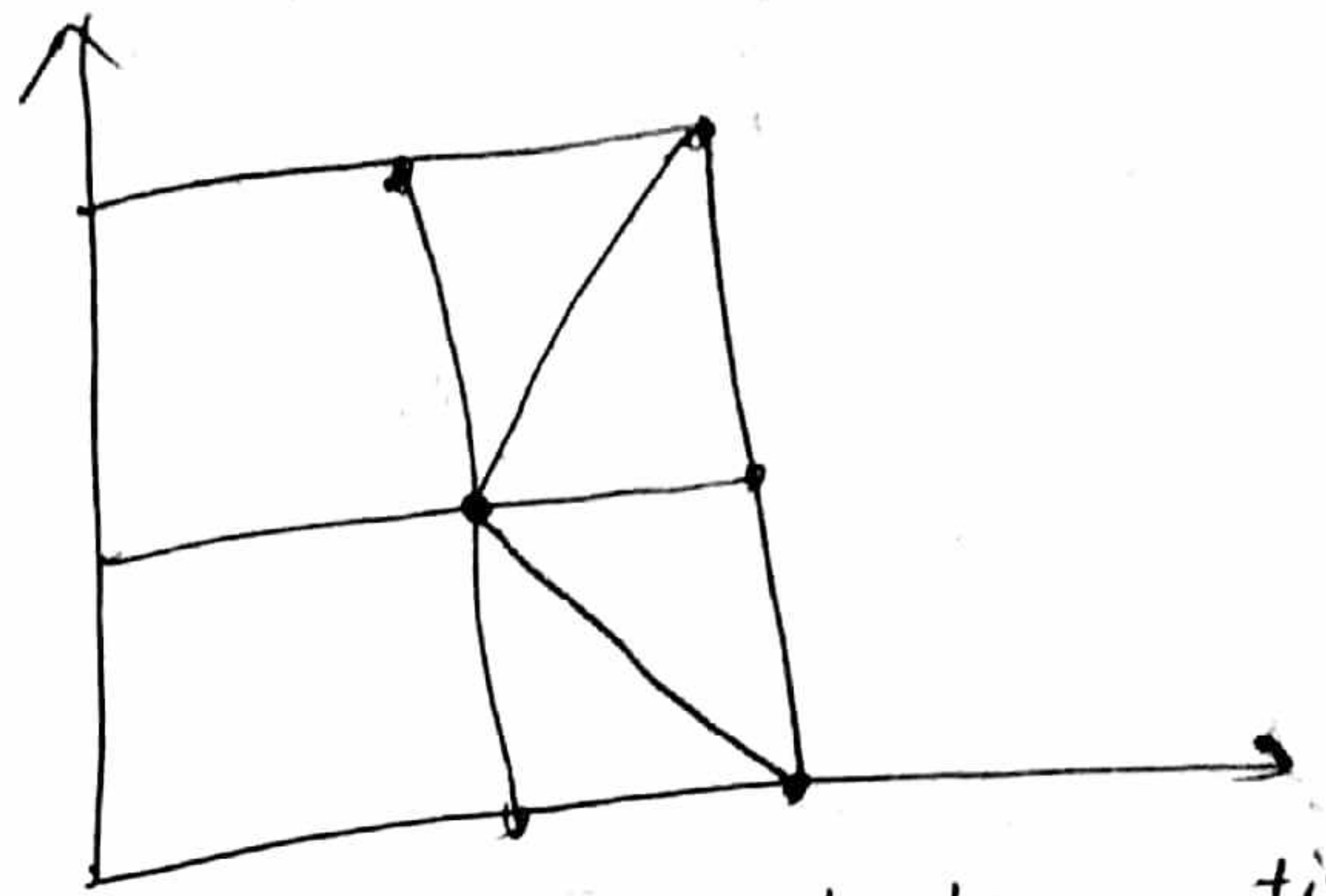
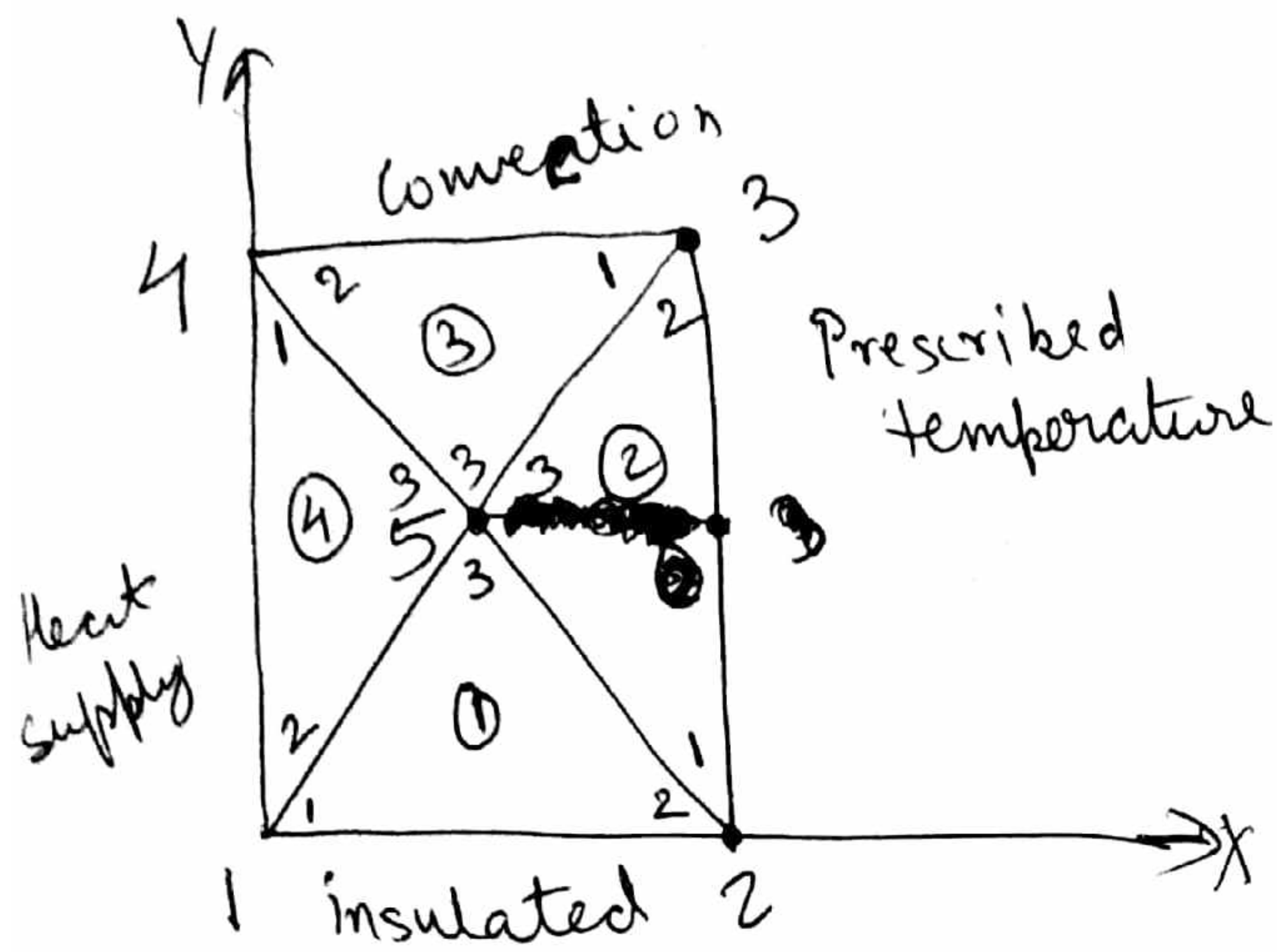
Class of interpolation functions belonging to Lagrange interpolation function:

$$\int_{\Omega} (\Psi_1^m \Psi_2^n \Psi_3^k) dx dy = \frac{m! n! k!}{(m+n+k+2)!} 2A$$



$$-\frac{\partial}{\partial x} \left( k_x \frac{\partial T}{\partial x} \right) - \frac{\partial}{\partial y} \left( k_y \frac{\partial T}{\partial y} \right) = f(x, y)$$

$$k_x \frac{\partial T}{\partial x} \hat{n}_x + k_y \frac{\partial T}{\partial y} \hat{n}_y + \beta(T - T_\infty) = \hat{q}_n \leftarrow \text{convective boundary}$$



This type of discretisation is also allowed.

$$\int_{\Omega} \left( k_x \frac{\partial w}{\partial x} \frac{\partial T}{\partial x} + k_y \frac{\partial w}{\partial y} \frac{\partial T}{\partial y} \right) dx dy = \int_{\Omega} w f dx dy + \int_s w \left( k_x \frac{\partial T}{\partial x} \hat{n}_x + k_y \frac{\partial T}{\partial y} \hat{n}_y \right) ds$$

$$= \int_{\Omega} w f dx dy + \int_s w [\beta(T - T_\infty)] ds$$

$$\int_{\Omega} \left( k_x \frac{\partial \psi_i}{\partial x} \frac{\partial}{\partial x} \sum_{j=1}^3 \psi_j T_j + k_y \frac{\partial \psi_i}{\partial y} \frac{\partial}{\partial y} \sum_{j=1}^3 \psi_j T_j \right) dx dy = \int_{\Omega} \psi_i f dx dy + \int_s \psi_i \left[ \beta \left( \sum_{j=1}^n \psi_j T_j - T_\infty \right) \right] ds$$

$$\psi_i = \frac{1}{2A} (\alpha_i + \beta_i x + \gamma_i y)$$

$$\frac{\partial \psi_i}{\partial x} = \frac{1}{2A} \beta_i$$

$$\frac{\partial \psi_i}{\partial y} = \frac{\gamma_i}{2A}$$

$$k_{ij} = \int_{\Omega} \left( k_x \frac{\beta_i}{2A} \frac{\beta_j}{2A} + k_y \frac{\gamma_i}{2A} \frac{\gamma_j}{2A} \right) dx dy$$

$$k_{ij} = \int_{\Omega} \frac{1}{4A^2} (k_x \beta_i \beta_j + k_y \gamma_i \gamma_j) dx dy$$

$$k_{ij} = \frac{1}{4A} (k_x \beta_i \beta_j + k_y \gamma_i \gamma_j)$$

$\therefore \beta_i, \beta_j, \gamma_i, \gamma_j, k_x, k_y$  are constants Also,  $\int_{-a/2}^{a/2} dx dy = A$

$$\alpha_i = x_j y_k - x_k y_j$$

$$\beta_i = (y_j - y_k)$$

$$\gamma_i = -(x_j - x_k)$$

$$\alpha_1 + \alpha_2 + \alpha_3 = 2A$$

for the first element

$$\beta_1 = (0 - \frac{a}{2}) = -\frac{a}{2}$$

$$\beta_2 = (\frac{a}{2} - 0) = \frac{a}{2}$$

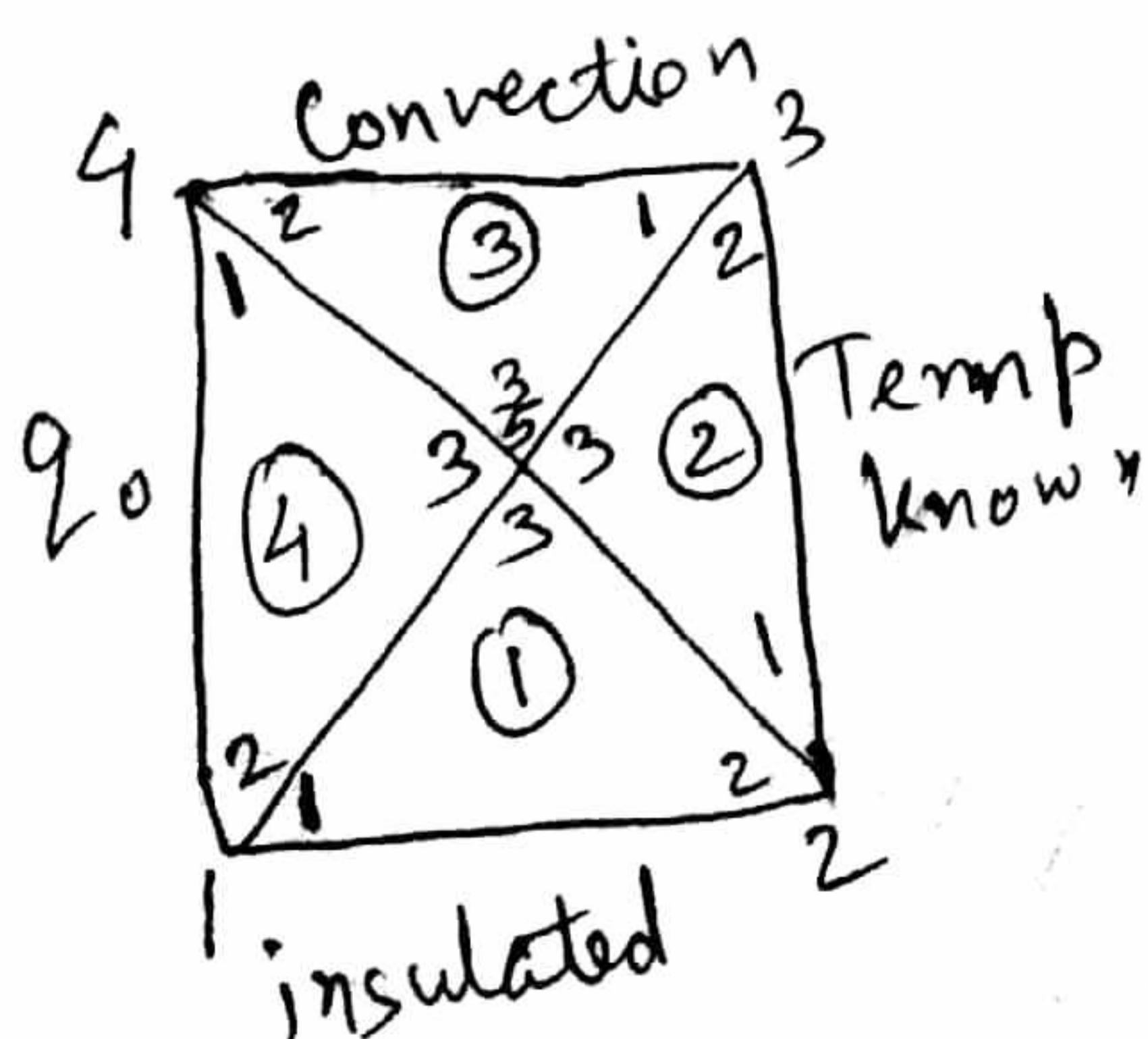
$$\beta_3 = (0 - 0) = 0$$

$$\gamma_1 = -(a - \frac{a}{2}) = -\frac{a}{2}$$

$$\gamma_2 = -(\frac{a}{2} - 0) = -\frac{a}{2}$$

$$\gamma_3 = -(0 - a) = a$$

$$A = \frac{1}{2} \times a \times \frac{a}{2} = \frac{a^2}{4}$$



Element

Local

Global

1

1-2-3

1-2-5

2

1-2-3

2-3-5

3

1-2-3

3-4-5

4

1-2-3

4-1-5

side EBC,  $T_2 = T_3 = 100$

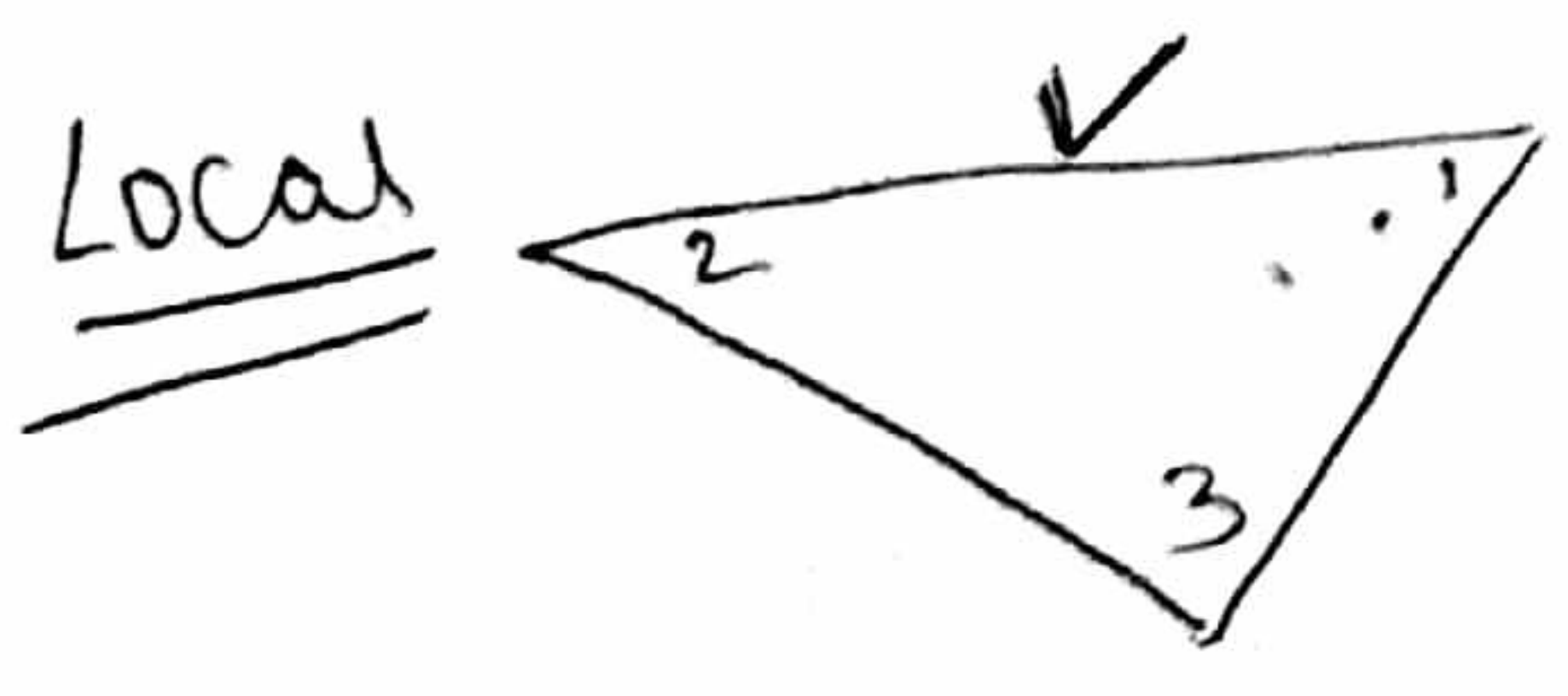
NBC - Side 1-2:  $k_x \frac{dT}{dx} \uparrow_x + k_y \frac{dT}{dy} \uparrow_y$  as insulated

Side 3-4  $\oint W [q_n - \beta (T - T_0)] ds \equiv \oint \psi_i \beta T_0 ds$  and

$q_n =$  External heat source on boundary 3-4  $\int \psi_i \beta T_0 ds$

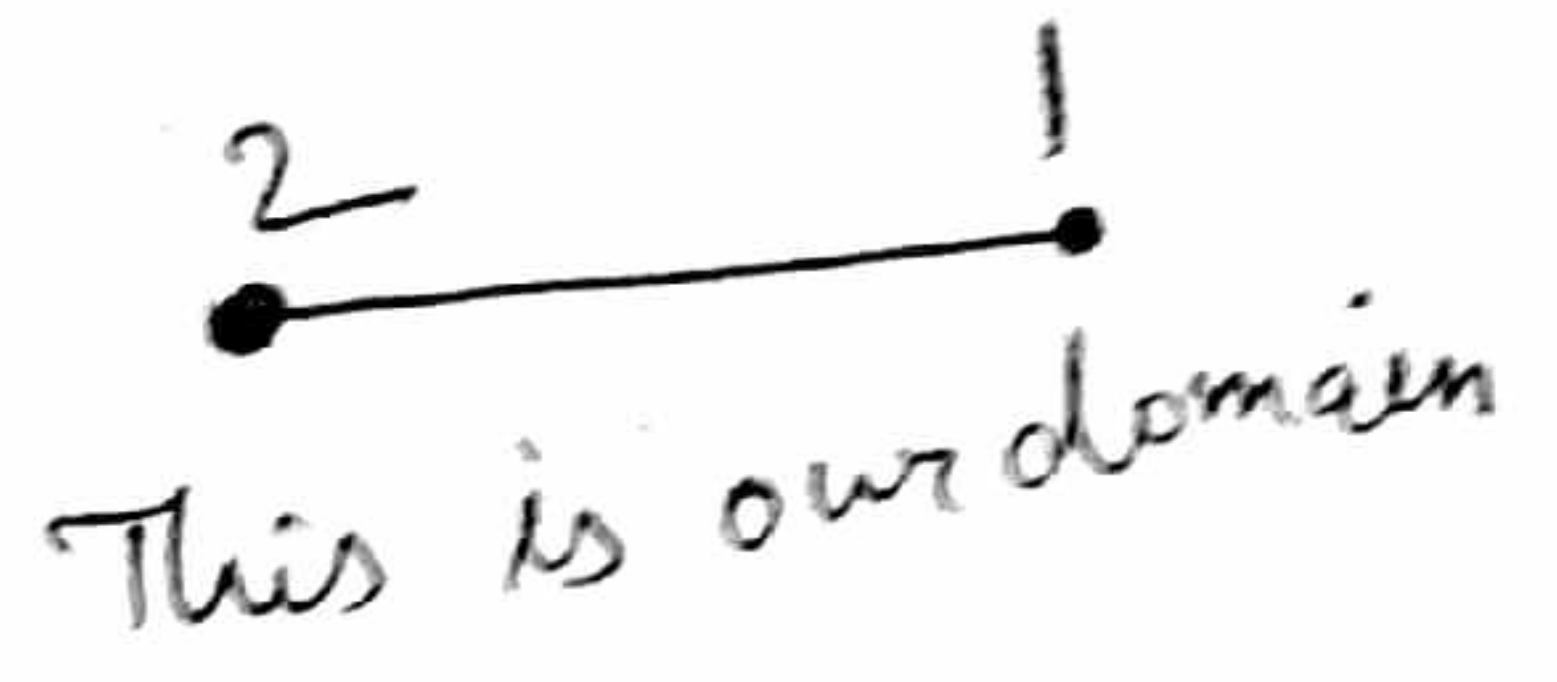
side 4-1 heat addition  $\equiv \oint \psi_i q \cdot ds$

Other integral internal heat generation  
 $f_i = \int_{\Omega} w T_0 dx dy = \int_{\Omega} \psi_i f_0 dx dy$  } heat generation inside the body



$$\begin{bmatrix} k_{11} & k_{12} & k_{13} \\ k_{21} & k_{22} & k_{23} \\ k_{31} & k_{32} & k_{33} \end{bmatrix} \begin{bmatrix} T_1 \\ T_2 \\ T_3 \end{bmatrix} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} \cdot \begin{bmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \end{bmatrix}$$

To find 'f' matrix



$\therefore f$  is on the line in this case.  
 Hence  $\psi$ 's to be put for calculating

$f$  will be of linear element Hence

$$f_1 = \int \psi_1 \beta T_0 ds = \frac{S\alpha}{2}$$

$$f_2 = \int \psi_2 \beta T_0 ds = \frac{S\alpha}{2}$$

$$f_3 = \int \psi_3 \beta T_0 ds = 0$$

$$d = \beta T_0$$

$$\psi_1 = 1 - \frac{s}{h}$$

$$\psi_2 = \frac{s}{h}$$

If  $3 \rightarrow 1$  was our domain then

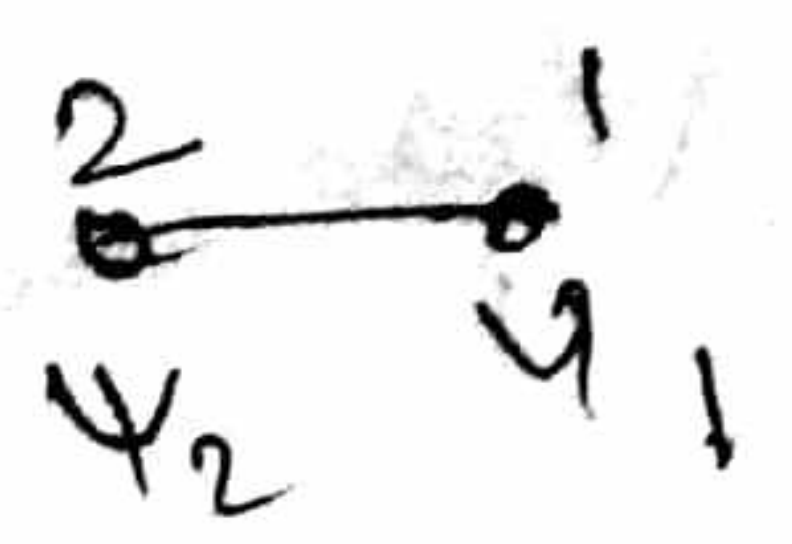
$$\psi_1 = 1 - \frac{s}{h}, \quad \psi_3 = \frac{s}{h}$$

$$f_1 = \int \psi_1 \beta T_0 ds = \frac{S\alpha}{2}$$

$$f_2 = \int \psi_2 \beta T_0 ds = 0$$

$$f_3 = \int \psi_3 \beta T_0 ds = \frac{S\alpha}{2}$$

$$\oint \psi_i \beta \sum_{j=1}^3 T_j \psi_j ds$$



$$\begin{bmatrix}
 c_{11} & c_{12} & 0 \\
 \int \psi_1 \psi_1 & \int \psi_1 \psi_2 & \int \psi_1 \psi_3 \rightarrow 0 \\
 c_{21} & c_{22} & 0 \\
 \int \psi_1 \psi_2 & \int \psi_2 \psi_2 & \int \psi_2 \psi_3 \rightarrow 0 \\
 0 & 0 & 0 \\
 \int \psi_1 \psi_3 & \int \psi_2 \psi_3 & \int \psi_3 \psi_3 \rightarrow 0
 \end{bmatrix}
 \begin{Bmatrix}
 T_1 \\
 T_2 \\
 T_3
 \end{Bmatrix}$$

$$\begin{bmatrix}
 k_{11} + c_{11} & k_{12} + c_{12} & k_{13} \\
 k_{21} + c_{21} & k_{22} + c_{22} & k_{23} \\
 k_{31} & k_{32} & k_{33}
 \end{bmatrix}
 \begin{Bmatrix}
 T_1 \\
 T_2 \\
 T_3
 \end{Bmatrix}
 =
 \begin{Bmatrix}
 f_1 \rightarrow \frac{\beta T_0 h}{2} \\
 f_2 \rightarrow \frac{\beta T_0 h}{2} \\
 f_3 \rightarrow 0
 \end{Bmatrix}$$

$$\begin{bmatrix}
 c_{11} & c_{12} \\
 c_{21} & c_{22}
 \end{bmatrix}
 = \frac{\beta h}{6} \begin{bmatrix}
 2 & 1 \\
 1 & 2
 \end{bmatrix}$$

for  $q_0$  boundary

$$\begin{aligned}
 f_1 &\rightarrow \frac{q_0 h}{2} \\
 f_2 &\rightarrow \frac{q_0 h}{2} \\
 f_3 &\rightarrow 0
 \end{aligned}$$

for internal heat generation boundary condition.

$$\int_{\Omega} \psi_i f_0 dx dy = \frac{f_0 A}{3}$$

As it is integration over domain.  $\psi_i$  here is triangular element interpolation function on

$$\int_{\Omega} \psi^m \psi^n \psi^p d\Omega = \frac{m! n! p!}{(m+n+p+2)!} \times 2A$$

Element  $k_{ij} = \frac{k_x \beta_i \beta_j + k_y \beta_i \beta_j}{4A}$

$$[k] = \frac{k}{2} \begin{bmatrix}
 1 & 0 & -1 \\
 0 & 1 & -1 \\
 -1 & -1 & 2
 \end{bmatrix}$$

↑  
local



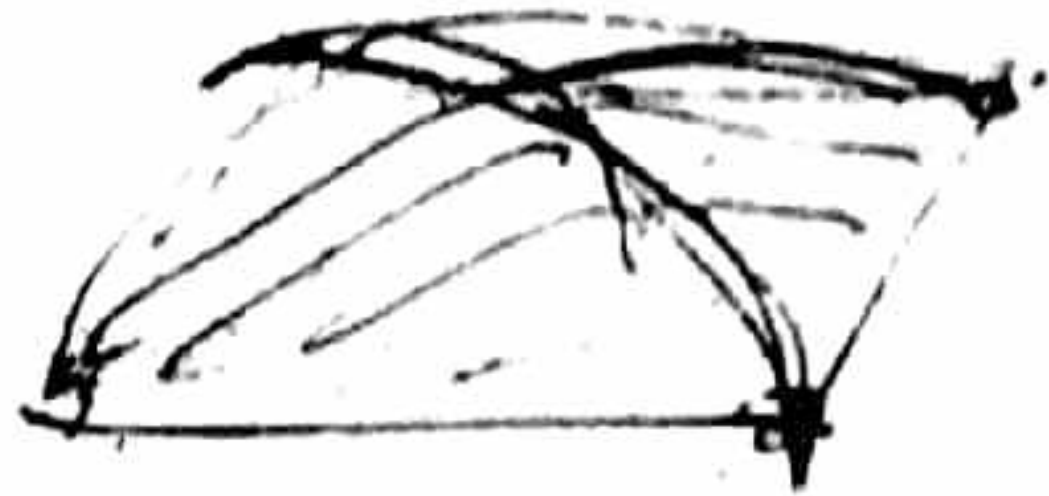
|   | 1                     | 2                     | 3                     | 4                     | 5   |
|---|-----------------------|-----------------------|-----------------------|-----------------------|---|
| 1 | $k_{11}^1 + k_{22}^4$ | $k_{12}^1$            | 0                     | $k_{21}^4$            | $k_{13}^1 + k_{23}^4$                       |
| 2 | $k_{21}^1$            | $k_{22}^1 + k_{11}^2$ | $k_{12}^2$            | 0                     | $k_{23}^1 + k_{13}^2$                       |
| 3 | 0                     | $k_{21}^2$            | $k_{22}^2 + k_{11}^3$ | $k_{12}^3$            | $k_{23}^2 + k_{13}^3$                       |
| 4 | $k_{12}^4$            | 0                     | $k_{21}^3$            | $k_{11}^4 + k_{22}^3$ | $k_{13}^4 + k_{23}^3$                       |
| 5 | $k_{31}^1 + k_{32}^4$ | $k_{32}^1 + k_{31}^2$ | $k_{32}^2 + k_{31}^3$ | $k_{31}^4 + k_{32}^3$ | $k_{33}^1 + k_{33}^2 + k_{33}^4 + k_{33}^3$ |

$[k] =$   
↑  
Global

$$k_{33}^1 + k_{33}^2 + k_{33}^4 + k_{33}^3 = 4k$$

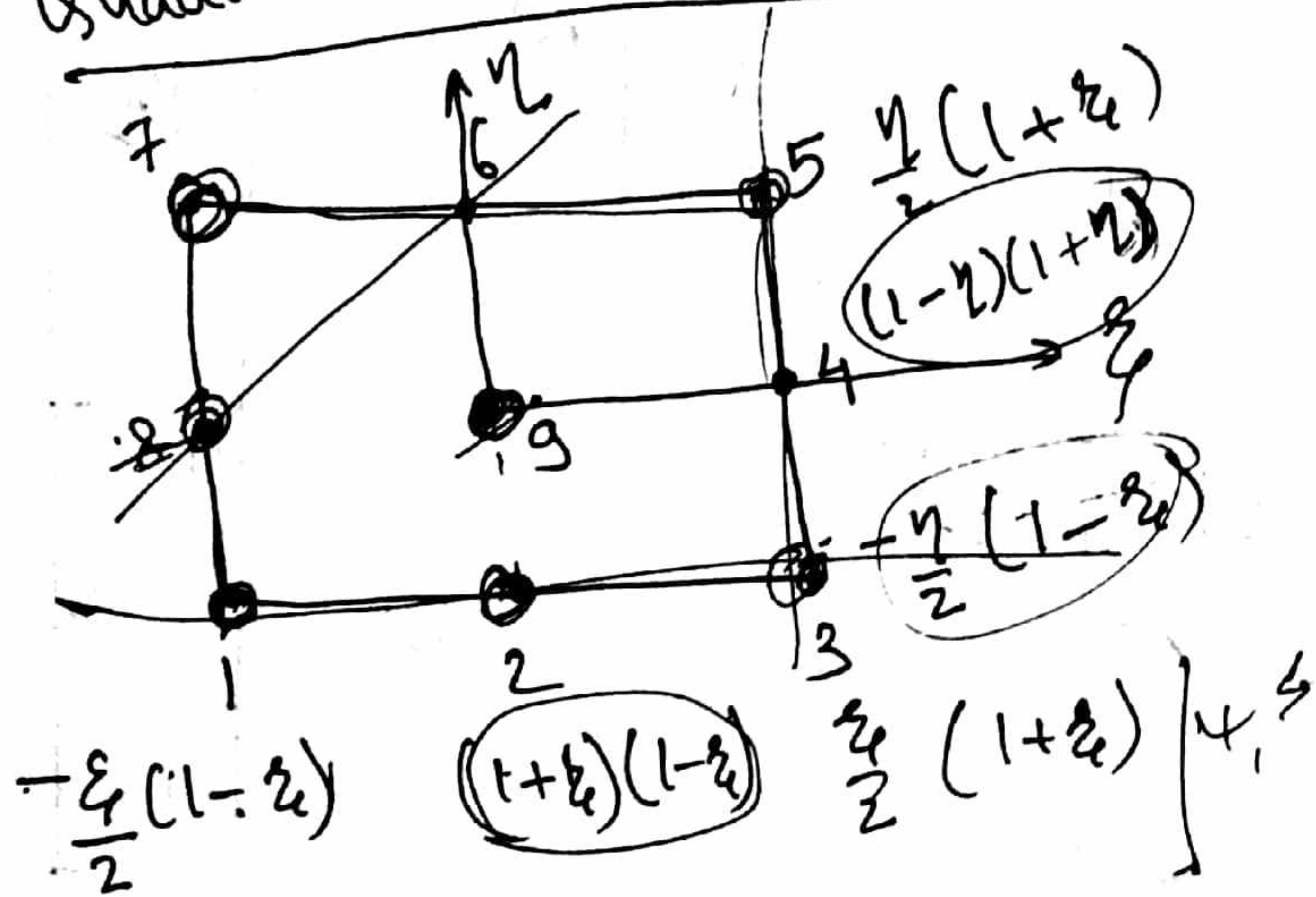
$$[k] \begin{matrix} \uparrow \\ \text{Global} \end{matrix} \begin{bmatrix} T_1 \\ T_2 \\ T_3 \\ T_4 \\ T_5 \end{bmatrix}$$

$$= \begin{cases} f_1^1 + f_2^4 \\ f_2^1 + f_1^2 \\ f_2^3 + f_1^3 \\ f_2^3 + f_1^4 \\ f_3^1 + f_3^2 + f_3^3 + f_3^4 \end{cases}$$

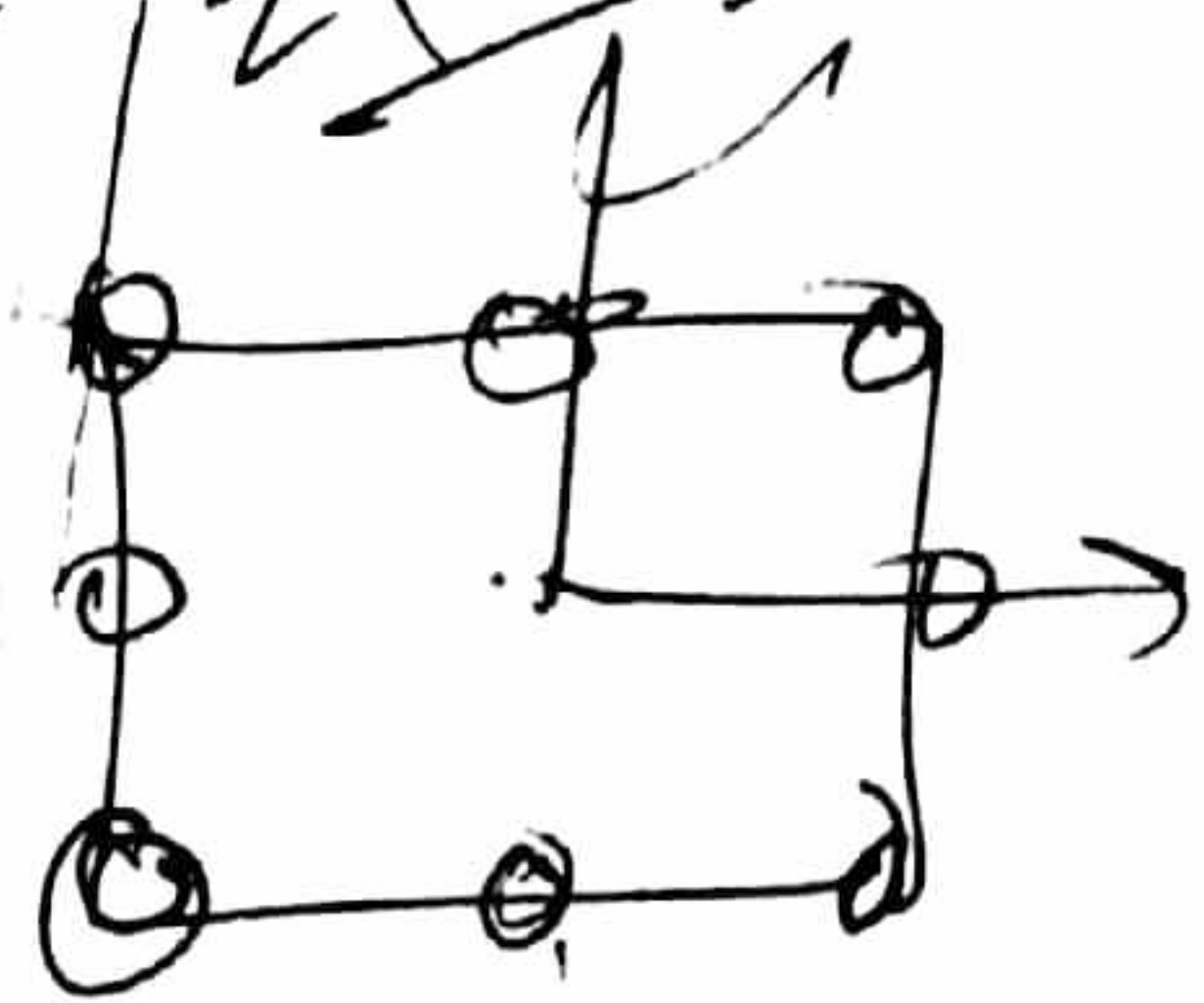


## Interpolation Function

### Quadratic interpolation (rectangular)



$$\Psi_8 = \frac{1}{8} (1-xi)(1-eta)(1+eta)$$



$$\Psi_9 = (1-xi)(1+xi)(1-eta)(1+eta)$$

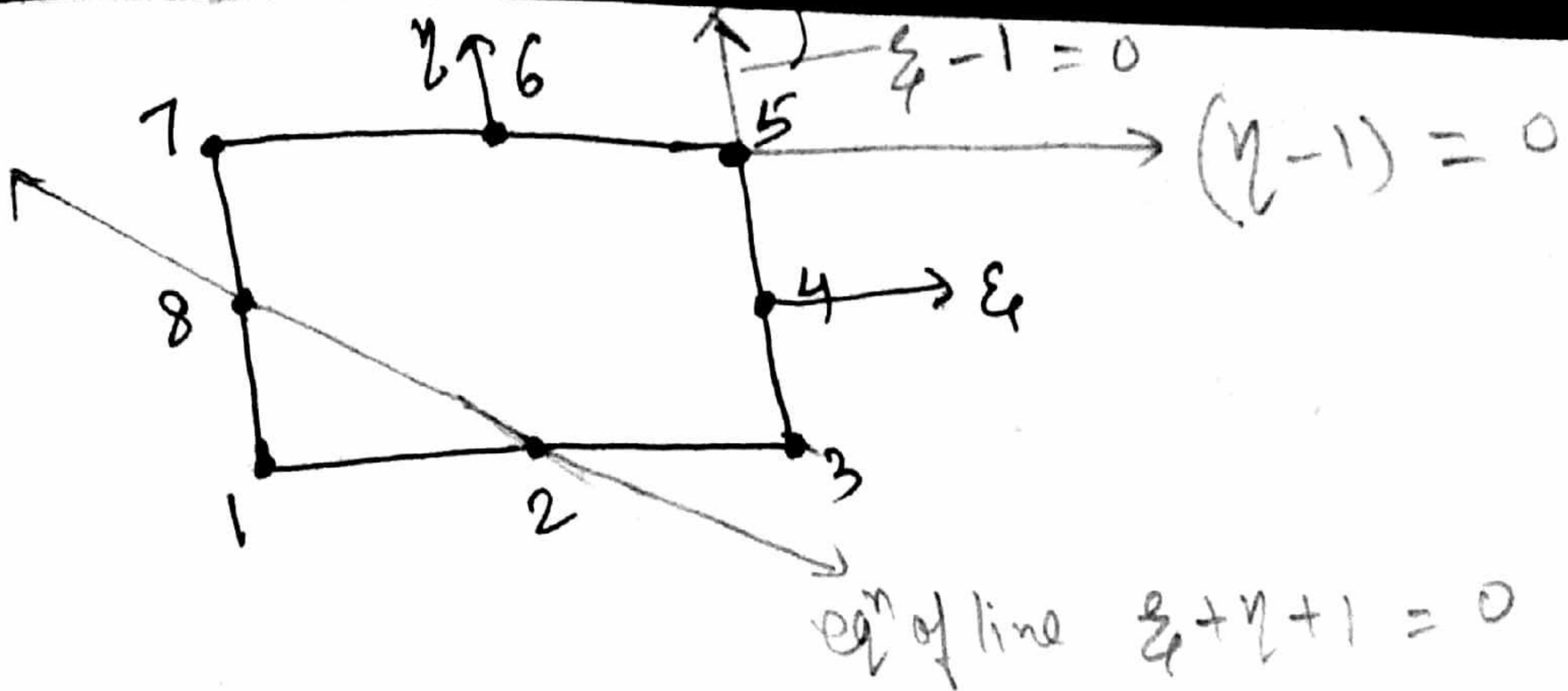
$$\Psi_9 = (1-xi^2)(1-eta^2)$$

$$\Psi_3 = \frac{xi}{2}(1+xi)\left(-\frac{eta}{2}\right)(1+eta) = -\frac{xi\eta}{4}(1+xi)(1+eta)$$

This particular with 8 node number element is very common in Ansys & abacus & serendipity element.

How do you evaluate the problem with 8 node points.

Polynomial method  $\rightarrow$  By using the properties of interpolation function we derive the interpolation function.



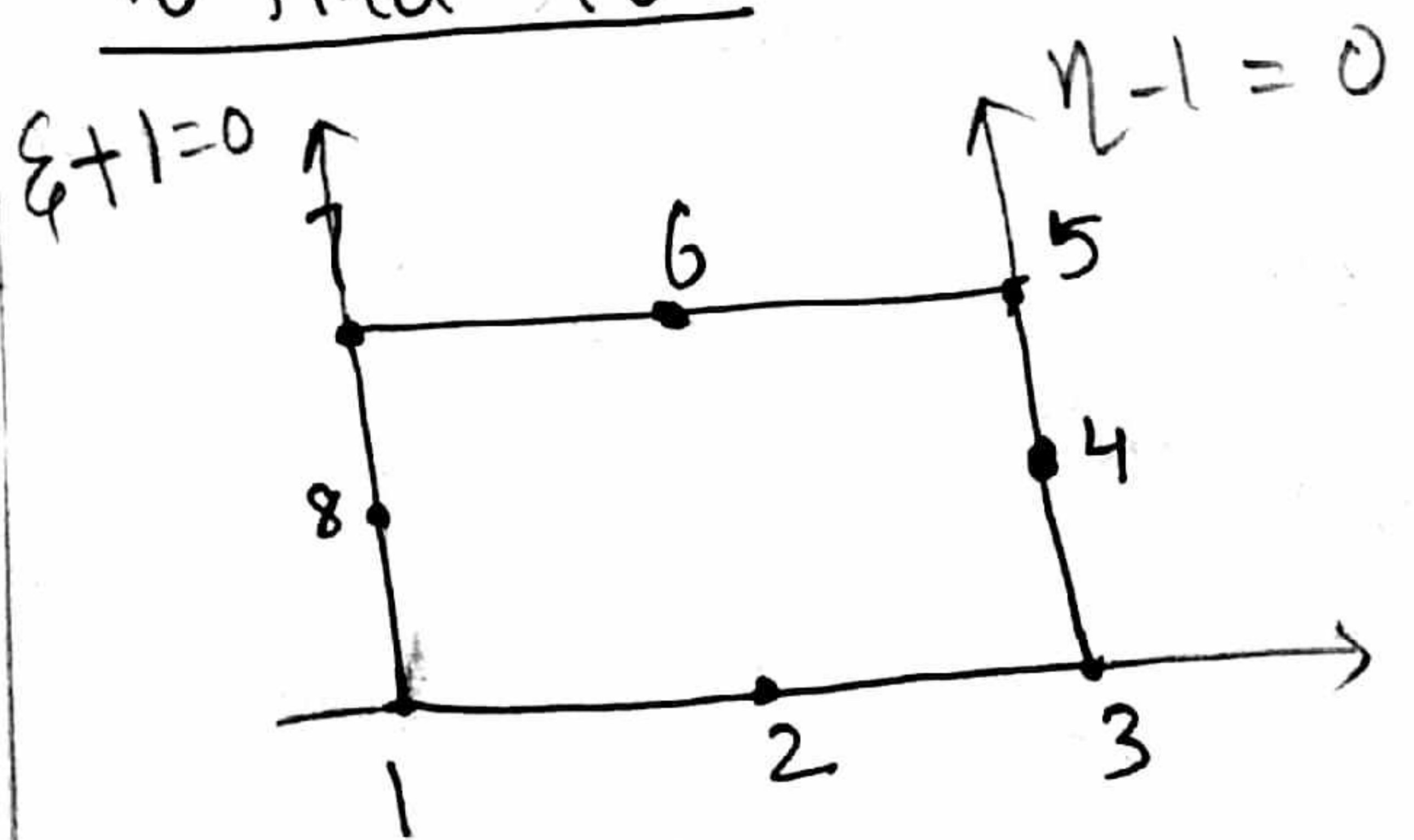
$$\psi_1 = c_1 (\xi + \eta + 1) (\xi - 1) (\eta - 1)$$

at  $(\xi, \eta) = (-1, -1)$

$$\psi_1(-1, -1) = 1 = c_1 (-1 - 1 + 1) (-1 - 1) (-1 - 1)$$

$$\Rightarrow c_1 = -\frac{1}{4}$$

To find  $\psi_6$



$$\psi_6 = c_6 (\xi + 1) (\eta + 1) (\eta - 1)$$

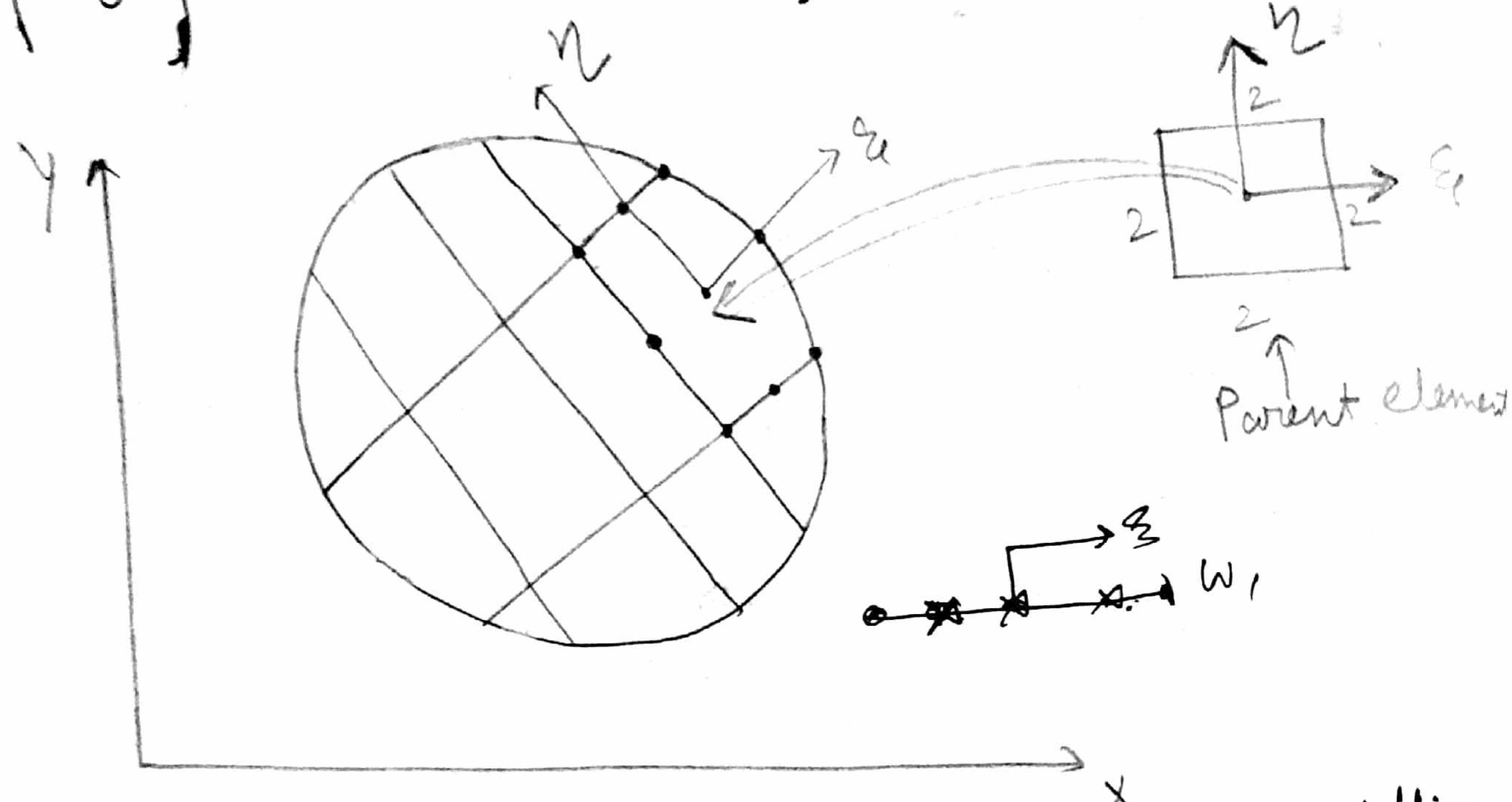
$$\psi_6(0, 1) = c_6 (1) (2) (-1)$$

$$c_6 = -\frac{1}{2}$$

$$\psi_6 = -\frac{1}{2} (\eta^2 - 1) (1 + \xi)$$

$$\begin{pmatrix} \frac{\eta c}{i h e} \\ \frac{\eta c}{h e} \\ \frac{\eta c}{i h e} \end{pmatrix} = \begin{pmatrix} \frac{\eta c}{h e} \\ \frac{\eta c}{h e} \\ \frac{\eta c}{i h e} \end{pmatrix} + \begin{pmatrix} \frac{\eta c}{i h e} \\ \frac{\eta c}{h e} \\ \frac{\eta c}{i h e} \end{pmatrix} = \frac{\eta c}{i h e}$$

$$\begin{Bmatrix} \frac{\partial \psi_i}{\partial x} \\ \frac{\partial \psi_i}{\partial y} \end{Bmatrix} = [J]^{-1} \begin{Bmatrix} \frac{\partial \psi_i}{\partial \xi} \\ \frac{\partial \psi_i}{\partial \eta} \end{Bmatrix}$$



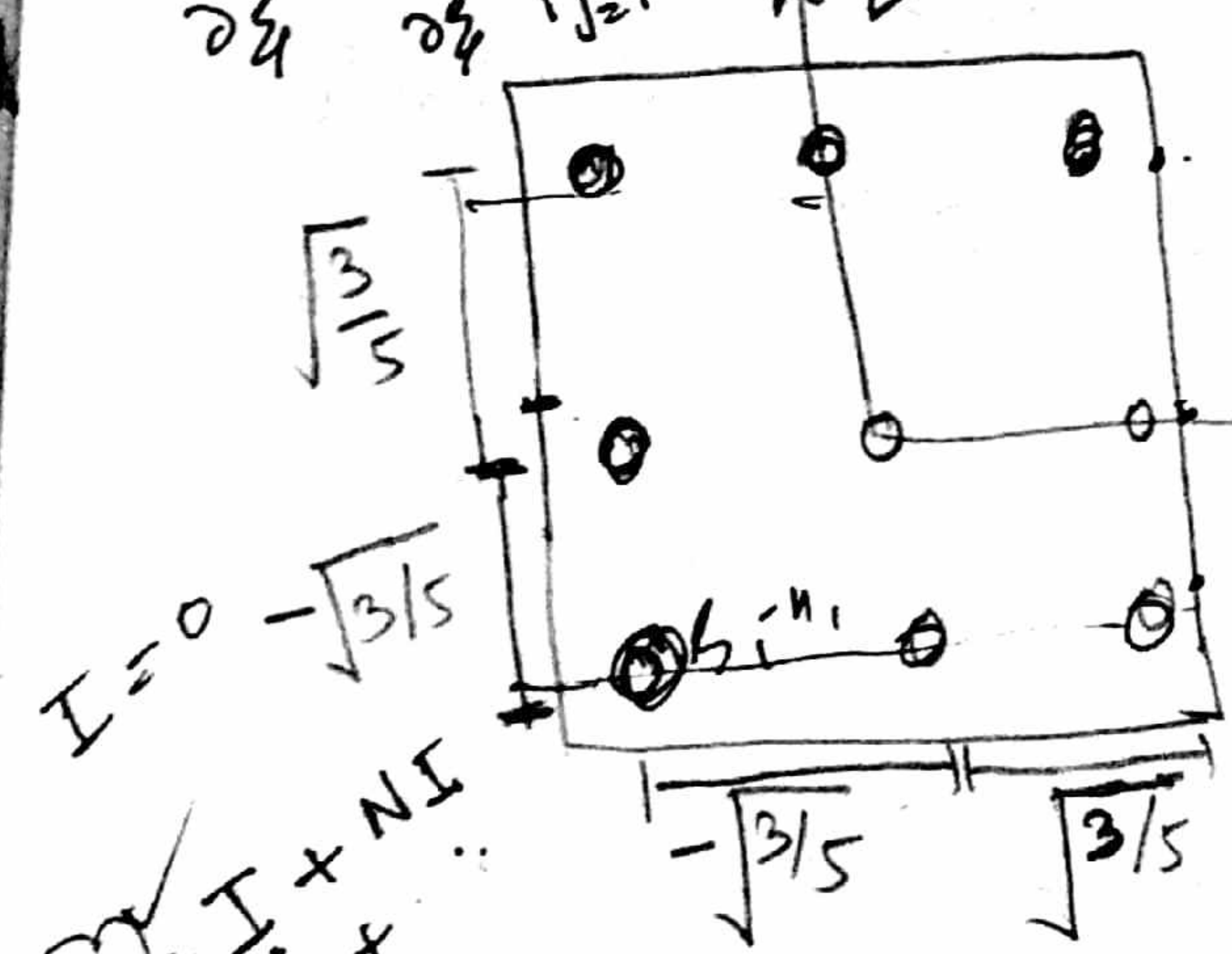
$$u = x_1 \psi_1 + x_2 \psi_2 + \dots + x_g \psi_g$$

$$y = y_1 \psi_1 + y_2 \psi_2 + \dots + y_g \psi_g$$

$$\frac{\partial y}{\partial \xi} = \frac{\partial}{\partial \xi} \left\{ \sum_{j=1}^n \psi_j(\xi, \eta) y_j \right\}$$

$$\int_{\xi_1}^{\xi_2} \int_{\eta_1}^{\eta_2} f(\xi, \eta) d\xi d\eta \approx \sum_{i=1}^n w_i f(\xi_i, \eta_i)$$

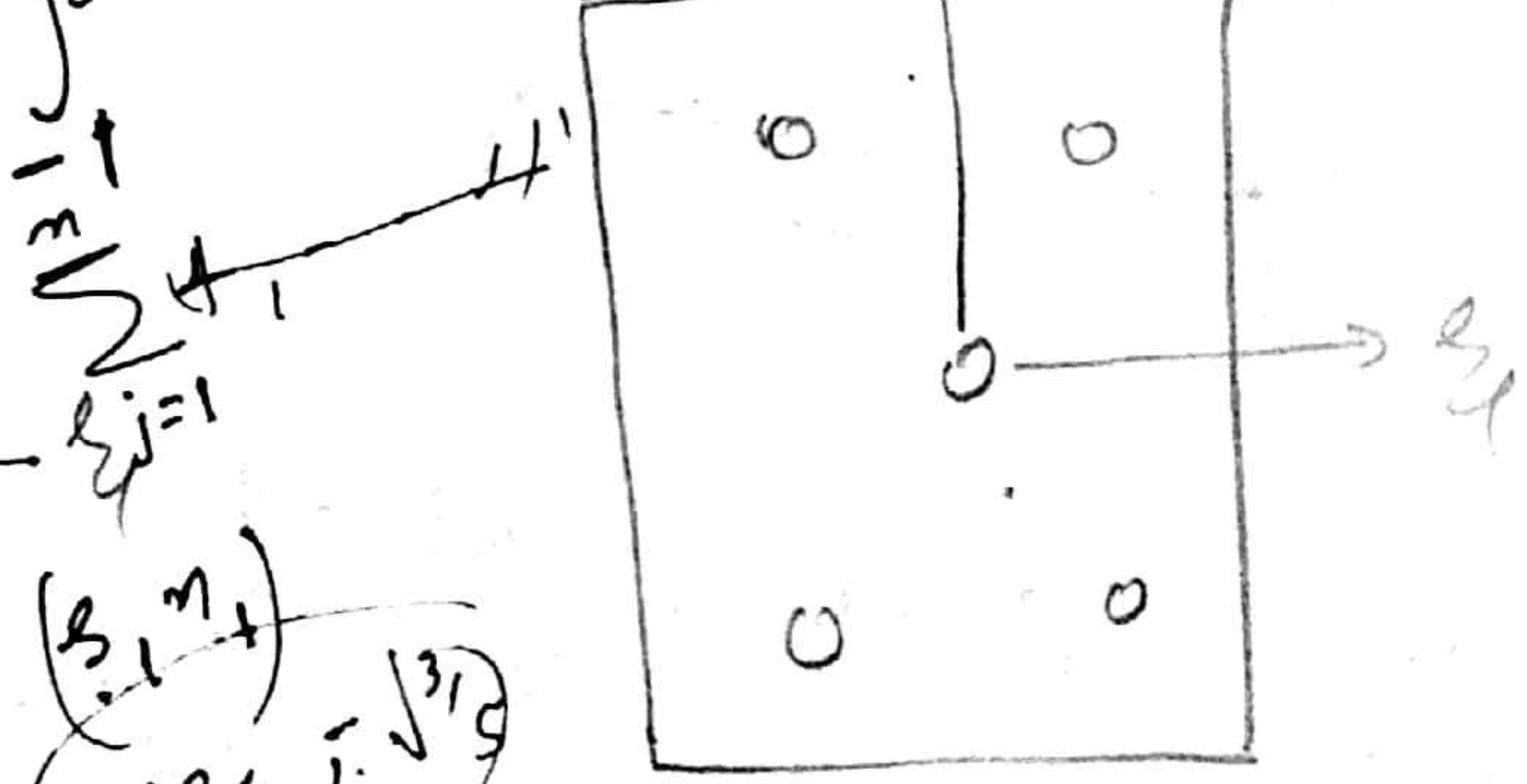
← Geometric modelling



$$I = \int_{-1}^1 \int_{-1}^1 f(\xi, \eta) d\xi d\eta$$

$$I \approx \sum_{i=1}^n w_i f(\xi_i, \eta_i)$$

3-point integration.



2-point integration

Functions are to be evaluated at  $\xi_i$  &  $\eta_j$  so integrand will be  $f(\xi_i, \eta_j) * w_i * w_j$

# Area coordinates

$$A_1 + A_2 + A_3 = A$$

$$A_1 = \frac{1}{2} b s$$

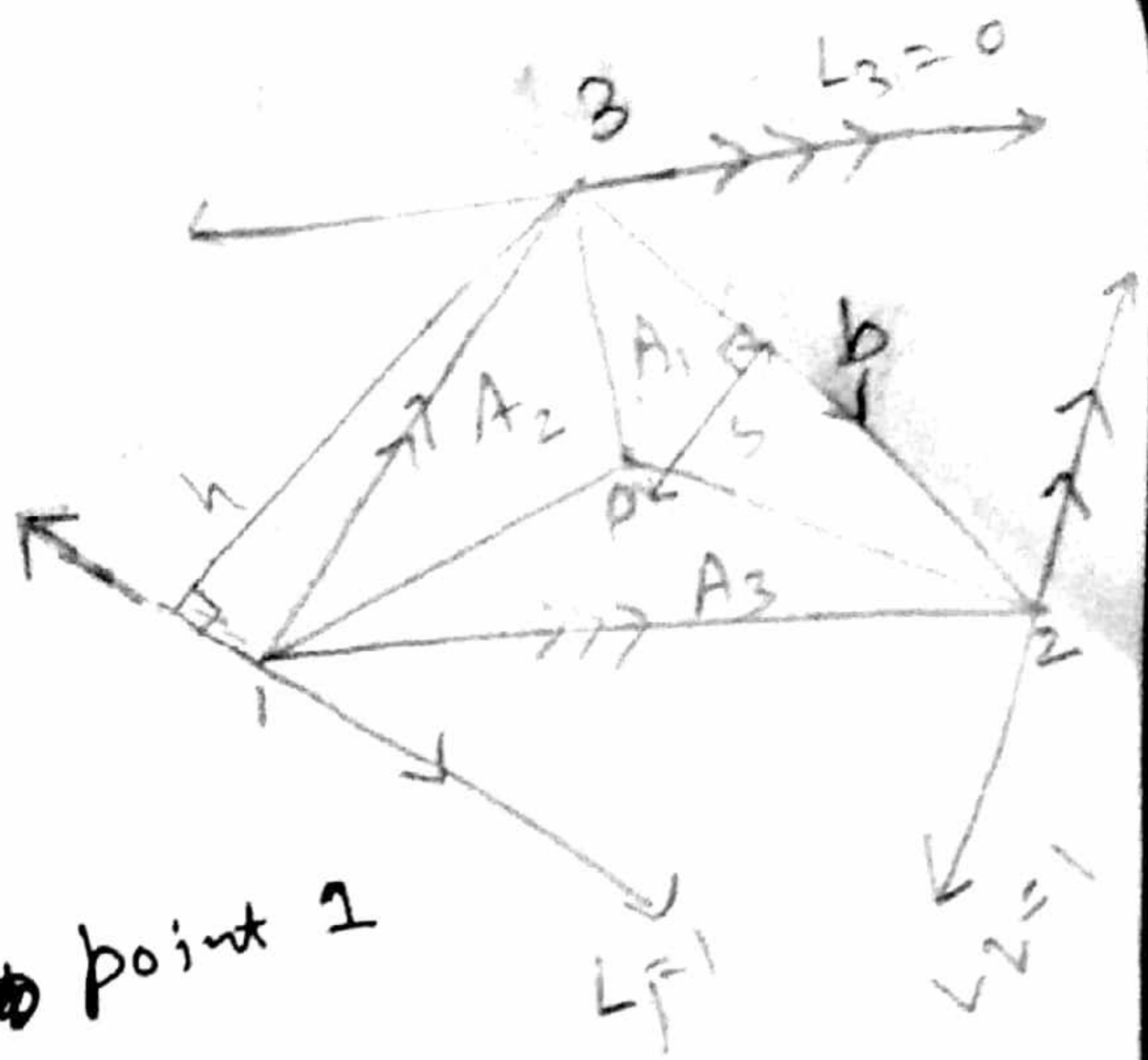
$$A = \frac{1}{2} b h$$

$$\frac{A_1}{A} = \frac{s}{h} = L_1$$

$h$  is  $\perp$  distance between 2-3 and point 1

$b$  = length of side 2-3

$s$  =  $\perp$  distance b/w side 2-3 and P



Quadratic element

$$\psi_1 = C_1 \left( L_1 - \frac{1}{2} \right) (L_1 - 0)$$

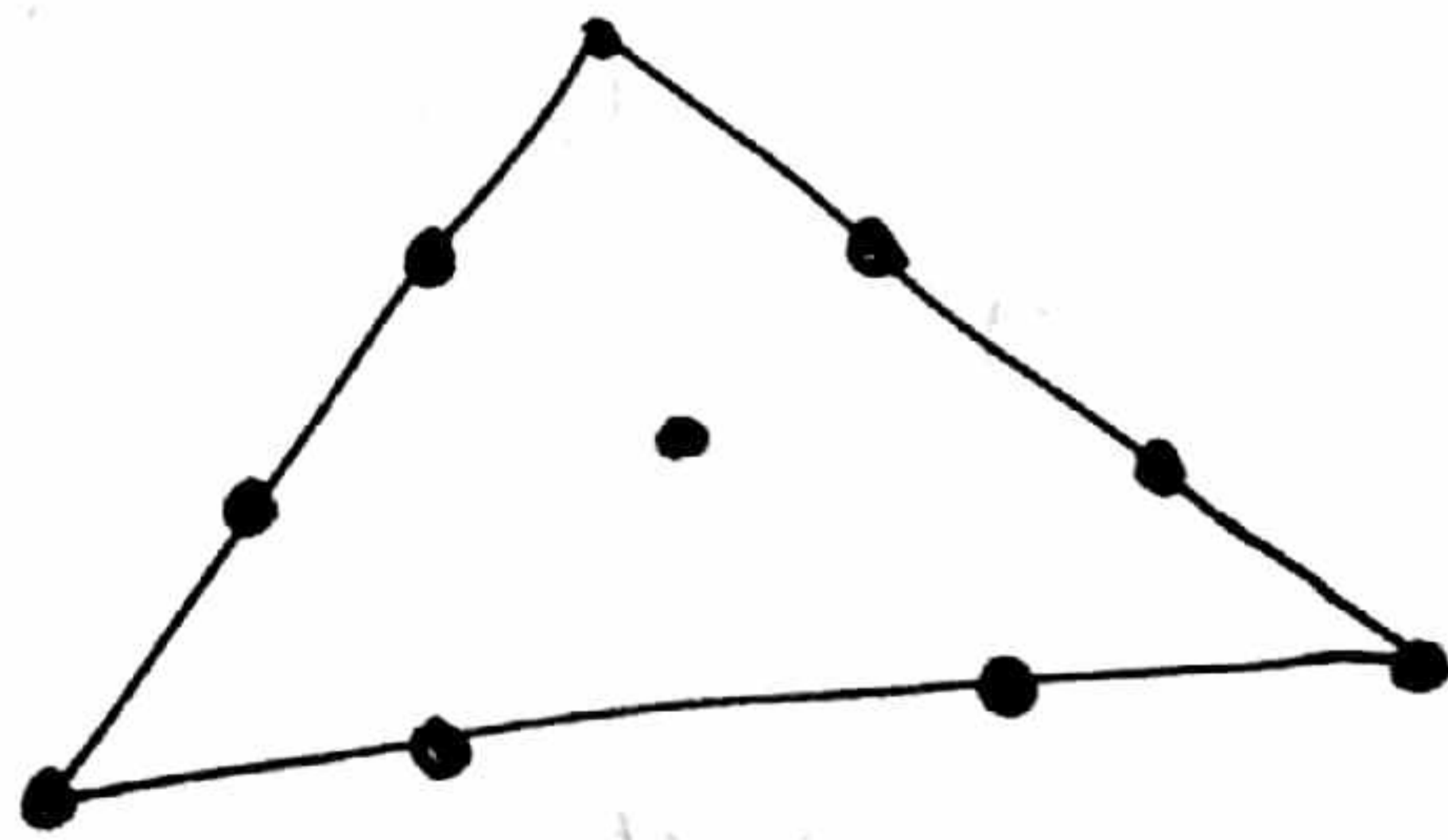
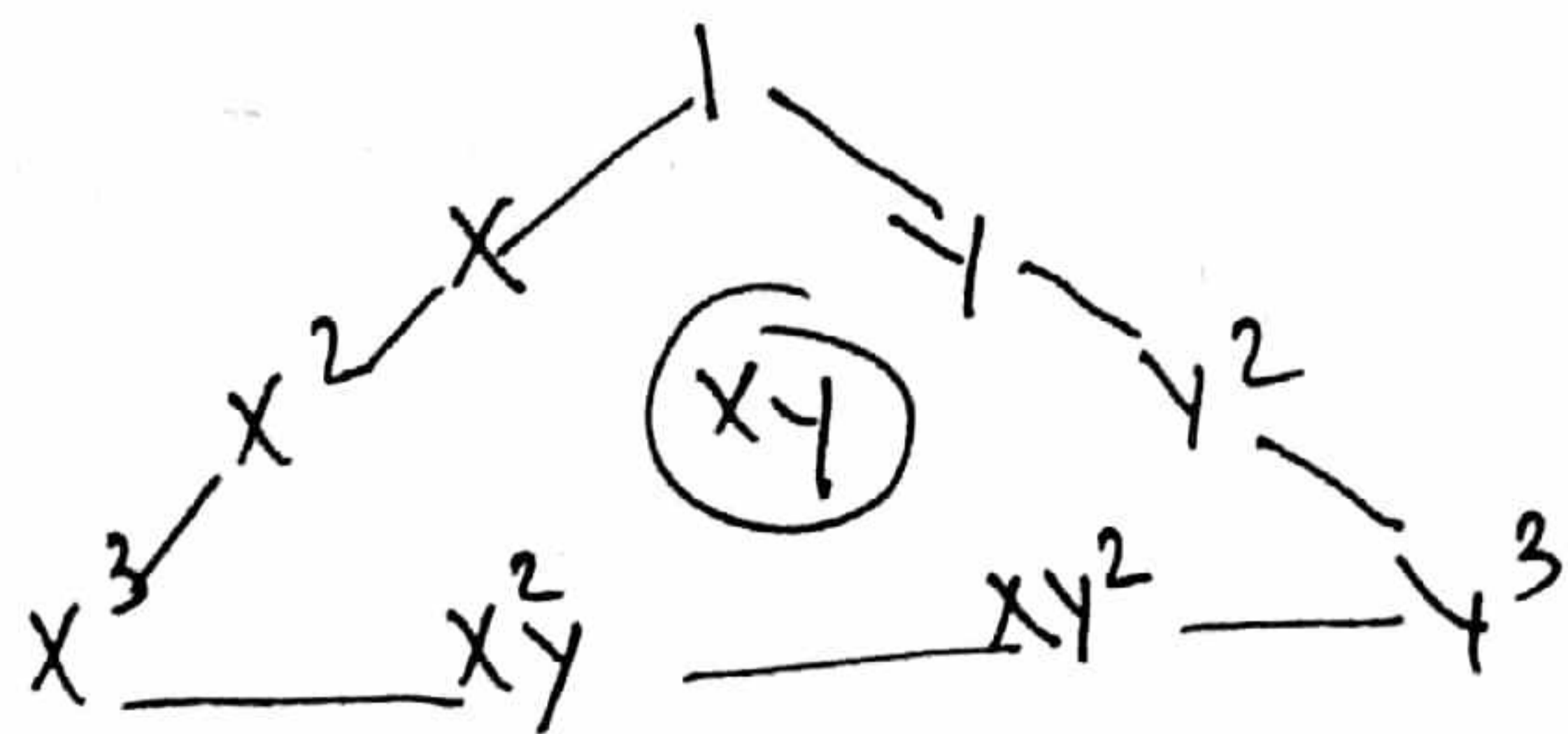
$$1 = C_1 \left( 1 - \frac{1}{2} \right) (1 - 0) \Rightarrow C_1 = 2$$

$$\Rightarrow \psi_1 = L_1 (2L_1 - 1)$$

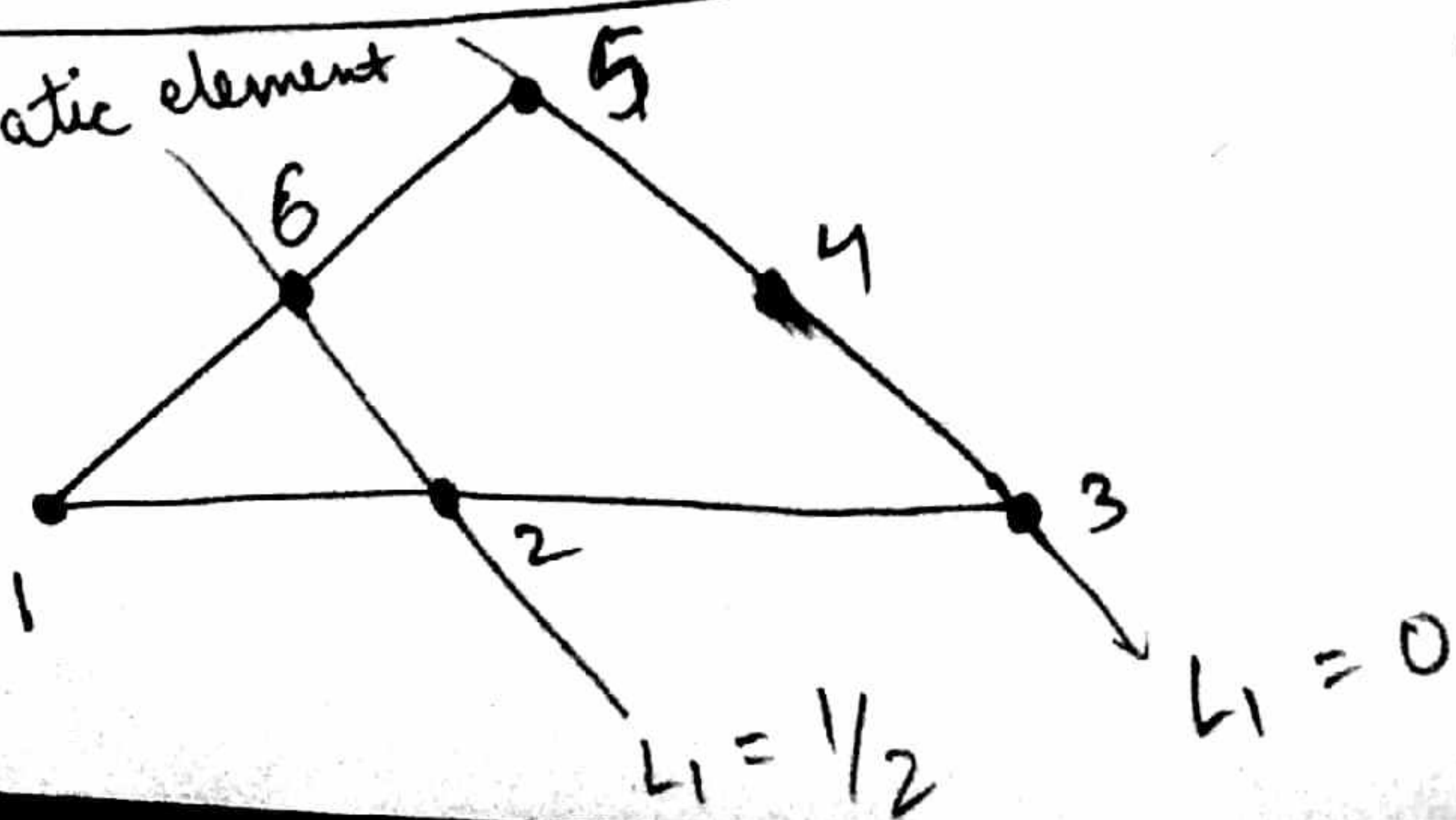
$$L_1 + L_2 + L_3 = 1$$

,  $L_1, L_2, L_3$  are called area coordinates.

Notes:- Cubic triangular element



Quadratic element



$$\psi_1 = C_1 L_1 \left( L_1 - \frac{1}{2} \right)$$

$$1 = C_1 \left( \frac{1}{2} \right) \left( \frac{1}{2} \right)$$

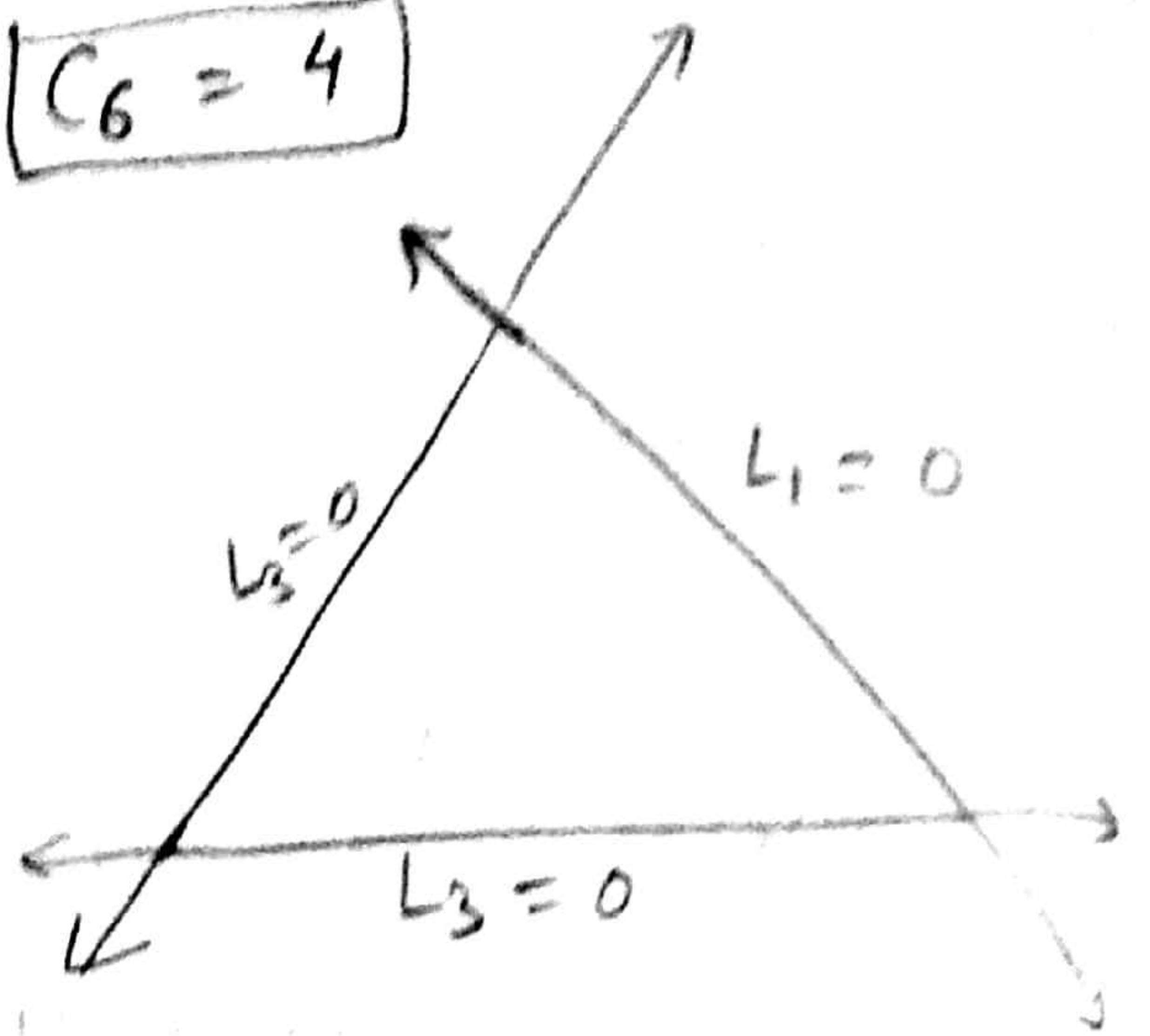
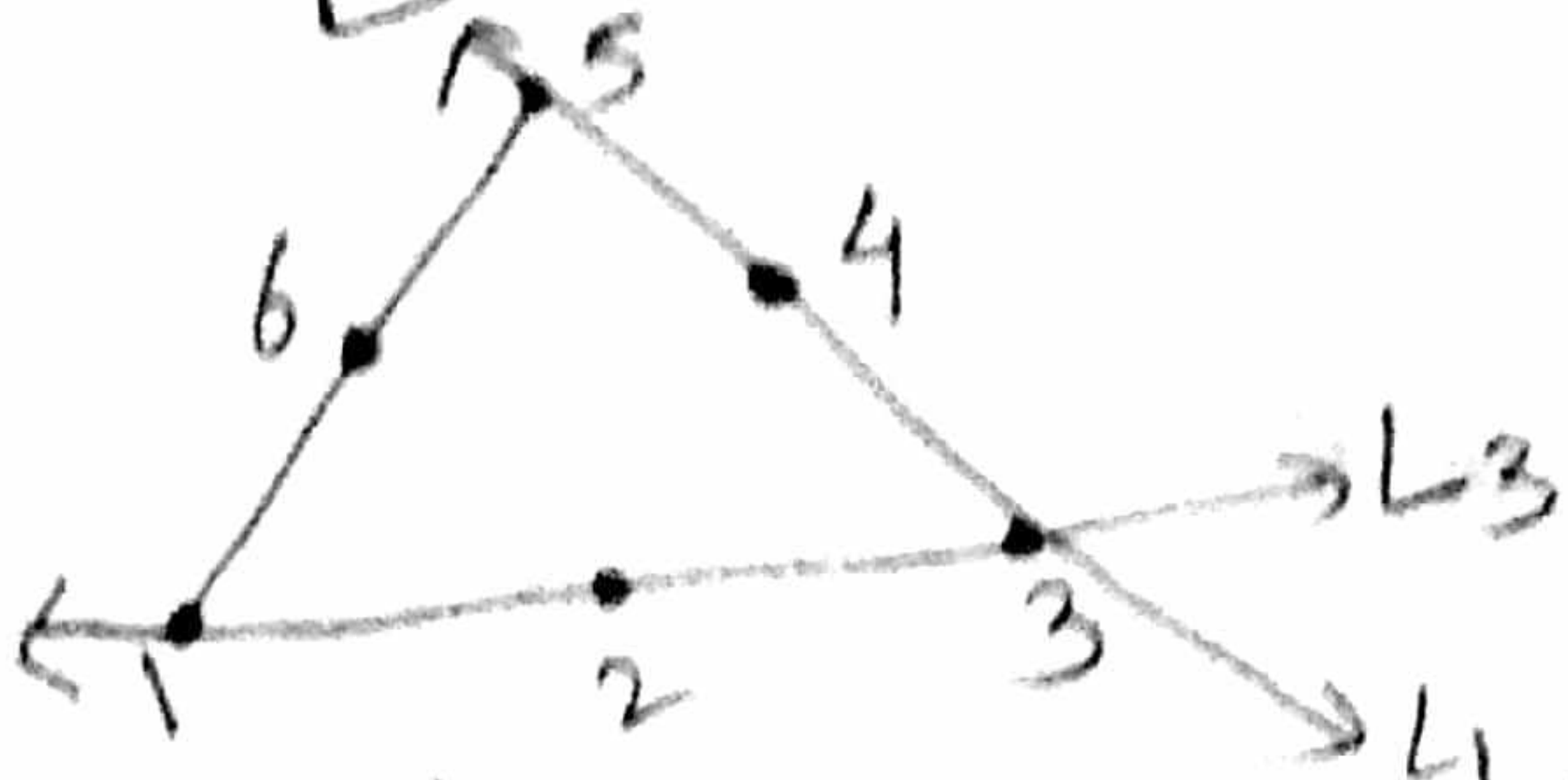
$$\Rightarrow C_1 = 2$$

$$\psi_1 = L_1 (2L_1 - 1)$$

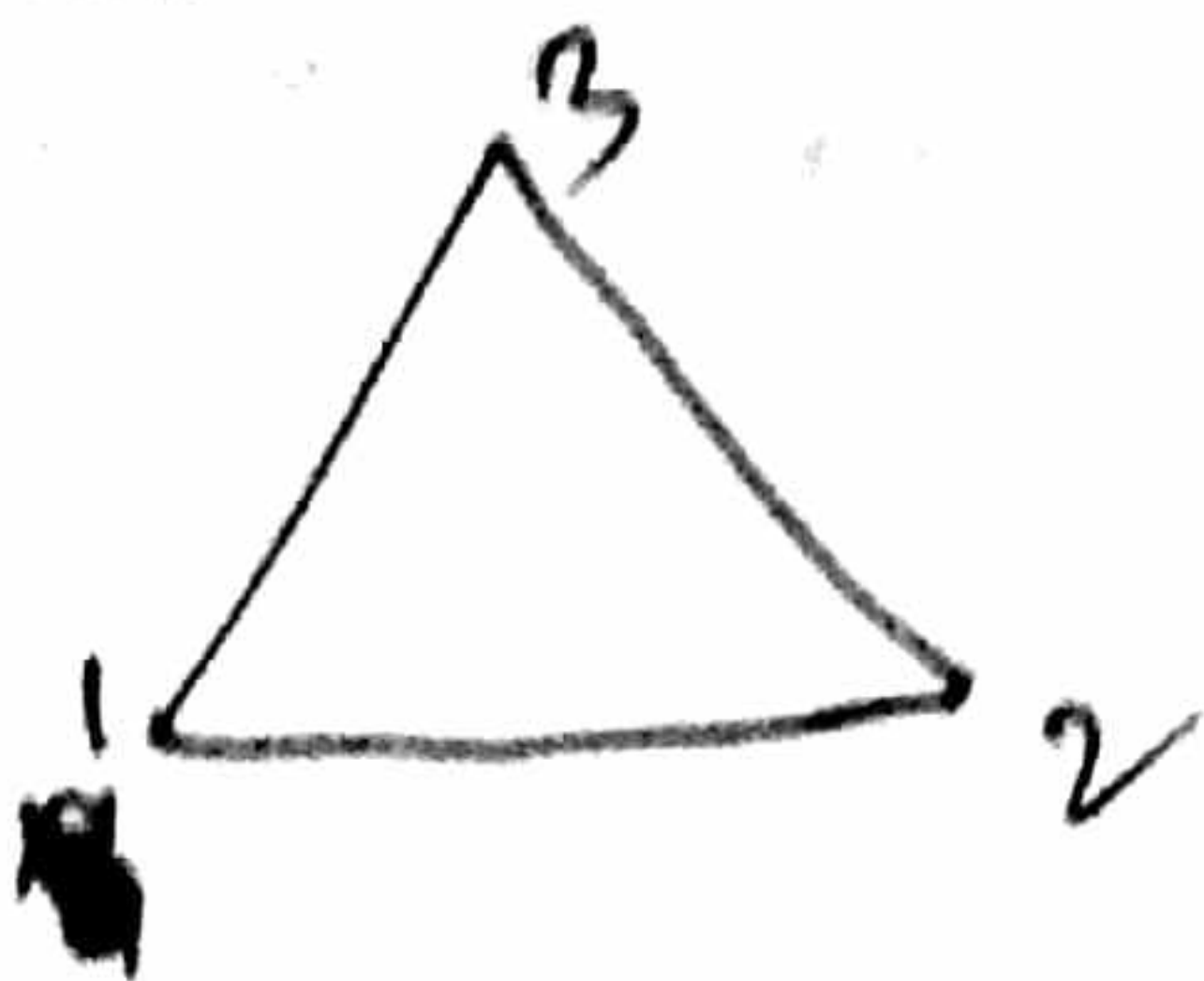
$$\Psi_6 = C_6 * L_1 * L_3$$

$$1 = C_6 * \frac{1}{2} * \frac{1}{2} \Rightarrow C_6 = 4$$

$$\Psi_6 = 4L_1L_3$$



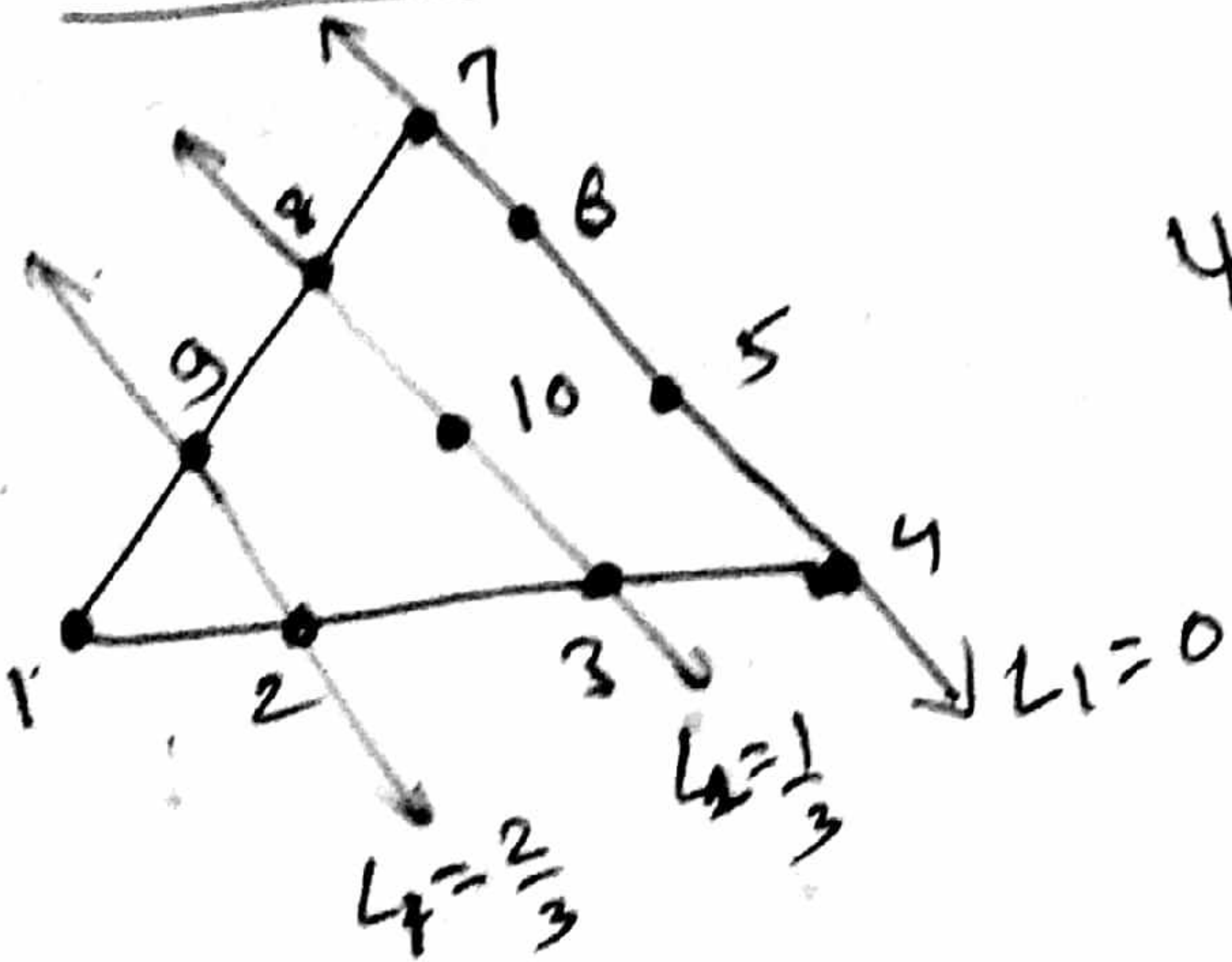
For linear triangular element



$$\Psi_1 = C_1 L_1$$

$$C_1 = 1 \Rightarrow \Psi_1 = L_1$$

For cubic triangular element



$$\Psi_1 = C_1 (L_1 - \frac{2}{3})(L_1 - \frac{1}{3}) L_1$$

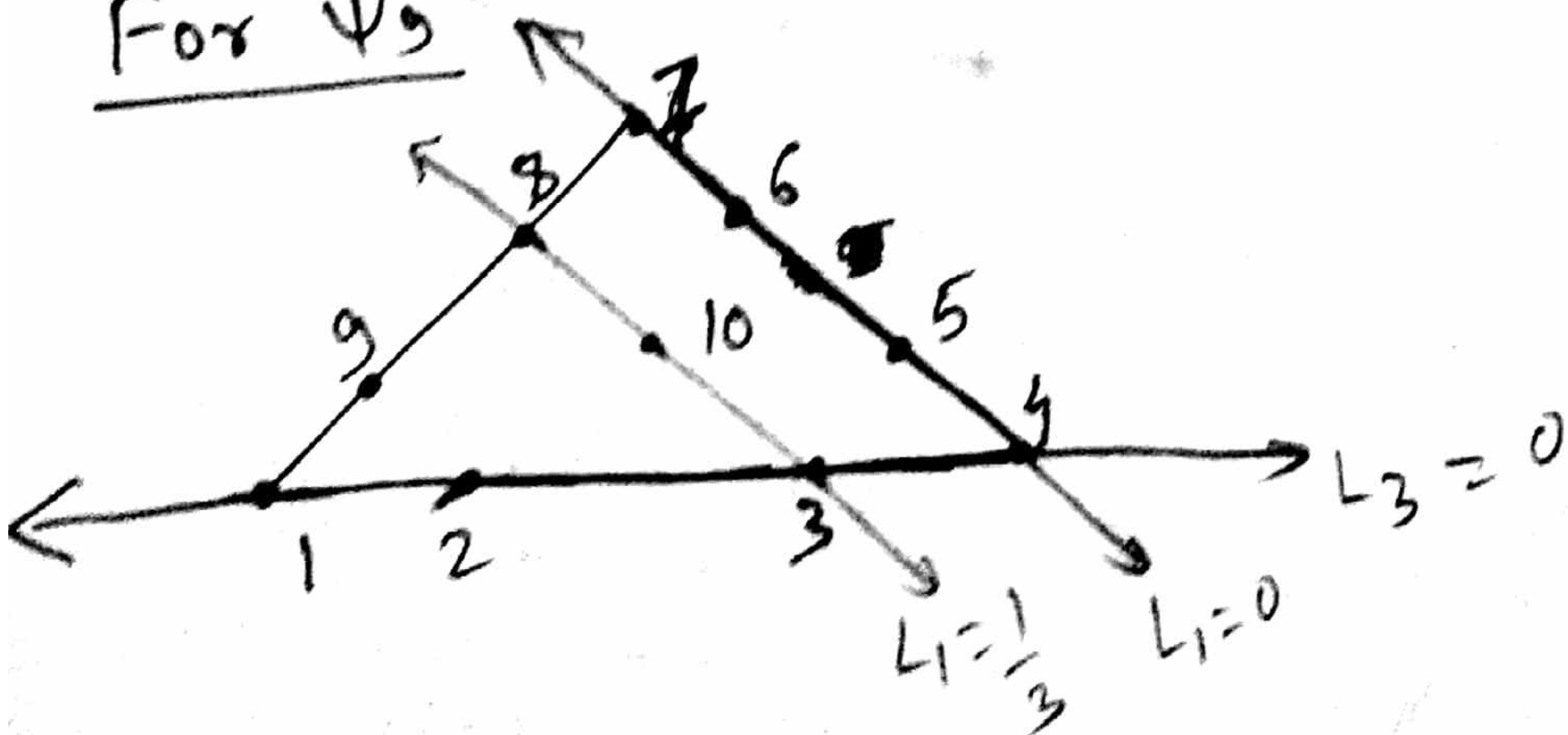
$$\Psi_1(1) = C_1 (\frac{1}{3})(\frac{2}{3})(1)$$

$$C_1 = \frac{9}{2}$$

$$\Psi_1 = \frac{(3L_1 - 2)(3L_1 - 1)L_1}{2}$$

$$\Psi_9 = C_1 L_1 (L_1 - \frac{1}{3}) L_3$$

For  $\Psi_9$



For triangular element

$$\psi_i = f(L_1, L_2, L_3)$$

$$L_1 + L_2 + L_3 = 1 \Rightarrow L_3 = 1 - L_1 - L_2$$

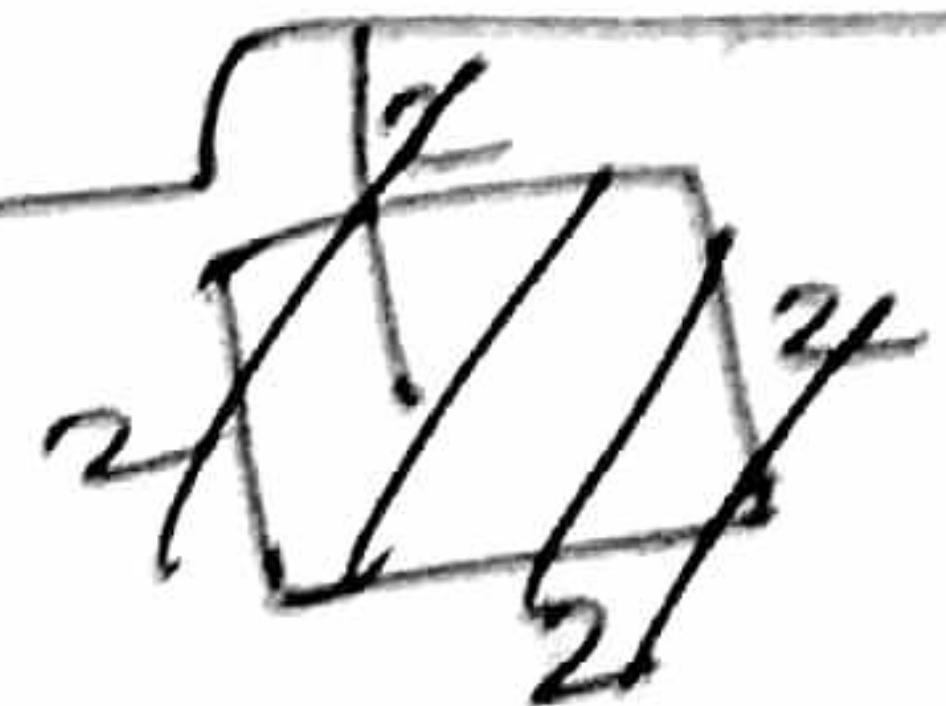
$$\psi_i = f(L_1, L_2, 1 - L_1 - L_2)$$

$$\frac{\partial \psi_i}{\partial L_1} = \frac{\partial \psi_i}{\partial x} \frac{\partial x}{\partial L_1} + \frac{\partial \psi_i}{\partial y} \frac{\partial y}{\partial L_1}$$

$$\frac{\partial \psi_i}{\partial L_2} = \frac{\partial \psi_i}{\partial x} \frac{\partial x}{\partial L_2} + \frac{\partial \psi_i}{\partial y} \frac{\partial y}{\partial L_2}$$

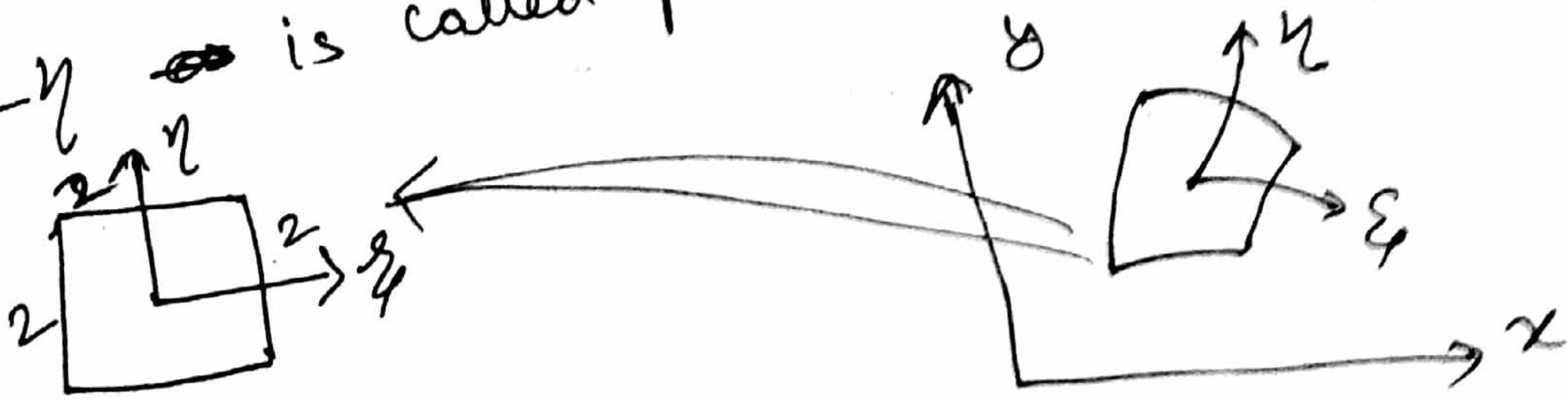
$$\begin{Bmatrix} \frac{\partial \psi_i}{\partial L_1} \\ \frac{\partial \psi_i}{\partial L_2} \\ \frac{\partial \psi_i}{\partial L_3} \end{Bmatrix} = \begin{Bmatrix} \frac{\partial x}{\partial L_1} & \frac{\partial y}{\partial L_1} \\ \frac{\partial x}{\partial L_2} & \frac{\partial y}{\partial L_2} \\ \frac{\partial x}{\partial L_3} & \frac{\partial y}{\partial L_3} \end{Bmatrix} \begin{Bmatrix} \frac{\partial \psi_i}{\partial x} \\ \frac{\partial \psi_i}{\partial y} \end{Bmatrix}$$

$$\begin{Bmatrix} \frac{\partial \psi_i}{\partial L_1} \\ \frac{\partial \psi_i}{\partial L_2} \\ \frac{\partial \psi_i}{\partial L_3} \end{Bmatrix} = \begin{bmatrix} \frac{\partial x}{\partial L_1} & \frac{\partial y}{\partial L_1} \\ \frac{\partial x}{\partial L_2} & \frac{\partial y}{\partial L_2} \\ \frac{\partial x}{\partial L_3} & \frac{\partial y}{\partial L_3} \end{bmatrix} \begin{Bmatrix} \frac{\partial \psi_i}{\partial x} \\ \frac{\partial \psi_i}{\partial y} \end{Bmatrix}$$



Notes

$\xi - \eta$  is called parent element.



Parent element.

Obtuse angle normally should not be more than  $160^\circ$   
 Acute angle should be more than  $20^\circ$ .  
 Aspect ratio in rectangle element should be ~~blw~~  $1:1$  than  
 $1:20$

Time dependent

$$u(x,t) = U(x,t) = \sum_{j=1}^n \psi_j(t) \psi_j(x)$$

↓  
analytical or actual solution

↓  
approximate sol<sup>n</sup>

$$u(x,t) = T(t) X(x)$$

① spatial approximation  
- solving time problem at particular time step.

② Time approximation → How you relate the solution at given time step and next time step.  ~~$\frac{u_j}{\Delta t}$~~   $\frac{u_j}{\Delta t}$

$$\{u\}_s \xrightarrow{\text{small time step}} \{u\}_{s+1}$$

$$b = EI$$

Weak form

$$-\frac{\partial}{\partial x} \left( a \frac{\partial u}{\partial x} \right) + \frac{\partial^2}{\partial x^2} \left( b \frac{\partial^2 u}{\partial x^2} \right) + c_0 u + c_1 \frac{\partial u}{\partial t} + c_2 \frac{\partial^2 u}{\partial t^2} = f(x,t)$$

subjected B.C.s

EBC. ①  $u(x,t)$  or  $-\frac{\partial u}{\partial x}(x,t) + \frac{\partial}{\partial x} \left( b \frac{\partial^2 u}{\partial x^2} \right)$

②  $\frac{\partial u}{\partial x}(x,t)$  or  $b \frac{\partial^2 u}{\partial x^2}$

(iii) initial condition  $c_2 u(x,0)$  and  $c_2 \dot{u}(x,0) + c_1 \dot{u}(x,0)$

Semi-discrete FE formulation

$$\int_{\Omega} w \left[ -\frac{\partial}{\partial x} \left( a \frac{\partial u}{\partial x} \right) + \frac{\partial^2}{\partial x^2} \left( b \frac{\partial^2 u}{\partial x^2} \right) + c_0 u + c_1 \frac{\partial u}{\partial t} + c_2 \frac{\partial^2 u}{\partial t^2} - f \right] dx = 0$$

differentiated twice

$$\Rightarrow \int_{\Omega} \left[ a \frac{\partial w}{\partial x} \frac{\partial u}{\partial x} + b \frac{\partial^2 w}{\partial x^2} \frac{\partial^2 u}{\partial x^2} + c_0 w u + c_1 w \frac{\partial u}{\partial t} + c_2 w \frac{\partial^2 u}{\partial t^2} - w f \right] dx = 0$$

$$-\hat{\Theta}_1 w(x_A) - \hat{\Theta}_3 w(x_B) - \hat{\Theta}_2 \left( -\frac{\partial w}{\partial x} \right) \Big|_{x_A} - \hat{\Theta}_4 \left( \frac{\partial w}{\partial x} \right) \Big|_{x_B} = 0$$

$$\hat{\Theta}_1 = \left[ -a \frac{\partial w}{\partial x} + \frac{\partial}{\partial x} \left( b \frac{\partial^2 w}{\partial x^2} \right) \right]_{x_A}, \quad \hat{\Theta}_2 = \left( b \frac{\partial^2 w}{\partial x^2} \right) \Big|_{x_A}$$



$$\hat{Q}_3 = - \left[ -a \frac{\partial u}{\partial x} + \frac{\partial}{\partial x} b \left( \frac{\partial^2 u}{\partial x^2} \right) \right] x_B, \quad \hat{Q}_4 = - \left( b \frac{\partial^2 u}{\partial x^2} \right) x_B$$

$$u(x,t) = \sum_{j=1}^n u_j(t) \psi_j(x)$$

$$0 = \int_{\Omega} \left\{ a \frac{d\psi_i}{dx} \sum_{j=1}^n \frac{d\psi_j}{dx} u_j + b \frac{\partial^2 \psi_i}{\partial x^2} \left( \sum_{j=1}^n \left( \frac{\partial^2 \psi_j}{\partial x^2} u_j \right) \right) + c_0 \psi_i \left( \sum_{j=1}^n \psi_j u_j \right) + c_1 \psi_i \left( \sum_{j=1}^n \psi_j(x) \dot{u}_j(t) \right) + c_2 \psi_i \left( \sum_{j=1}^n \psi_j(x) \ddot{u}_j(t) \right) - \psi_i f \right\} d\Omega - \hat{Q}_i$$

$$0 = \sum_{j=1}^n \left[ (k_{ij}^1 + k_{ij}^2 + M_{ij}^0) u_j + M_{ij}^1 \frac{du_j}{dt} + M_{ij}^2 \frac{\partial^2 u_j}{\partial t^2} - f_i \right]$$

$$M_{ij}^0 = \int c_0 \psi_i \psi_j dx$$

$$M_{ij}^1 = \int c_1 \psi_i \psi_j dx$$

$$k_{ij}^1 = \int a \frac{d\psi_i}{dx} \frac{d\psi_j}{dx} dx$$

$$M_{ij}^2 = \int c_2 \psi_i \psi_j dx$$

$$k_{ij}^2 = \int b \frac{\partial^2 \psi_i}{\partial x^2} \frac{\partial^2 \psi_j}{\partial x^2} dx$$

$$f_i = \int \psi_i f dx + \hat{Q}_i$$

$$[M][\ddot{u}] + [K][u] = \{f\}$$

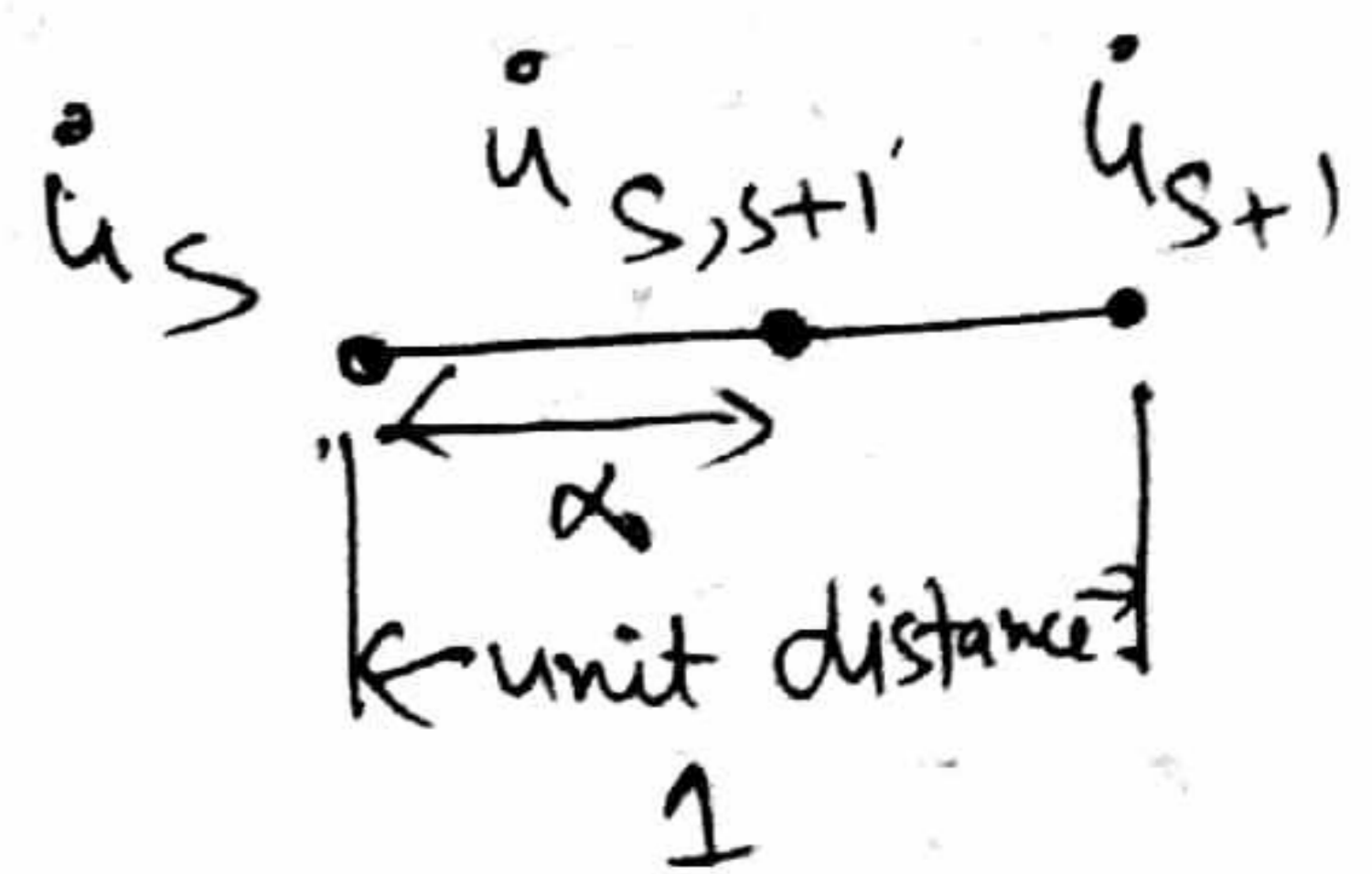
(Parabolic)

↑  
Mass matrix

Time approximation

$$\dot{u} = \frac{u_2 - u_1}{t_2 - t_1} \quad \{ \text{from FDM} \}$$

$$\dot{u}_{s, s+1} = \frac{u_{s+1} - u_s}{\Delta t}$$



$\Delta t \rightarrow$  time step assumed to be equal

$$(y_2 - y_1) = \frac{y_2 - y_1}{x_2 - x_1} * (x_2 - x_1)$$

$$\dot{u}_{s+1} - \dot{u}_{s,s+1} = \frac{(\dot{u}_{s+1} - \dot{u}_s) (1-\alpha)}{1}$$

$$\dot{u}_{s,s+1} = \alpha \dot{u}_{s+1} + (1-\alpha) \dot{u}_s$$

$$\alpha \dot{u}_{s+1} + (1-\alpha) \dot{u}_s = \frac{u_{s+1} - u_s}{\Delta t} \quad \text{--- (1)}$$

$\alpha$  family of approximation

$$[M] \{\dot{u}\}_s + [K]_s \{u\}_s = \{F\}_s \quad \text{--- (2)}$$

$$[M] \{\dot{u}\}_{s+1} + [K]_{s+1} \{u\}_{s+1} = \{F\}_{s+1} \quad \text{--- (3)}$$

Now pre-multiplying eq (1) by  $[M] \Delta t$

$$\Delta t [M] \alpha \dot{u}_{s+1} + \Delta t [M] (1-\alpha) \dot{u}_s = [M] \{u\}_{s+1} - [M] \{u\}_s \quad \text{--- (4)}$$

From (3), (2) & (4)

$$\Delta t \alpha \left\{ \{F\}_{s+1} - [K]_{s+1} \{u\}_{s+1} \right\} + \Delta t (1-\alpha) \left\{ \{F\}_s - [K]_s \{u\}_s \right\}$$

$$= [M] \{u\}_{s+1} - [M] \{u\}_s$$

$$\Rightarrow [M] \{u\}_{s+1} + \Delta t \alpha [K]_{s+1} \{u\}_{s+1} = [M] \{u\}_s - \Delta t (1-\alpha) [K]_s \{u\}_s + \Delta t (1-\alpha) \{F\}_s + \Delta t \alpha \{F\}_{s+1}$$

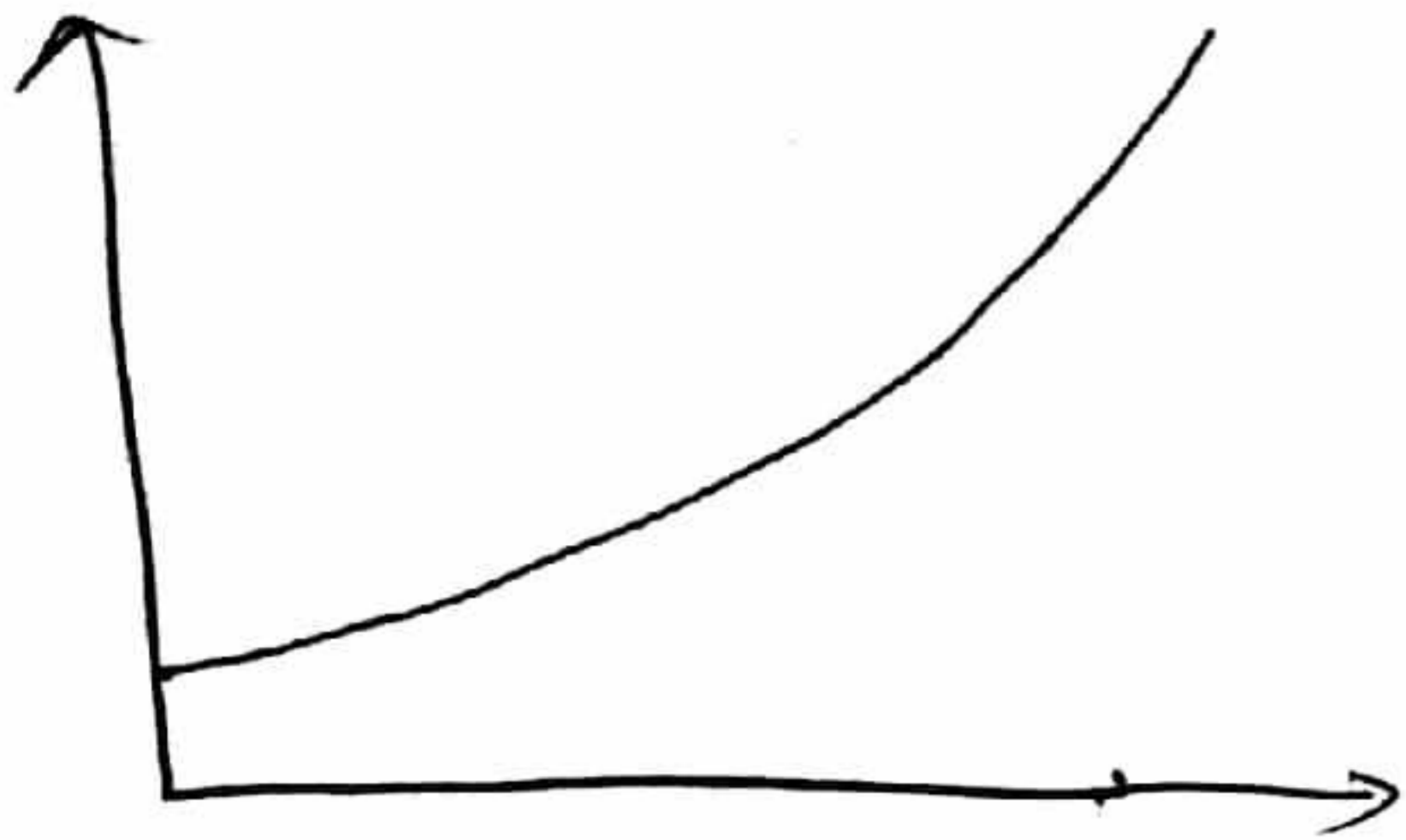
$$\Rightarrow [\hat{K}] \{u\}_{s+1} = [\hat{K}] \{u\}_s + \{\hat{F}\}_{s,s+1}$$

Algorithm

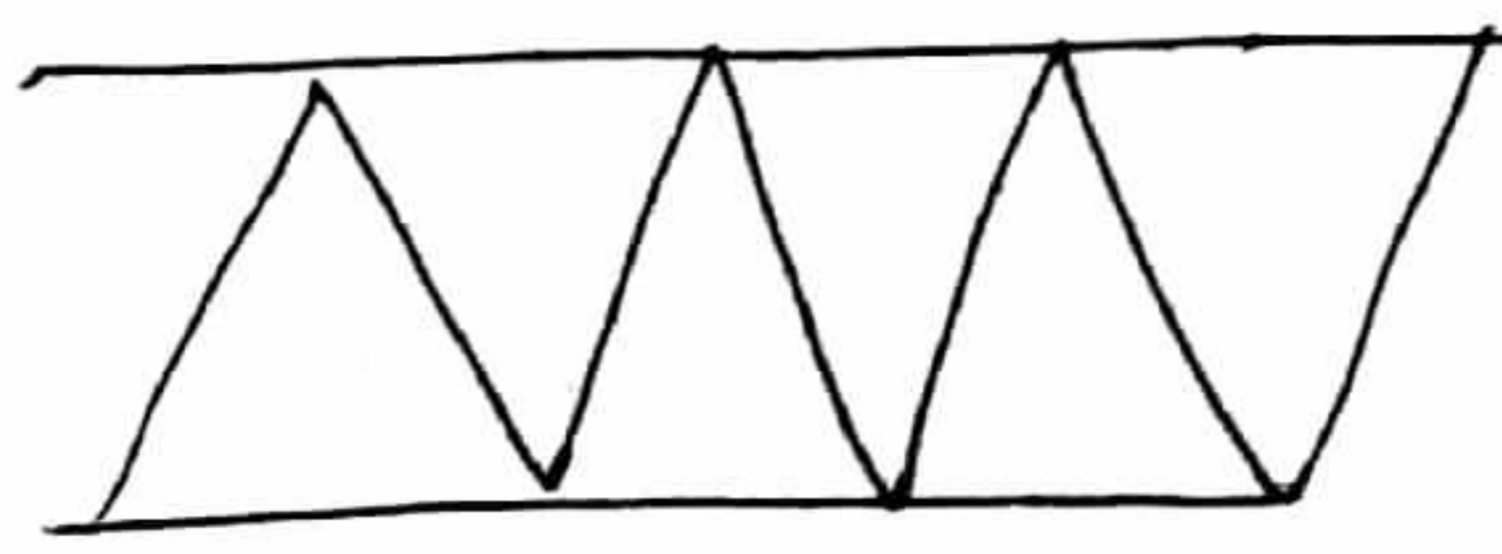
$$([M] + \Delta t \alpha [K]) \{u\}_{s+1} = ([M] - \Delta t (1-\alpha) [K]) \{u\}_s + \Delta t (\alpha \{F\}_{s+1} + (1-\alpha) \{F\}_s)$$

$\{u\}_s = \text{sol}^n$  from initial cond<sup>n</sup>.

unstable sol<sup>n</sup>



Bounded



- $\alpha = 0 \rightarrow$  Forward difference scheme
- $\alpha = 1 \rightarrow$  Backward
- $\alpha = \frac{1}{2} \rightarrow$  Crank - Nicholas
- $\alpha = \frac{2}{3} \rightarrow$  Galerkin.

If  $\alpha = 0$ , the above eq<sup>n</sup> reduces to where  $\{F\}$  includes terms only at time ~~s~~ s.

$$[M] \{u\}_{s+1} = \{F\}$$

$$h \{u\}_{s+1} = [M]^{-1} \{F\}$$

$$\int_0^h c \psi_i \psi_j dx = \frac{ch}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \rightarrow \text{consistent mass matrix}$$

$$\rightarrow \begin{bmatrix} k_{11} & 0 \\ 0 & k_{22} \end{bmatrix} \leftarrow \text{lumped matrix.}$$

Adv: - You need not do matrix inversion.

There are two methods for making lumped matrix

① Row - sum lumping

$$M_{11} = \sum_{i=1}^N \int_0^h c \psi_i \psi_j dx = \int_0^h c \psi_i dx = \frac{ch}{2} \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix}$$

$\Delta t < t_{crit} = \frac{2}{(1-2\alpha)\lambda}$ , for all  $\alpha < 1/2$   
 ↑  
 timestep  $\lambda \rightarrow$  Max<sup>m</sup> eigen value of the problem

If  $\Delta t > t_{crit} \rightarrow$  The problem will be unbounded.  
 When  $\alpha = 0$ , it is called explicit scheme.  
 When  $\alpha \neq 0$ , it is called implicit scheme.  
 $\lambda =$  max<sup>m</sup> eigenvalue of the  $k$  matrix

If  $\alpha = \frac{1}{2}$ ,  $\Delta t = ? \rightarrow$  assess  $\Delta t$ .

$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0 \quad 0 < x < 1$   
 B.C.  $u(0, t) = 0 \quad \frac{\partial u}{\partial t}(1, t) = 0$   
 I.C.  $u(0, 0) = 0 \quad u(1, 0) = 1$   
 $u_x(1, 0) = 1$

Let's consider  $\alpha = \frac{1}{2}$

Mass matrix  $\rightarrow$  consistent mass matrix

$$\int_0^h w \left( \frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} \right) dx = \int_0^h \psi_i \psi_j \{ \dot{u} \} dx - \int_0^h \frac{\partial^2 u}{\partial x^2} w dx$$

on integrating we get  $[M] \{ \dot{u} \} + [K] [u] = [Q]$

$$\int_0^h w \frac{\partial u}{\partial t} dx = \int_0^h \psi_i \frac{d}{dt} (\sum \psi_j u_j) dx = \int_0^h \psi_i \psi_j \{ \dot{u} \} dx$$

we get  $\frac{h}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{Bmatrix} \dot{u}_1 \\ \dot{u}_2 \end{Bmatrix} + \frac{1}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix} = \begin{Bmatrix} Q_1 \\ Q_2 \end{Bmatrix}$

One element model

$$\begin{bmatrix} \frac{h}{3} + \alpha \frac{\Delta t}{h} & \frac{h}{6} - \frac{\alpha \Delta t}{h} \\ \frac{h}{6} - \alpha \Delta t & \frac{h}{3} + \frac{\alpha \Delta t}{h} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}_{s+1} = ([M] + \alpha \Delta t [K]) \{u\}_{s+1}$$

$$([M] - (1-\alpha) \Delta t [K]) = \begin{bmatrix} \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} & \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} \\ \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} & \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} \end{bmatrix}$$

If  $\{F\}$  does not change with time then  $\alpha \{F\}_{s+1} + \{F\}_s (1-\alpha) = \{F\}_s \Rightarrow \{F\}_s = \{0\}?$

Given  $\frac{\partial u}{\partial t}(1,0) = 0 \Rightarrow Q_2 = 0$

~~u~~

$$\begin{bmatrix} \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} & \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} \\ \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} & \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}_s + \begin{bmatrix} Q_1 \\ Q_2 \end{bmatrix}$$

$$\begin{bmatrix} \frac{h}{3} + \frac{\alpha \Delta t}{h} & \frac{h}{6} - \frac{\alpha \Delta t}{h} \\ \frac{h}{6} - \frac{\alpha \Delta t}{h} & \frac{h}{3} + \frac{\alpha \Delta t}{h} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}_{s+1} = \begin{bmatrix} \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} & \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} \\ \frac{h}{6} + \frac{(1-\alpha) \Delta t}{h} & \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \end{bmatrix}_s + \begin{bmatrix} Q_1 \\ Q_2 \end{bmatrix}$$

$$\left( \frac{h}{3} + \frac{\alpha \Delta t}{h} \right) (u_2)_{s+1} = \frac{h}{3} - (1-\alpha) \frac{\Delta t}{h} (u_2)_s$$

if  $\Delta t = 0.05$ ,  $\alpha = 0.5$

Do using One quadratic element,

Eigen value problems

$$A(u) = \lambda B(u)$$

$$-\frac{d^2 u}{dx^2} = \lambda u$$

A & B are differential operators

DE  $\rho c A \frac{\partial u}{\partial t} - \frac{\partial}{\partial x} \left( k A \frac{\partial u}{\partial x} \right) = q(x,t)$   $\rightarrow$  sol<sup>n</sup>

$\rightarrow$  homogeneous form  
(from variable separation method)

$$u(x,t) = U(x) e^{-\lambda t}$$

Parabolic eqn type.

$$-\frac{d}{dx} \left( k A \frac{dU(x)}{dx} \right) - \lambda B C A U(x) = 0$$

Governing eq<sup>n</sup> for FEM sol<sup>n</sup>.

On substituting the sol<sup>n</sup> in DE.

$\uparrow$  Eigen value problem after substituting the sol<sup>n</sup>

$$\int_0^h w \left( -\frac{d}{dx} \left( k A \frac{dU(x)}{dx} \right) - \lambda B C A U(x) \right) dx = 0$$

weighted integral form

8.2  $\rho A \frac{\partial^2 u}{\partial t^2} - \frac{\partial}{\partial x} \left( EA \frac{\partial u}{\partial x} \right) = q(x,t)$  ← hyperbolic eq<sup>n</sup>

Sol<sup>n</sup> →  $u(x,t) = U(x) e^{-i\omega t}$   $i = \sqrt{-1}$

Substitute the sol<sup>n</sup> in Differential eq<sup>n</sup>.

$$\left[ \rho A U(x) (+i^2 \omega^2) - \frac{d}{dx} EA \frac{dU}{dx} \right] e^{-i\omega t} = 0$$

$$\boxed{-\rho A U(x) \omega^2 - \frac{d}{dx} EA \frac{dU(x)}{dx} = 0}$$

$\omega =$  natural frequency.

domain  
 $0 \leq x \leq 1$

$$-\frac{d^2 u}{dx^2} + \frac{\partial u}{\partial t} = 0$$

$$u(0,t) = 0 \quad ; \quad \left( \frac{\partial u}{\partial x} + u \right) \Big|_{x=1} = 0$$

$$u(x,t) = U(x) e^{-\lambda t}$$

$$u(0,t) = U(0) e^{-\lambda t} = 0 \Rightarrow \boxed{U(0) = 0}$$

$$\left( \frac{dU(x)}{dx} + U(x) \right) \Big|_{x=1} = 0 \Rightarrow \frac{dU(x)}{dx} \Big|_{x=1} + U(1) = 0$$

$$Q_2 = -Q_3 = Q_3$$

Substitute  $u(x,t) = U(x)e^{-\lambda t}$  in DE

$$-\frac{d^2 U}{dx^2} - \lambda U = 0$$

$$\int_0^h w \left( -\frac{d^2 U}{dx^2} - \lambda U \right) dx = 0$$

$$\int_0^h \frac{\partial w}{\partial x} \frac{\partial U}{\partial x} dx - \int_0^h w \lambda U dx - \int_0^h w \frac{dU}{dx} dx = 0$$

$$\int_0^h \left( \frac{d\psi_i}{dx} \frac{d}{dx} \sum_{j=1}^n \psi_j U_j \right) dx - \lambda \int_0^h \psi_i \left( \sum_{j=1}^n U_j \psi_j \right) dx = Q_i$$

$$[k_{ij}] \{U_j\} - \frac{\lambda h}{6} [M_{ij}] \{U_j\} = \{B_i\}$$

Considering linear elements (minimum elements are 2) 1 2 3

$$\frac{1}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} - \frac{\lambda h}{6} \begin{bmatrix} 2 & 1 \\ 1 & 2 \end{bmatrix} \begin{bmatrix} U_1 \\ U_2 \end{bmatrix} = \begin{bmatrix} B_1 \\ B_2 \end{bmatrix}$$

$$\left\{ \frac{1}{h} \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix} - \frac{\lambda h}{6} \begin{bmatrix} 2 & 1 & 0 \\ 1 & 4 & 1 \\ 0 & 1 & 2 \end{bmatrix} \right\} \begin{bmatrix} U_1 \\ U_2 \\ U_3 \end{bmatrix} = \begin{bmatrix} B_1 \\ B_2 + B_1 \\ B_2 \end{bmatrix}$$

Condensed matrix directly for  $U_2$  and  $U_3$

$$\begin{bmatrix} \frac{2}{h} - \frac{4\lambda h}{6} & -\frac{1}{h} - \frac{\lambda h}{6} \\ -\frac{1}{h} - \frac{\lambda h}{6} & \frac{1}{h} - \frac{2\lambda h}{6} \end{bmatrix} \begin{bmatrix} U_2 \\ U_3 \end{bmatrix} = \begin{bmatrix} 0 \\ -U_3 \end{bmatrix}$$

eqn of  $\lambda$  is called determinant of the ~~matrix~~ characteristic equation. matrix below gives a 2nd order

$$\begin{bmatrix} \frac{2}{h} - \frac{4\lambda h}{6} & -\frac{1}{h} - \frac{\lambda h}{6} \\ -\frac{1}{h} - \frac{\lambda h}{6} & \frac{1}{h} + 1 - \frac{2\lambda h}{6} \end{bmatrix} \begin{bmatrix} U_2 \\ U_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} (a_{11} - b_{11}\lambda) & (a_{12} - b_{12}\lambda) \\ (a_{21} - b_{21}\lambda) & (a_{22} - b_{22}\lambda) \end{bmatrix} \begin{bmatrix} U_2 \\ U_3 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

For  $\lambda = \lambda_1$  (1st eigen value)

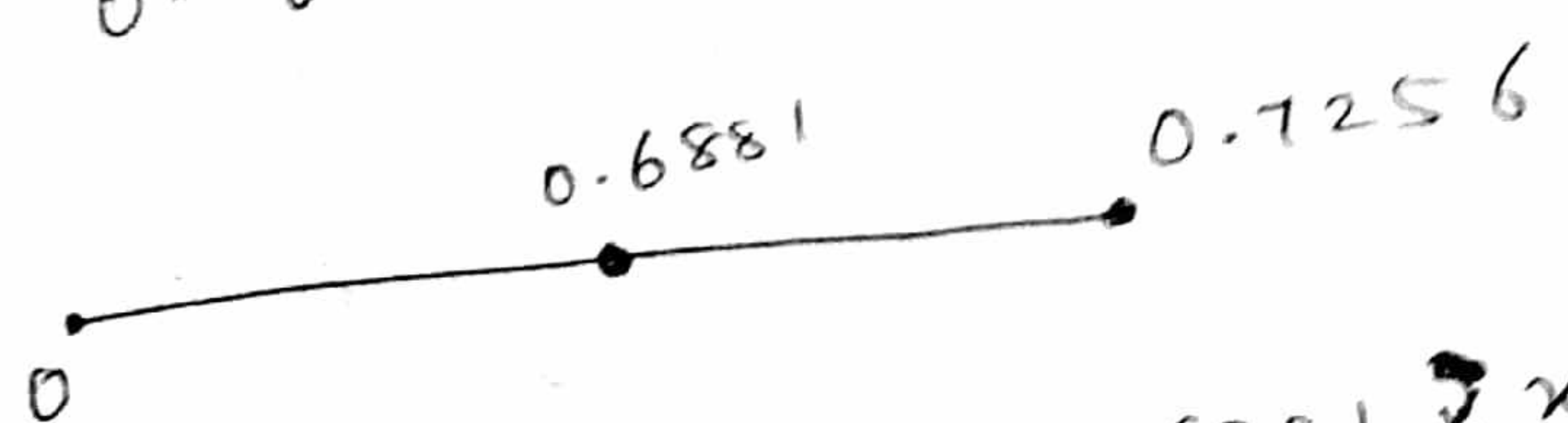
$$(a_{11} - b_{11}\lambda_1) U_2^{(1)} + (a_{12} - b_{12}\lambda_1) U_3^{(1)} = 0$$



$$2.5033 U_2 - 2.3742 U_3 = 0$$

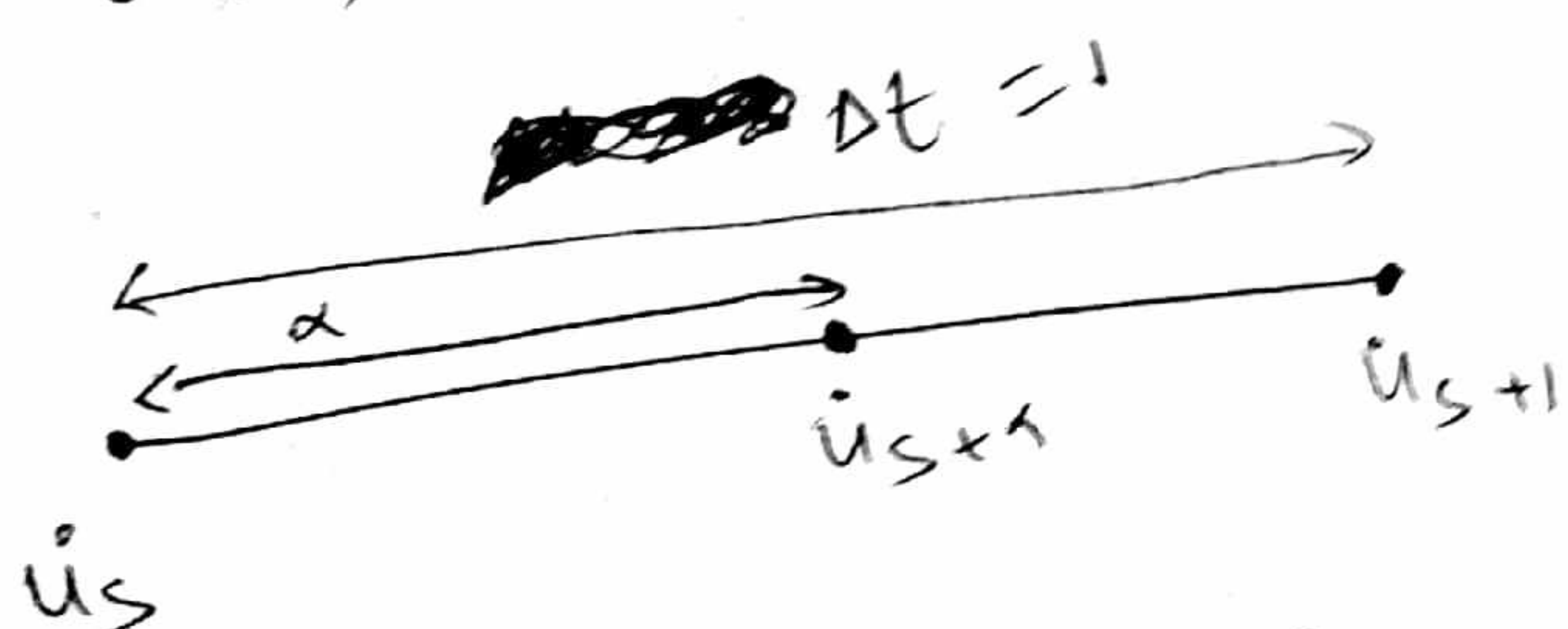
$$1.0 U_2 - 1.0544 U_3 = 0$$

$$0.6881 U_2 - 0.7256 U_3 = 0$$



$$U'(x) = 0 \left(1 - \frac{x}{h}\right) + 0.6881 \frac{x}{h}$$

$$U'(x) = 0.6881 \left(1 - \frac{x}{h}\right) + 0.7256 \frac{x}{h}$$



$$\frac{du}{dt} \approx \frac{u_{s+1} - u_s}{\Delta t} = \dot{u}_{s+1} - \dot{u}_s$$

$$\frac{\dot{u}_{s+\alpha} - \dot{u}_s}{(\alpha - 0)} = \frac{\dot{u}_{s+1} - \dot{u}_s}{(1 - 0)}$$

$$\dot{u}_{s+\alpha} = \dot{u}_s (1 - \alpha) + \dot{u}_{s+1} \alpha$$

$$\frac{u_{s+\alpha} - u_s}{\Delta t} =$$

$$u_s (1 - \alpha) + u_{s+1} \alpha$$

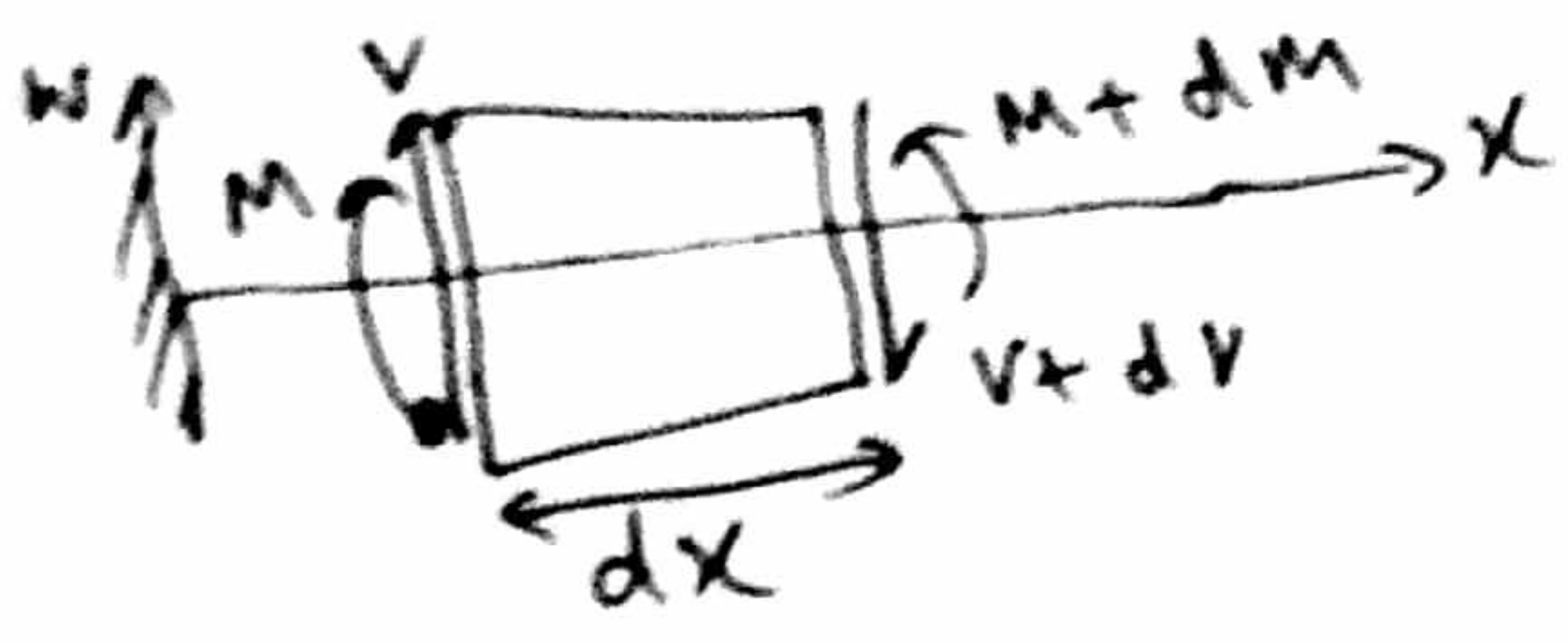


$$b = EI$$

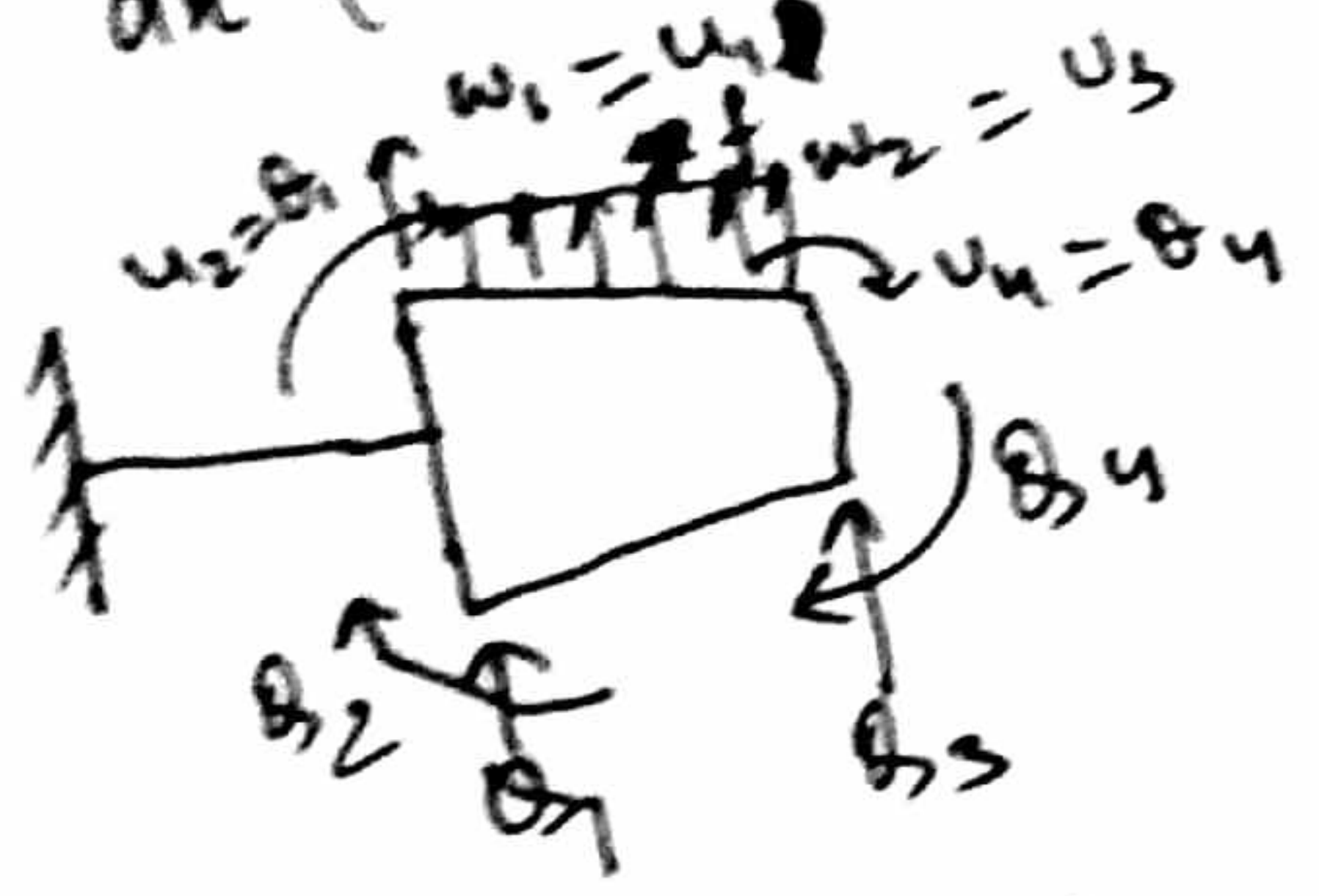
$$M = b \frac{d^2 w}{dx^2}$$

$$V = b \frac{dM}{dx}$$

$$\frac{dV}{dx} = f(x)$$



$$\frac{d^2}{dx^2} \left( b \frac{d^2 w}{dx^2} \right) = f(x)$$



$$\int_{x_e}^{x_{e+1}} v \left[ \frac{d^2}{dx^2} b \frac{d^2 w}{dx^2} - f \right] dx = 0$$

$v \rightarrow$  weight function

$$\int_{x_e}^{x_{e+1}} \left[ -\frac{dv}{dx} \frac{d}{dx} b \frac{d^2 w}{dx^2} - v f \right] dx + \left[ v \frac{d}{dx} b \frac{d^2 w}{dx^2} \right]_{x_e}^{x_{e+1}} = 0$$

$\theta_1, \theta_4$  are the slopes  
 $\theta_2, \theta_3$  are forces  
 $\theta_1, \theta_3$  are forces  
 $\theta_2, \theta_4$  are the moments

2nd  $\rightarrow$   $\int_{x_e}^{x_{e+1}} \left[ b \frac{d^2 v}{dx^2} \frac{d^2 w}{dx^2} - v f \right] dx + \left[ v \frac{d}{dx} b \frac{d^2 w}{dx^2} - \frac{dv}{dx} b \frac{d^2 w}{dx^2} \right]_{x_e}^{x_{e+1}} = 0$   
 Coefficient of functions of weight function gives natural boundary condition

where is field variable

NBC  $\rightarrow \frac{d}{dx} b \frac{d^2 w}{dx^2}$   
 Shear force  $\rightarrow \frac{d}{dx} b \frac{d^2 w}{dx^2}$   
 Moment  $\rightarrow b \frac{d^2 w}{dx^2}$

The weight  $f^n$  written in terms of field variable gives Essential boundary conditions

EBc  $v$   
 $\frac{dv}{dx}$

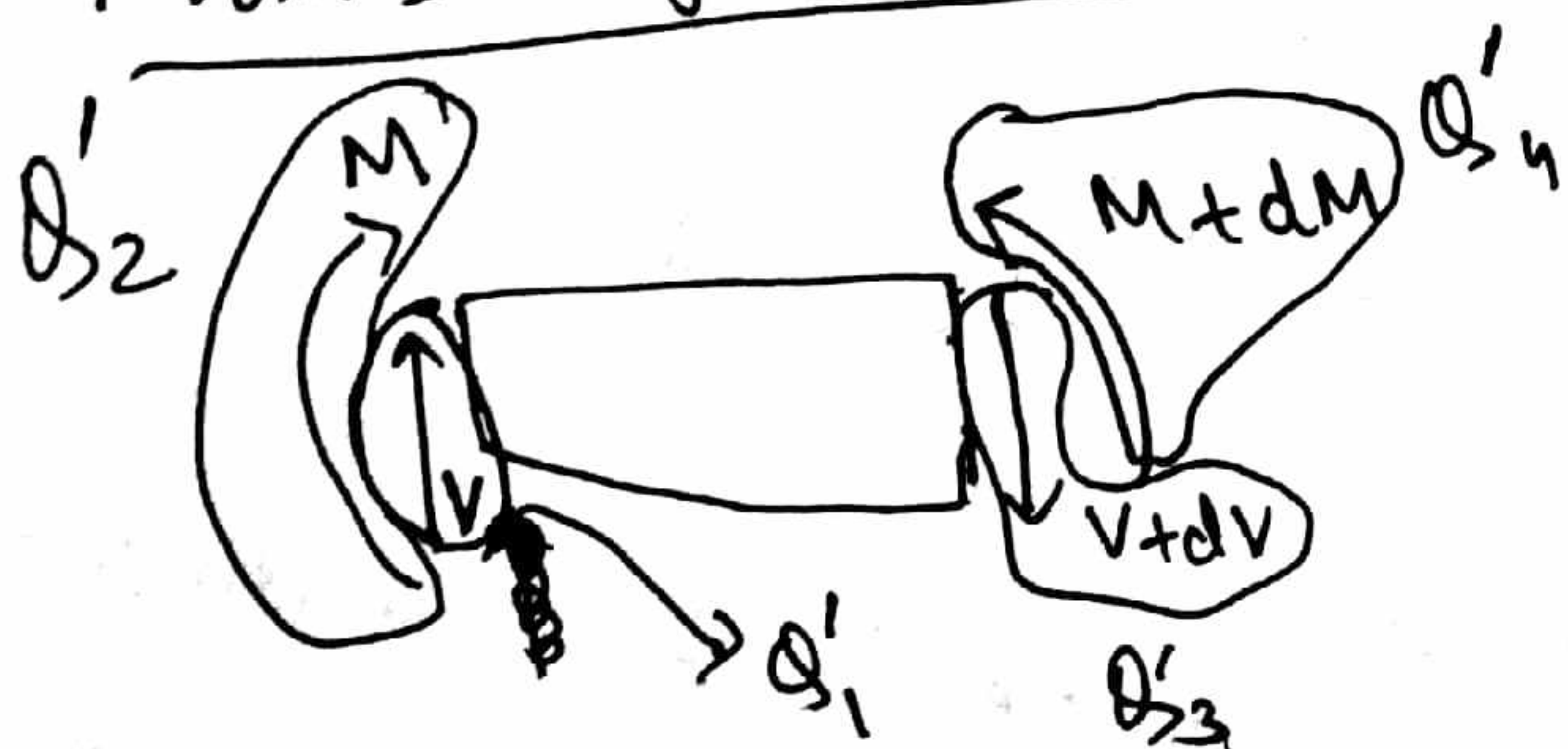
$$\delta_1 = \left[ \frac{d}{dx} \left( b \frac{d^2 w}{dx^2} \right) \right]_{x^e}$$

$$\delta_3 = - \frac{d}{dx} \left( b \frac{d^2 w}{dx^2} \right)_{x^{e+1}}$$

$$\delta_2 = \left( b \frac{d^2 w}{dx^2} \right)_{x^e}$$

$$\delta_4 = - \left( b \frac{d^2 w}{dx^2} \right)_{x^{e+1}}$$

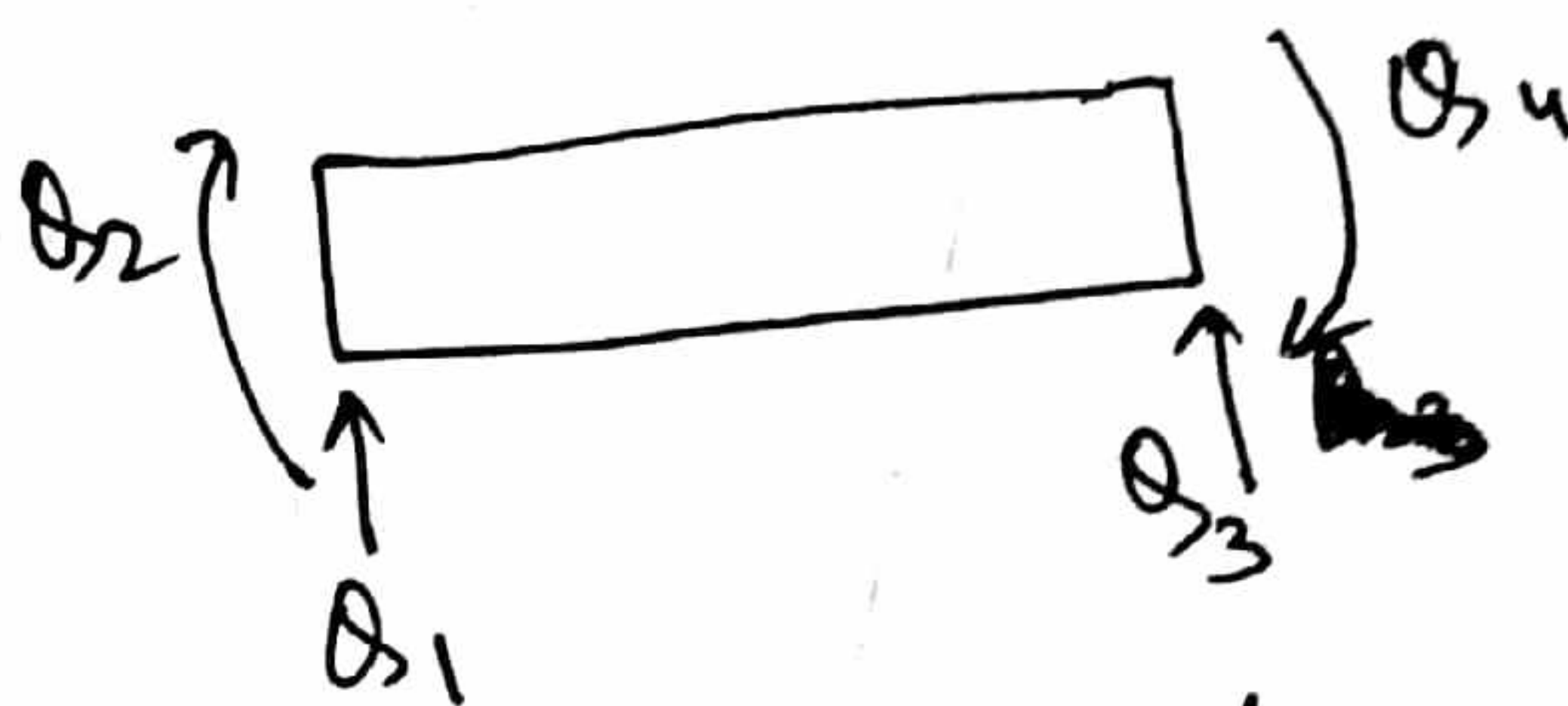
From strengths of materials



$$\delta_3' = \frac{d}{dx} \left( b \frac{d^2 w}{dx^2} \right)_{x^{e+1}}$$

$$\delta_4' = - \left( b \frac{d^2 w}{dx^2} \right)_{x^{e+1}}$$

For FEM



so,  $\delta_2 = -\delta_3$   
 $\delta_4 = -\delta_4$

$$0 = \int_{x^e}^{x^{e+1}} \left( b \frac{d^2 v}{dx^2} \frac{d^2 w}{dx^2} - v f \right) dx - v(x^e) \delta_1 - \left( -\frac{dv}{dx} \right) \Big|_{x^e} \delta_2 - v(x^{e+1}) \delta_3 - \left( -\frac{dv}{dx} \right) \Big|_{x^{e+1}} \delta_4$$

Weak form

Total no. of nodal variables at each node = 2 (two degrees of freedom per node)

Total no. of nodal variables = 4 (for one linear element)

$u_1, \theta_1$        $u_2, \theta_2$

$$w = c_1 + c_2 x + c_3 x^2 + c_4 x^3$$

$$\theta = - \frac{dw}{dx}$$

$$u = w$$

$$U_1 = W(x^e), \quad U_2 = -\left.\frac{dW}{dx}\right|_{x=x^e}, \quad U_3 = W(x^{e+1}), \quad U_4 = \left.\left(-\frac{dW}{dx}\right)\right|_{x^{e+1}}$$

Generalised displacements  $u, \theta$  represented by  $U$ .

$$U_1 = C_1 + C_2 \cdot x^e + C_3 (x^e)^2 + C_4 (x^e)^3$$

$$U_2 = 0 - C_2 - 2C_3 x^e - 3C_4 (x^e)^2$$

$$U_3 = C_1 + C_2 x^{e+1} + C_3 (x^{e+1})^2 + C_4 (x^{e+1})^3$$

$$U_4 = 0 - C_2 - 2C_3 x^{e+1} - 3C_4 (x^{e+1})^2$$

$$\{U\} = \begin{bmatrix} 1 & x^e & (x^e)^2 & (x^e)^3 \\ 0 & -1 & -2x^e & -3(x^e)^2 \\ 1 & x^{e+1} & (x^{e+1})^2 & (x^{e+1})^3 \\ 0 & -1 & -2x^{e+1} & -3(x^{e+1})^2 \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \\ C_3 \\ C_4 \end{bmatrix}$$

$$W(x) = U_1 \phi_1 + U_2 \phi_2 + U_3 \phi_3 + U_4 \phi_4$$

$$\phi_1 = 1 - 3\left(\frac{\bar{x}}{h}\right)^2 + 2\left(\frac{\bar{x}}{h}\right)^3$$

$$\phi_2 = -\bar{x} \left(1 - \frac{\bar{x}}{h}\right)^2$$

$$\phi_3 = 3\left(\frac{\bar{x}}{h}\right)^2 - 2\left(\frac{\bar{x}}{h}\right)^3$$

$$\phi_4 = -\bar{x} \left[\left(\frac{\bar{x}}{h}\right)^2 - \left(\frac{\bar{x}}{h}\right)\right]$$

$$\frac{d\phi_1}{dx} = -\frac{6}{h} \frac{\bar{x}}{h} \left(1 - \frac{\bar{x}}{h}\right)$$

$$\frac{d\phi_2}{dx} = -\left[1 + 3\left(\frac{\bar{x}}{h}\right)^2 - 4\left(\frac{\bar{x}}{h}\right)\right]$$

$$\frac{d\phi_3}{dx} = -\frac{d\phi_1}{dx}$$

$$\frac{d\phi_4}{dx} = -\frac{\bar{x}}{h} \left[\frac{3\bar{x}}{h} - 2\right]$$

Written in Local coordinates.  
No need to memorize,  
will be given in the exams

hermite cubic interpolation  
Functions

$$U_1 = w(x^e), \quad U_2 = -\frac{dw}{dx}\bigg|_{x=x^e}, \quad U_3 = w(x^{e+1}), \quad U_4 = \left(\frac{dw}{dx}\right)\bigg|_{x=x^{e+1}}$$

Generalised displacements  $u, \theta$  represented by  $U$ .

$$U_1 = C_1 + C_2 \cdot x^e + C_3 (x^e)^2 + C_4 (x^e)^3$$

$$U_2 = 0 - C_2 - 2C_3 x^e - 3C_4 (x^e)^2$$

$$U_3 = C_1 + C_2 x^{e+1} + C_3 (x^{e+1})^2 + C_4 (x^{e+1})^3$$

$$U_4 = 0 - C_2 - 2C_3 x^{e+1} - 3C_4 (x^{e+1})^2$$

$$\{U\} = \begin{bmatrix} 1 & x^e & (x^e)^2 & (x^e)^3 \\ 0 & -1 & -2x^e & -3(x^e)^2 \\ 1 & x^{e+1} & (x^{e+1})^2 & (x^{e+1})^3 \\ 0 & -1 & -2x^{e+1} & -3x^{e+1} \end{bmatrix} \begin{bmatrix} C_1 \\ C_2 \\ C_3 \\ C_4 \end{bmatrix}$$

$$w(x) = U_1 \phi_1 + U_2 \phi_2 + U_3 \phi_3 + U_4 \phi_4$$

$$\phi_1 = 1 - 3\left(\frac{\bar{x}}{h}\right)^2 + 2\left(\frac{\bar{x}}{h}\right)^3$$

$$\phi_2 = -\bar{x} \left(1 - \frac{\bar{x}}{h}\right)^2$$

$$\phi_3 = 3\left(\frac{\bar{x}}{h}\right)^2 - 2\left(\frac{\bar{x}}{h}\right)^3$$

$$\phi_4 = -\bar{x} \left[\left(\frac{\bar{x}}{h}\right)^2 - \left(\frac{\bar{x}}{h}\right)\right]$$

$$\frac{d\phi_1}{dx} = -\frac{6\bar{x}}{hh} \left(1 - \frac{\bar{x}}{h}\right)$$

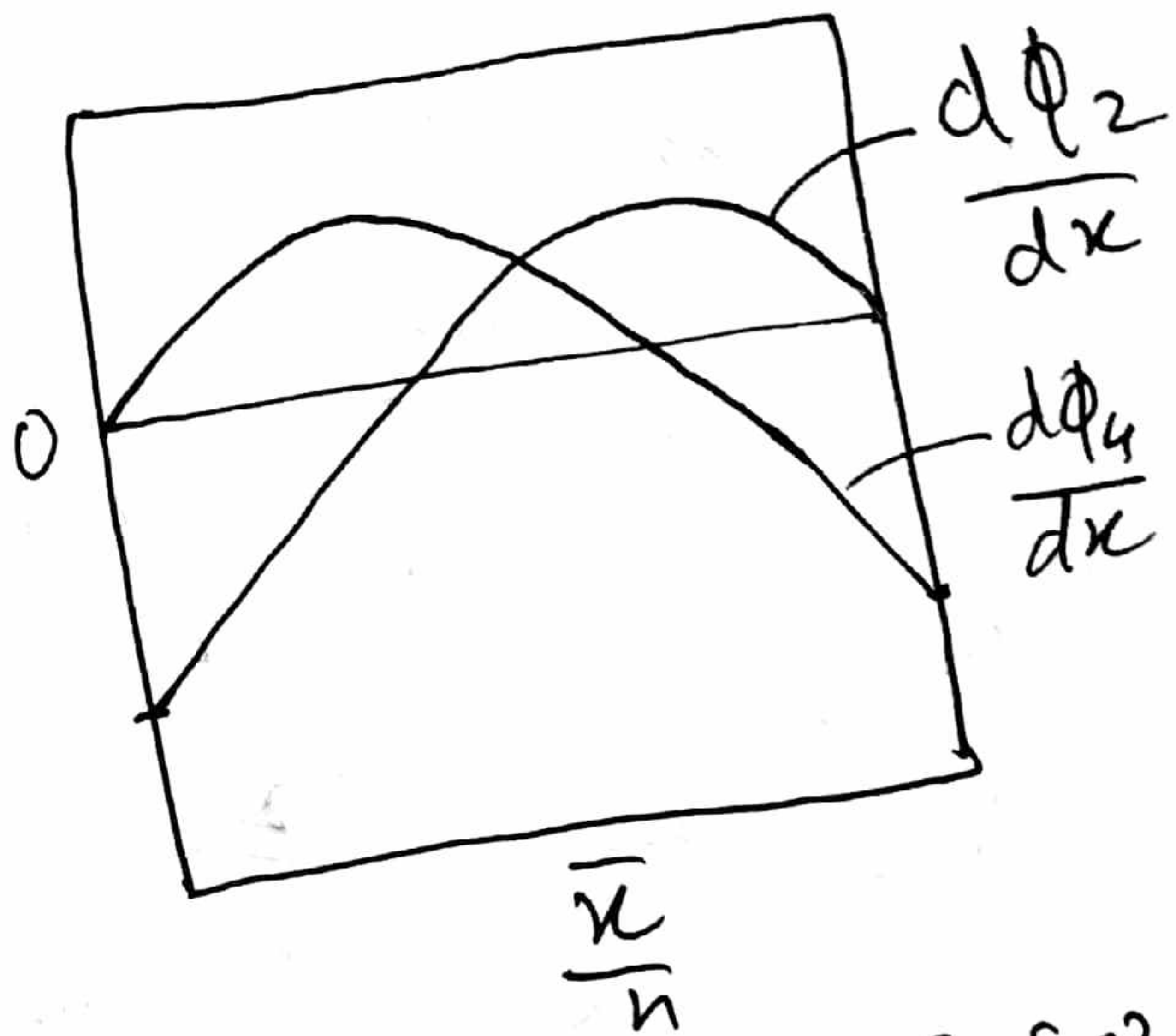
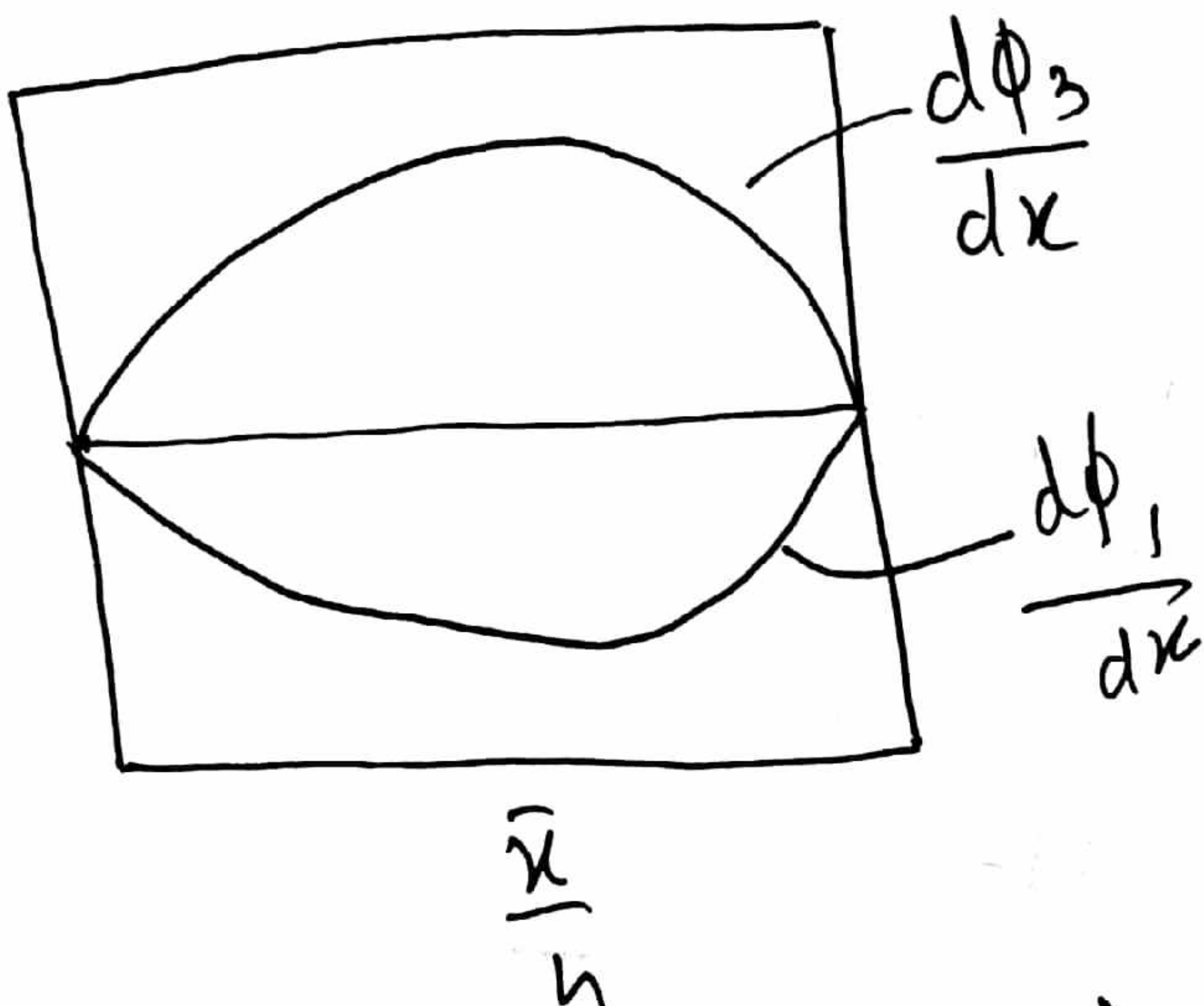
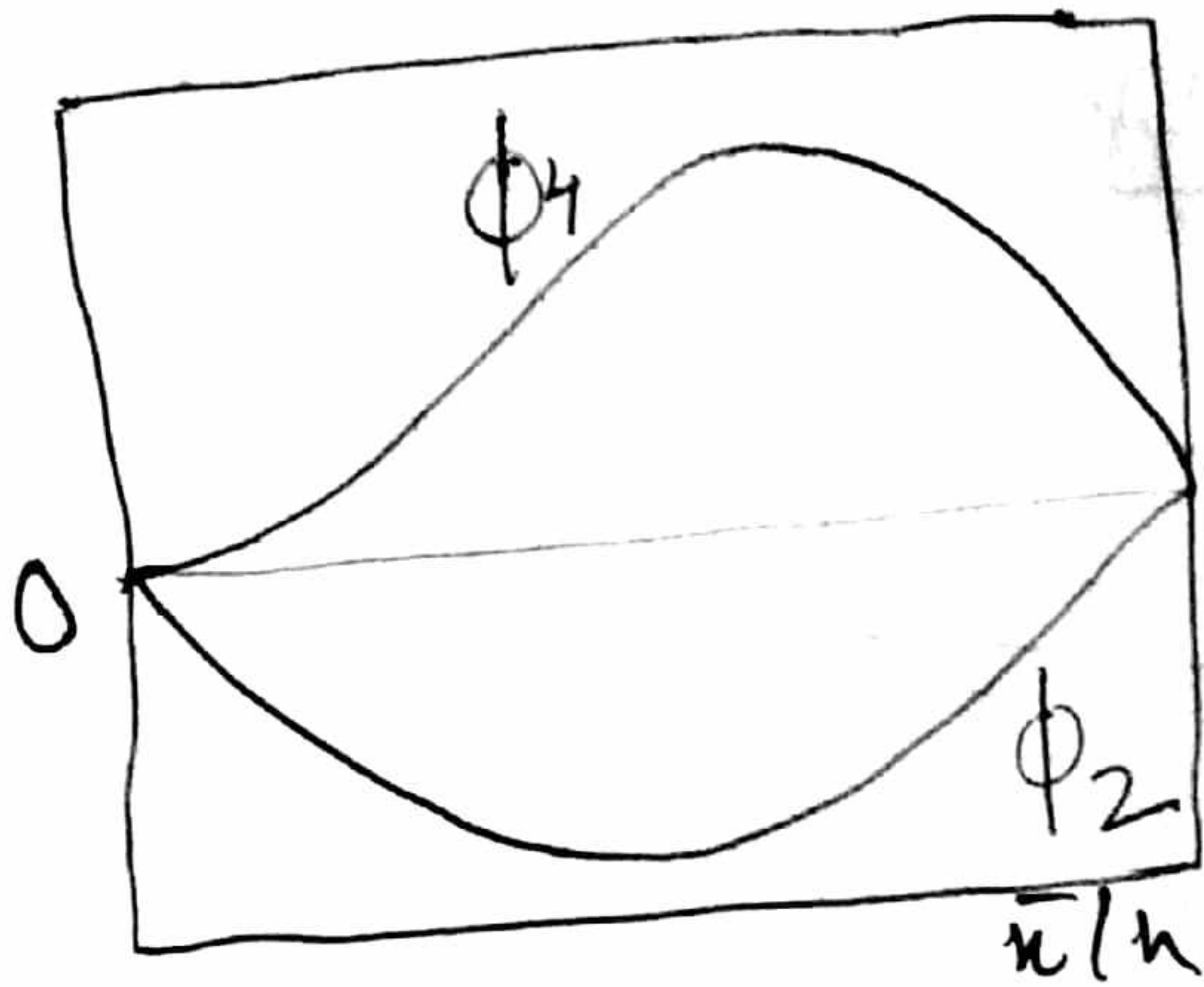
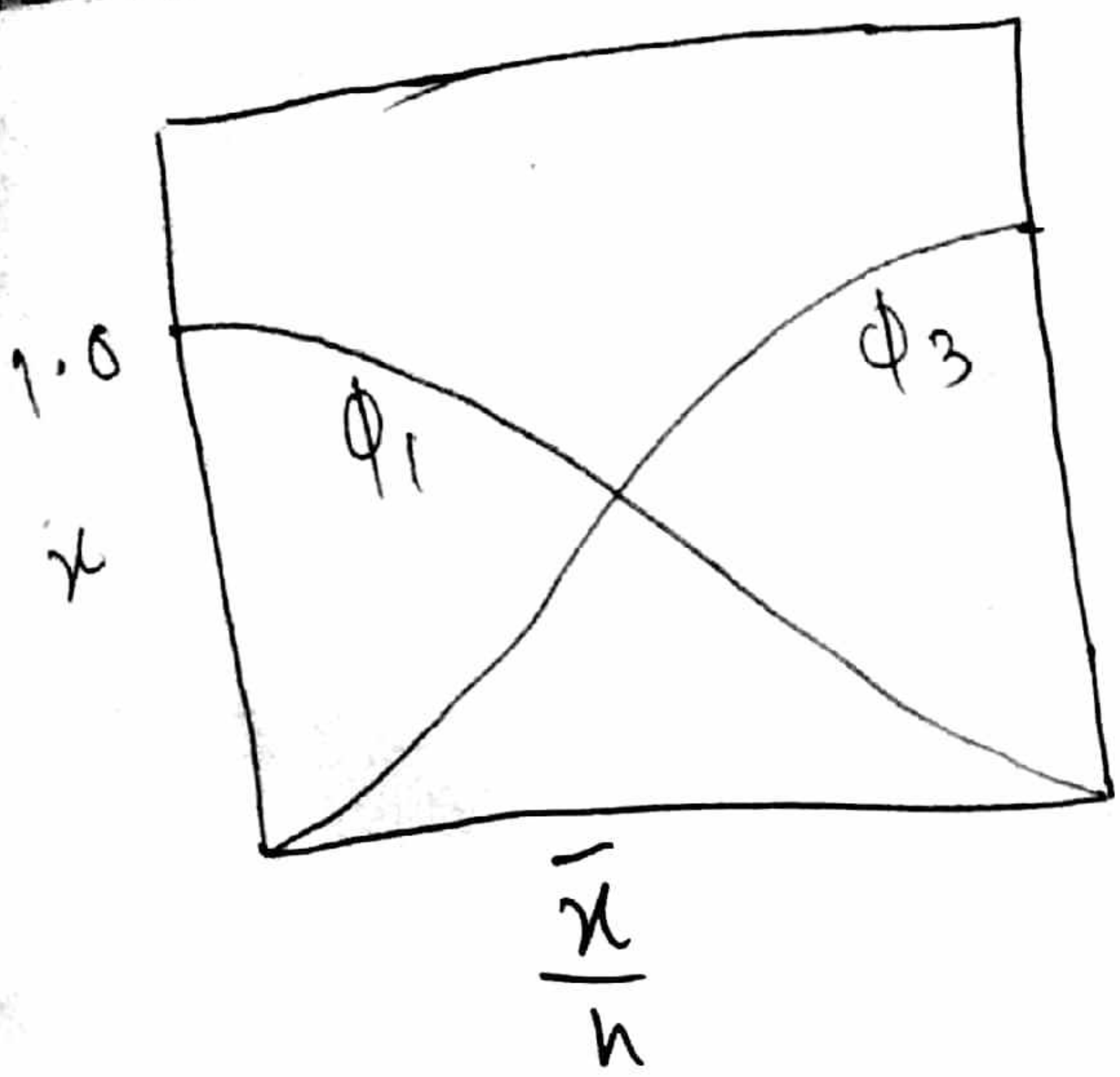
$$\frac{d\phi_2}{dx} = -\left[1 + 3\left(\frac{\bar{x}}{h}\right)^2 - 4\left(\frac{\bar{x}}{h}\right)\right]$$

$$\frac{d\phi_3}{dx} = -\frac{d\phi_1}{dx}$$

$$\frac{d\phi_4}{dx} = -\frac{\bar{x}}{h} \left[\frac{3\bar{x}}{h} - 2\right]$$

Written in Local coordinates.  
No need to memorize,  
will be given in the exams

hermite cubic interpolation  
Functions



$$\int_0^h b \frac{d^2 \phi_i}{dx^2} \sum_{j=1}^4 \frac{d^2 (u_j \phi_j)}{dx^2} dx - \int_0^h \phi_i f dx = \{Q\}$$

$$[K] \{u\} = \{F\} + \{Q\}$$

$$[K] = \frac{2b}{h^3} \begin{bmatrix} 6 & -3h & -6 & -3h \\ -3h & 2h^2 & 3h & h^2 \\ -6 & 3h & 6 & 3h \\ -3h & h^2 & 3h & 2h^2 \end{bmatrix}$$

$$\{u\} = \begin{cases} u_1 \rightarrow \text{displacement} \\ u_2 \rightarrow \text{slope} \\ u_3 \rightarrow \text{displ.} \\ u_4 \rightarrow \text{slope} \end{cases}$$

$$\{F\} = \frac{fh}{12} \begin{Bmatrix} 6 \\ -h \\ 6 \\ h \end{Bmatrix} + \begin{cases} Q_1 \rightarrow \text{shear force} \\ Q_2 \rightarrow \text{bending moment} \\ Q_3 \rightarrow \text{shear force} \\ Q_4 \rightarrow \text{bending moment} \end{cases}$$

$$\Rightarrow \{K\} \{u\} = \{F\}$$

# Assembly

Two linear elements

For 1 - DOF

$$\begin{bmatrix} k_{11} & k_{12} \\ k_{21} & k_{22} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix}$$

For one element

Assembly

$$\begin{bmatrix} k_{11}^1 & k_{12}^1 & 0 \\ k_{21}^1 & k_{22}^1 + k_{11}^2 & k_{12}^2 \\ 0 & k_{21}^2 & k_{22}^2 \end{bmatrix}$$

For 2 - DOF

$$\begin{bmatrix} k_{11} & k_{12} & k_{13} & k_{14} \\ k_{21} & k_{22} & k_{23} & k_{24} \\ k_{31} & k_{32} & k_{33} & k_{34} \\ k_{41} & k_{42} & k_{43} & k_{44} \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{Bmatrix}$$

K assembly

$$\begin{bmatrix} k_{11} & k_{12} & k_{13} & k_{14} & 0 & 0 & 0 & 0 \\ k_{21} & k_{22} & k_{23} & k_{24} & 0 & 0 & 0 & 0 \\ k_{31} & k_{32} & k_{33}^1 + k_{11}^2 & k_{34}^1 + k_{12}^2 & k_{13} & k_{14} & 0 & 0 \\ k_{41} & k_{42} & k_{43}^1 + k_{21}^2 & k_{44}^1 + k_{22}^2 & k_{23} & k_{24} & 0 & 0 \\ 0 & 0 & k_{32} & k_{31} & k_{33} & k_{34} & 0 & 0 \\ 0 & 0 & k_{42} & k_{41} & k_{43} & k_{44} & 0 & 0 \end{bmatrix}$$

|            |                       |            |
|------------|-----------------------|------------|
| $k_{11}^1$ | $k_{12}^1$            | 0          |
| $k_{22}^1$ | $k_{22}^1 + k_{11}^2$ | $k_{12}^2$ |
| 0          | $k_{21}^2$            | $k_{22}^2$ |



EI → frictional rigidity

~~[k]~~  $U_1 = 0, U_2 = 0$

$$\frac{2b}{h^3} \begin{bmatrix} 6 & -3h & -6 & -3h \\ -3h & 2h^2 & 3h & h^2 \\ -6 & 3h & 6 & 3h \\ -3h & h^2 & 3h & 2h^2 \end{bmatrix} \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ -P \\ 0 \end{bmatrix}$$

$$\frac{2b}{h^3} \begin{bmatrix} 6 & 3h \\ 3h & 2h^2 \end{bmatrix} \begin{bmatrix} u_3 \\ u_4 \end{bmatrix} = \begin{bmatrix} -P \\ 0 \end{bmatrix}$$

$$\frac{2b}{h^3} (6u_3 + 3hu_4) = -P$$

$$\frac{2b}{h^3} (3hu_3 + 2h^2u_4) = 0 \Rightarrow u_4 = -\frac{3}{2h}u_3$$

$$\frac{2b}{h^3} (6u_3 + 3h(-\frac{3}{2h}u_3)) = -\frac{Ph^3}{2b}$$

$$(6 - \frac{9}{2})u_3 = -\frac{Ph^3}{2b}$$

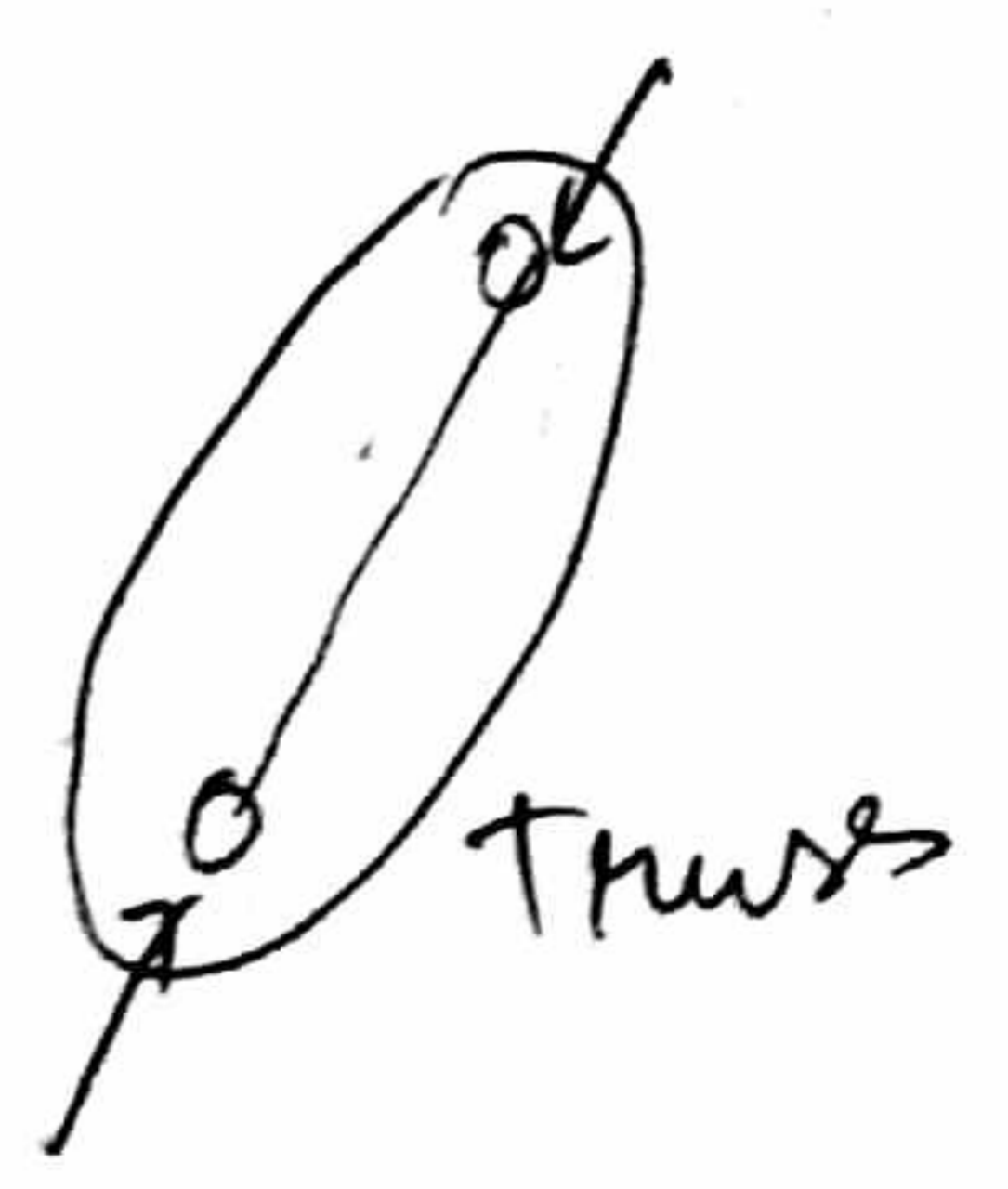
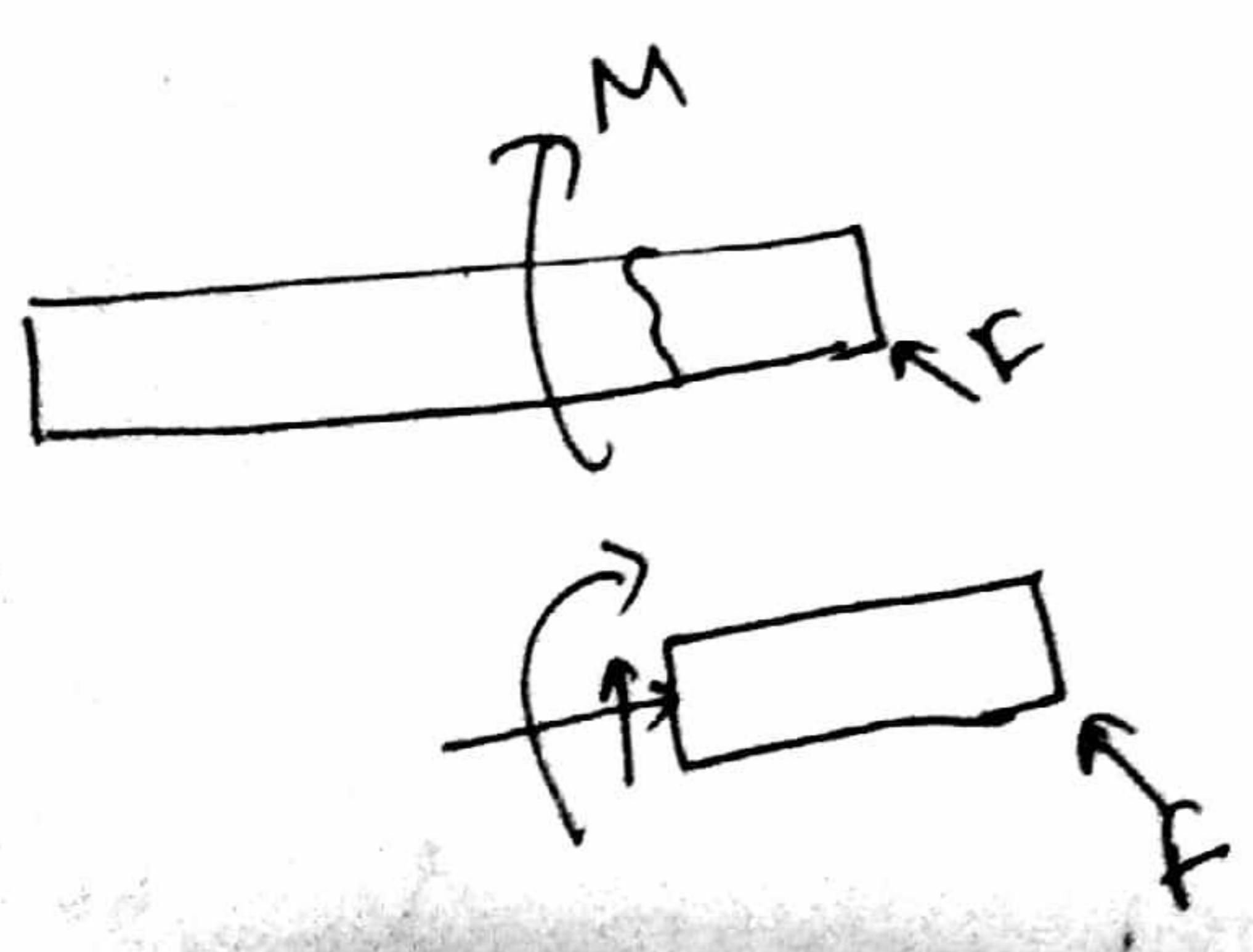
$$u_3 = -\frac{Ph^3}{3b}$$

Post processing

Difference b/w truss and ~~structure~~ frame

Truss is always a 2 force member.

Frame always has some moment.





# Generalised problem (generalised load)

$$[k]_{4 \times 4} = \frac{2EI}{h^3}$$

$$\begin{bmatrix} M & 0 & 0 & -M & 0 & 0 \\ 0 & 6 & -3h & 0 & -6 & -3h \\ 0 & -3h & 2h^2 & 0 & 3h & h^2 \\ -M & 0 & 0 & M & 0 & 0 \\ 0 & -6 & 3h & 0 & 6 & 3h \\ 0 & -3h & h^2 & 0 & 3h & 2h^2 \end{bmatrix}$$

$$\begin{bmatrix} U_1 \\ U_2 \\ U_3 \\ U_4 \\ U_5 \\ U_6 \end{bmatrix}$$

$$M = \frac{Ah^2}{2I}$$

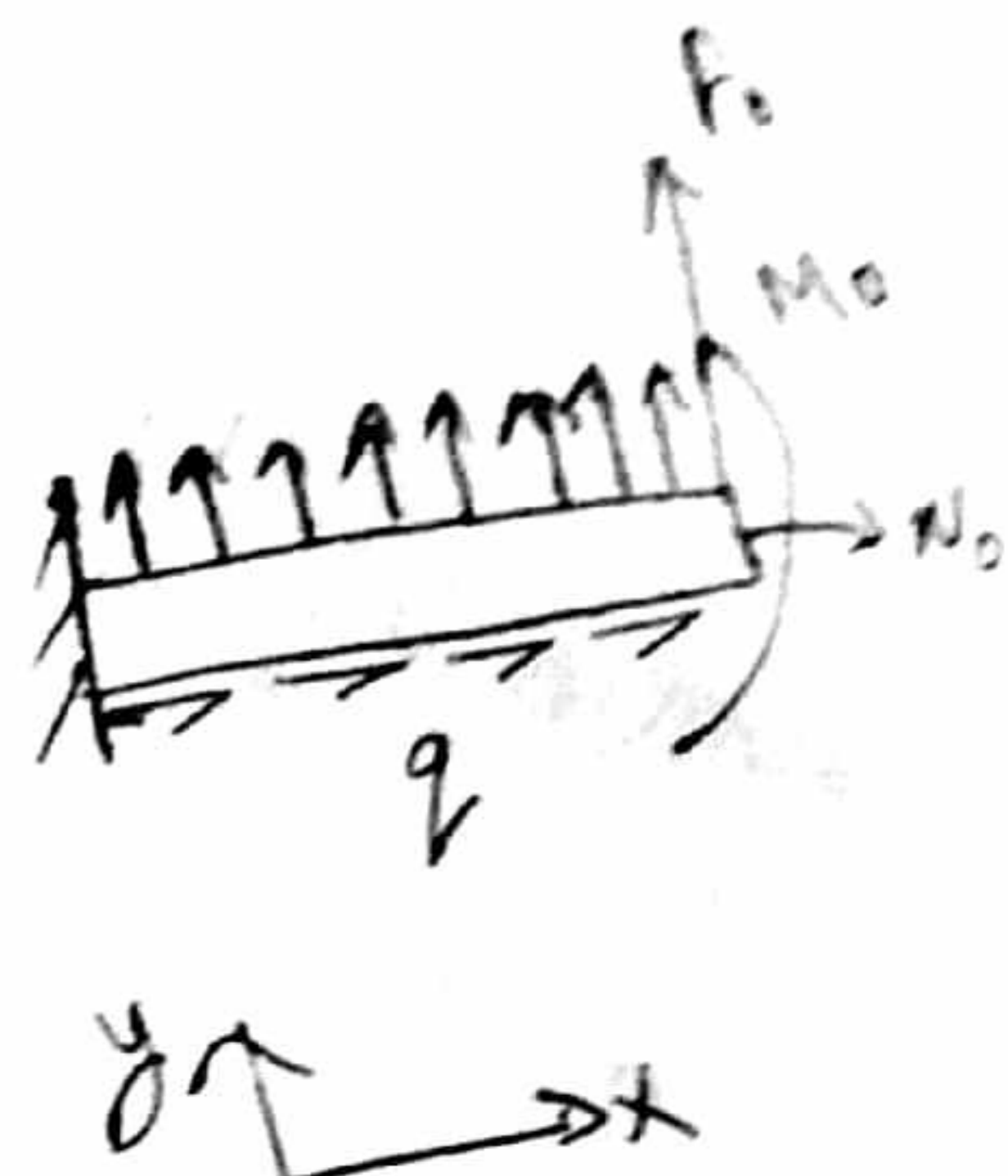
These eqns are only for conditions where deflection are not very large.

Two nodal element

1st & 4th are due to the axial loading

$$\{F\} = \begin{bmatrix} \frac{1}{2}qh \\ \frac{1}{2}fh \\ -\frac{1}{12}fh^2 \\ \frac{1}{2}qh \\ \frac{1}{2}fh \\ \frac{1}{12}fh^2 \end{bmatrix}$$

$$\{D\} = \begin{bmatrix} D_1 \\ D_2 \\ D_3 \\ D_4 \\ D_5 \\ D_6 \end{bmatrix}$$



For two elements

|            |            |
|------------|------------|
| $k_{11}^1$ | $k_{12}^1$ |
| $k_{21}^1$ | $k_{22}^1$ |

|            |            |
|------------|------------|
| $k_{11}^2$ | $k_{12}^2$ |
| $k_{21}^2$ | $k_{22}^2$ |

|            |                       |            |
|------------|-----------------------|------------|
| $k_{11}^1$ | $k_{12}^1$            | 0          |
| $k_{21}^1$ | $k_{22}^1 + k_{11}^2$ | $k_{12}^2$ |
| 0          | $k_{21}^2$            | $k_{22}^2$ |

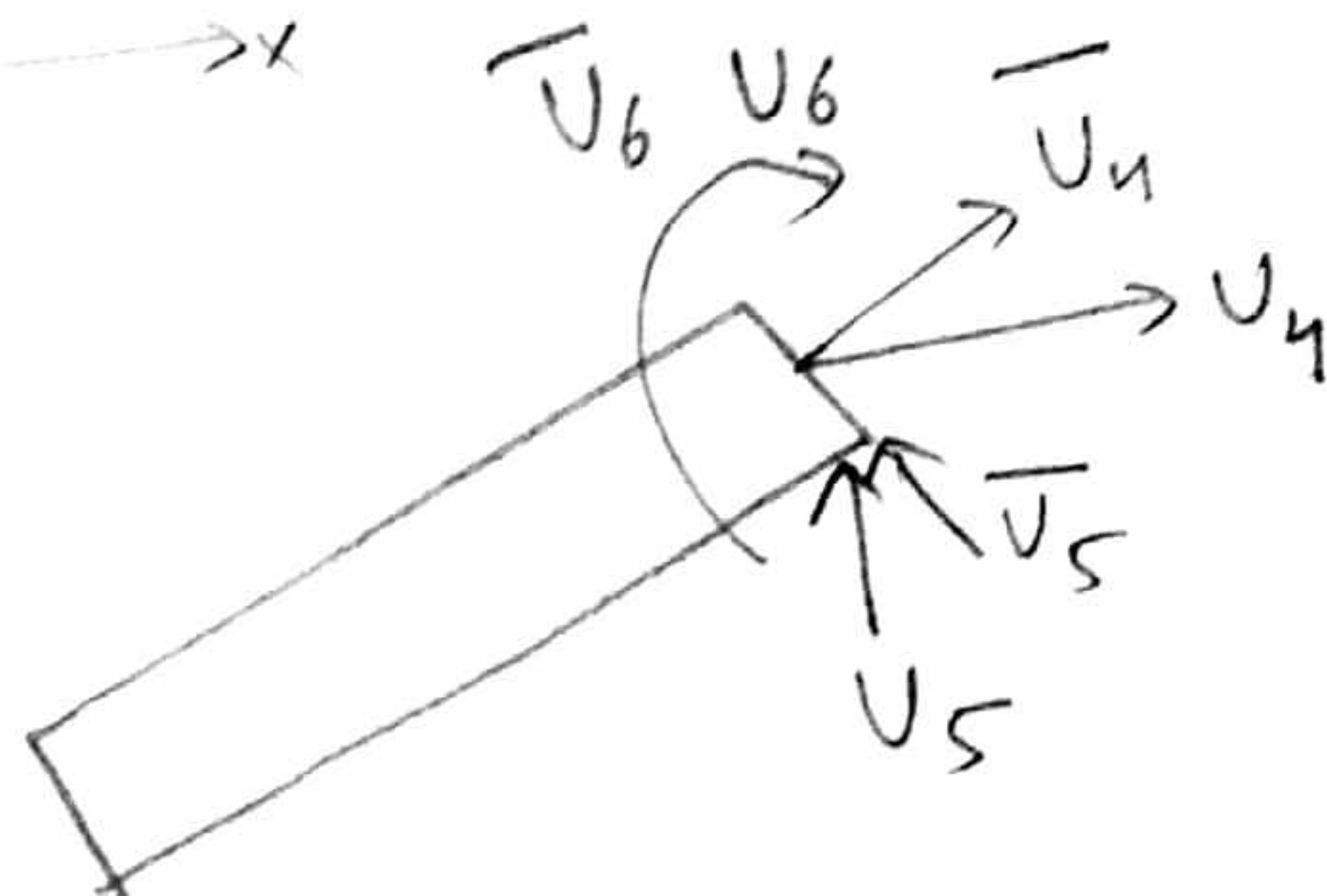
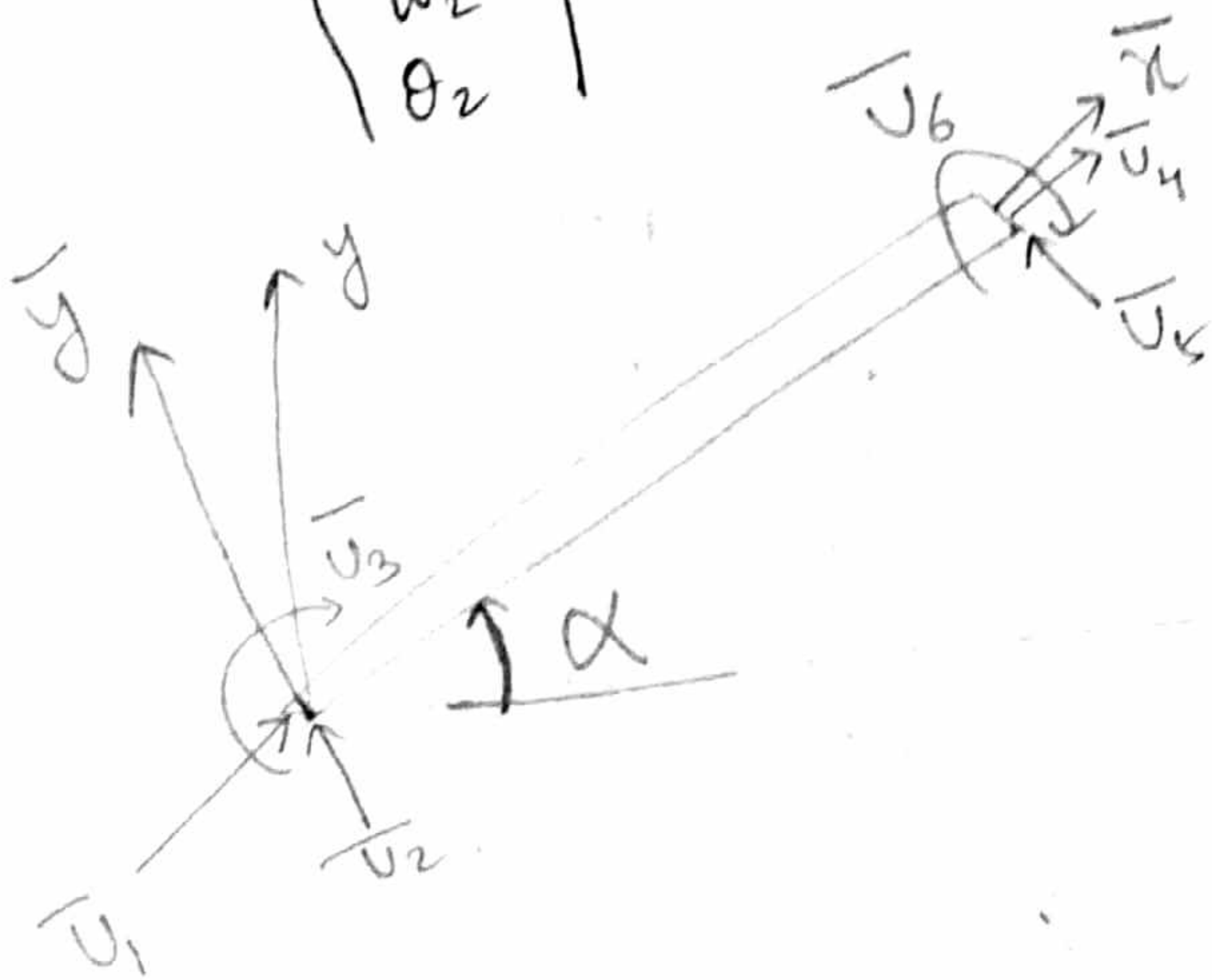
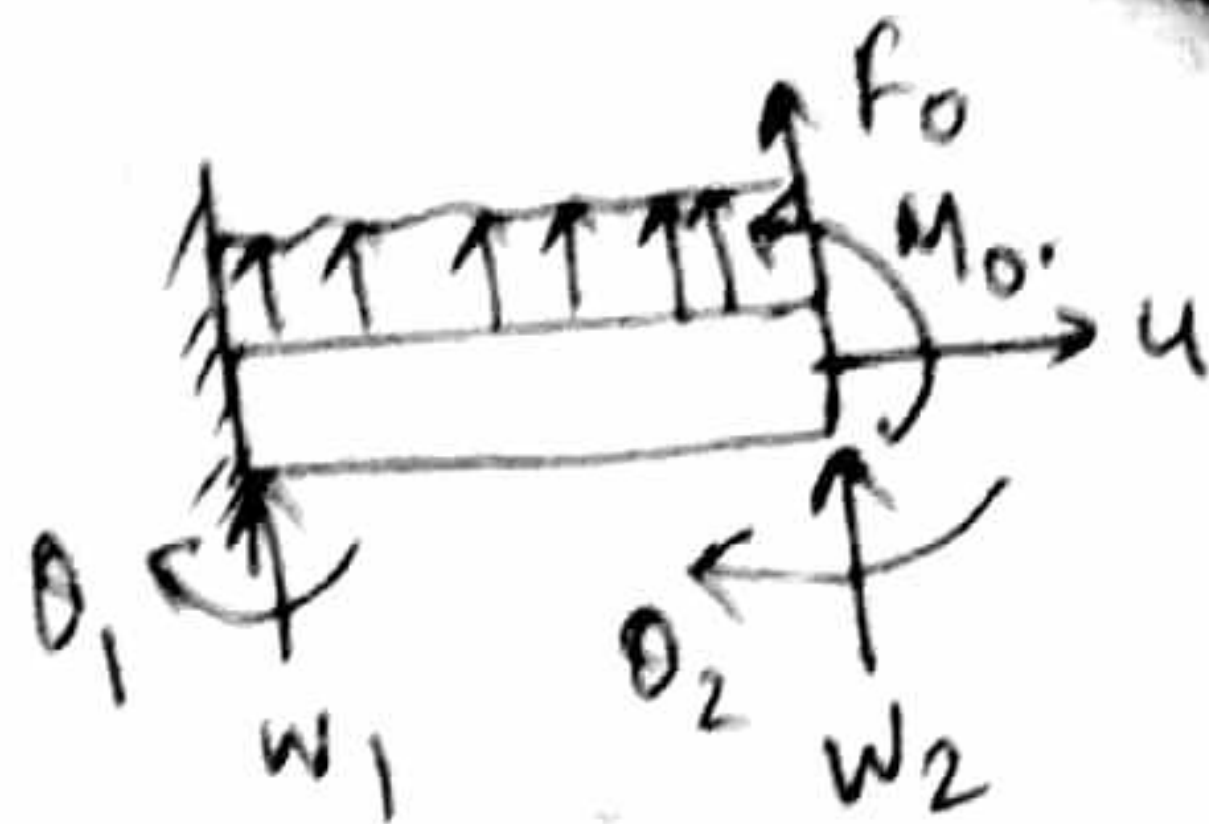
$$\{U\} = \begin{bmatrix} U_1 \\ U_2 \\ U_3 \\ U_4 \\ U_5 \\ U_6 \\ U_7 \\ U_8 \\ U_9 \end{bmatrix}$$

$$\{F\} = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 + f_1 \\ f_5 + f_2 \\ f_6 + f_3 \\ f_7 \\ f_8 \\ f_9 \end{bmatrix}$$

$$\{D\} = \begin{bmatrix} D_1 \\ D_2 \\ D_3 \\ D_4 + D_1 \\ D_5 + D_2 \\ D_6 + D_3 \\ D_7 \\ D_8 \\ D_9 \end{bmatrix}$$

$$\{u\} = \begin{Bmatrix} u_1 \\ w_1 \\ \theta_1 \\ u_2 \\ w_2 \\ \theta_2 \end{Bmatrix}$$

$$\{F\} = \begin{Bmatrix} N_1 \\ V_1 \\ M_1 \\ N_2 \\ V_2 \\ M_2 \end{Bmatrix}$$



$$\begin{bmatrix} \bar{u} \\ \bar{v} \\ \bar{w} \end{bmatrix} = \begin{bmatrix} \cos\alpha & \sin\alpha & 0 \\ -\sin\alpha & \cos\alpha & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix}$$

$$\begin{Bmatrix} \bar{u} \\ \bar{v} \end{Bmatrix} = \begin{bmatrix} \cos\alpha & \sin\alpha \\ -\sin\alpha & \cos\alpha \end{bmatrix} \begin{Bmatrix} u \\ v \end{Bmatrix} \Rightarrow \begin{bmatrix} u \\ v \end{bmatrix} = \begin{bmatrix} \cos\alpha & -\sin\alpha \\ \sin\alpha & \cos\alpha \end{bmatrix} \begin{Bmatrix} \bar{u} \\ \bar{v} \end{Bmatrix}$$

$$\{ \bar{u} \} = [T] \{ u \} \rightarrow \text{displacements}$$

$$\{ u \} = [T]^{-1} \{ \bar{u} \} = [T]^T \{ \bar{u} \}$$

$$\{ \bar{F} \} = [T] \{ F \} \rightarrow \text{forces}$$

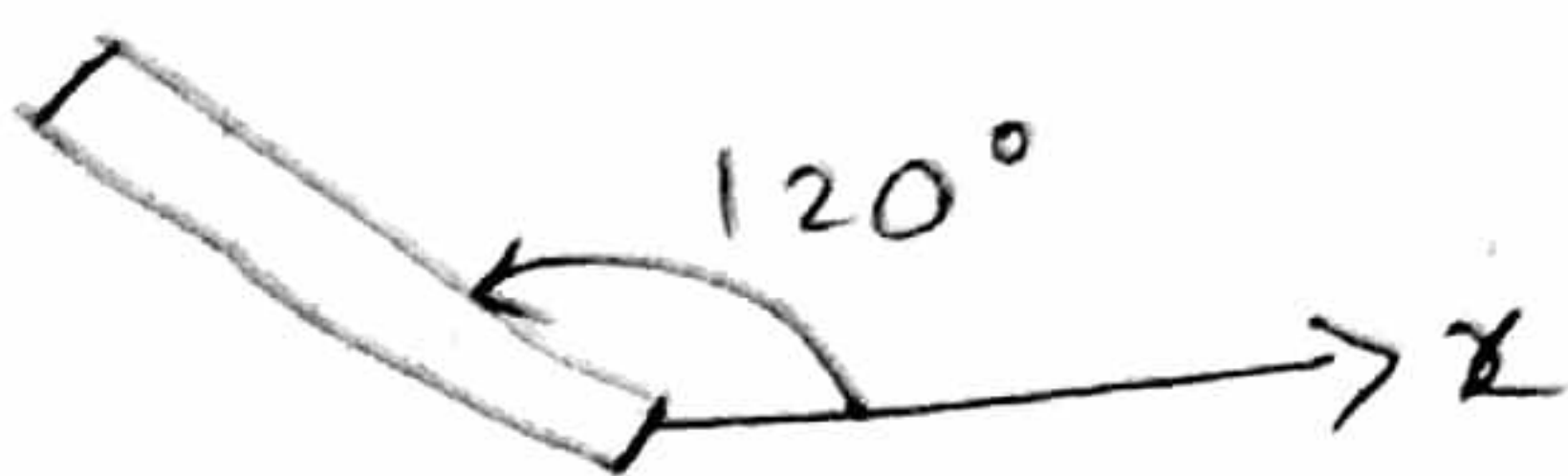
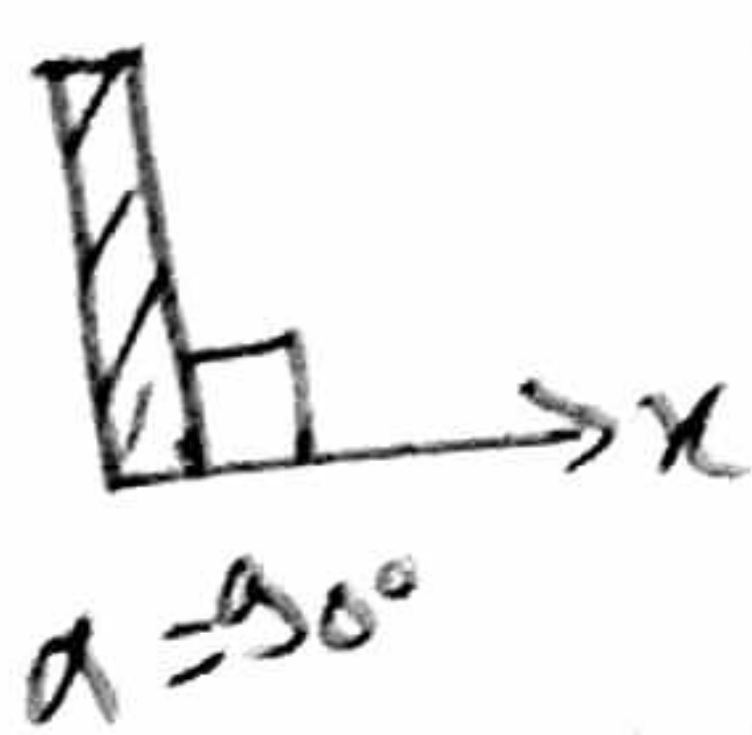
$$[K] [\bar{u}] = [\bar{F}]$$

$$[K] [T] [u] = [T] [F]$$

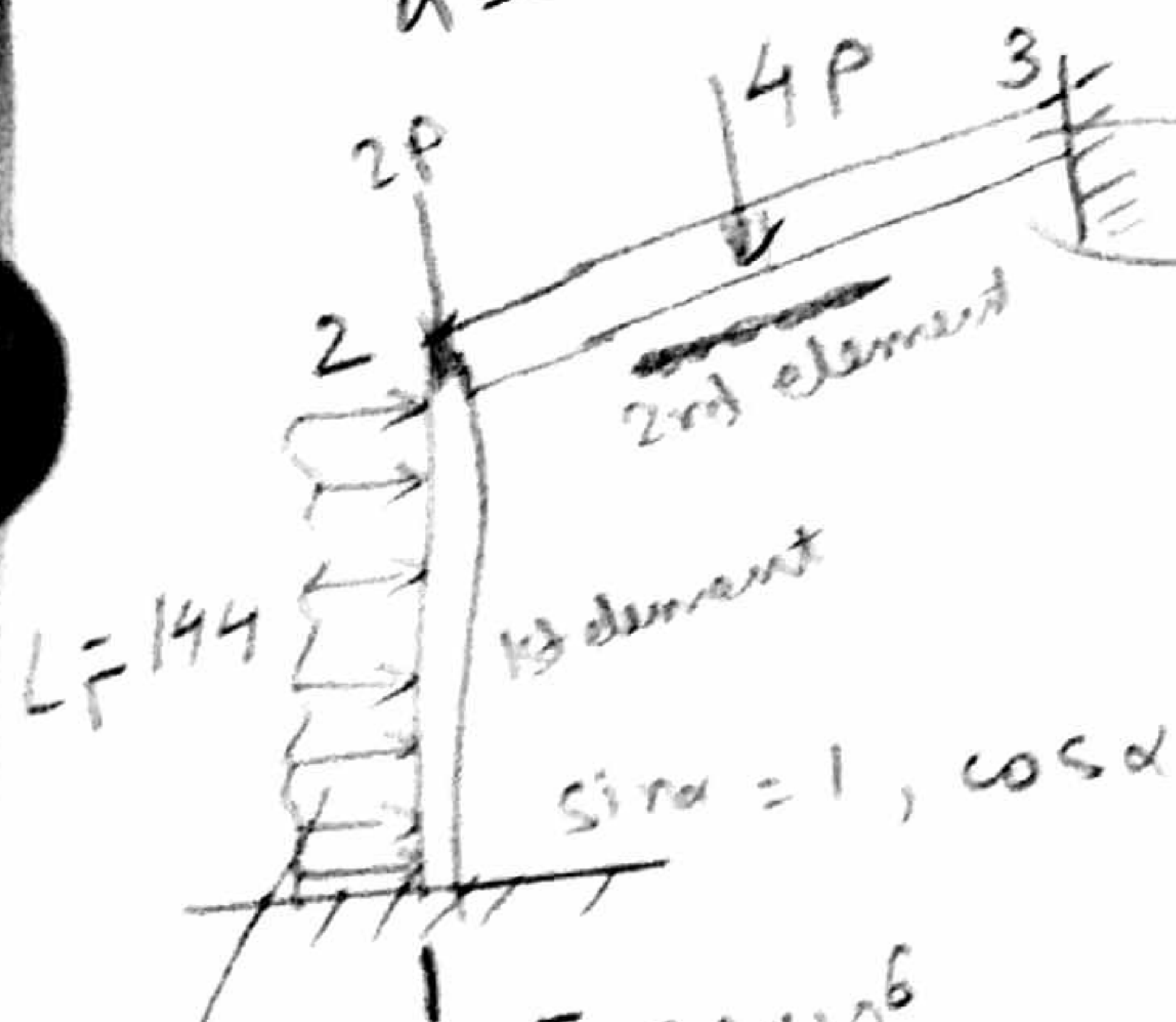
$$[T]^T [K] [T] \{ u \} = [T]^T [T] \{ F \}$$

$$[K] \{ u \} = \{ F \}$$

$\alpha$  is always measured from  $x$  axis in the counter ~~clockwise~~ clockwise direction.



Do not write  $\alpha = -60^\circ$



$L_2 = 180$   
 $\sin \alpha = 0.6$   
 $\cos \alpha = 0.8$

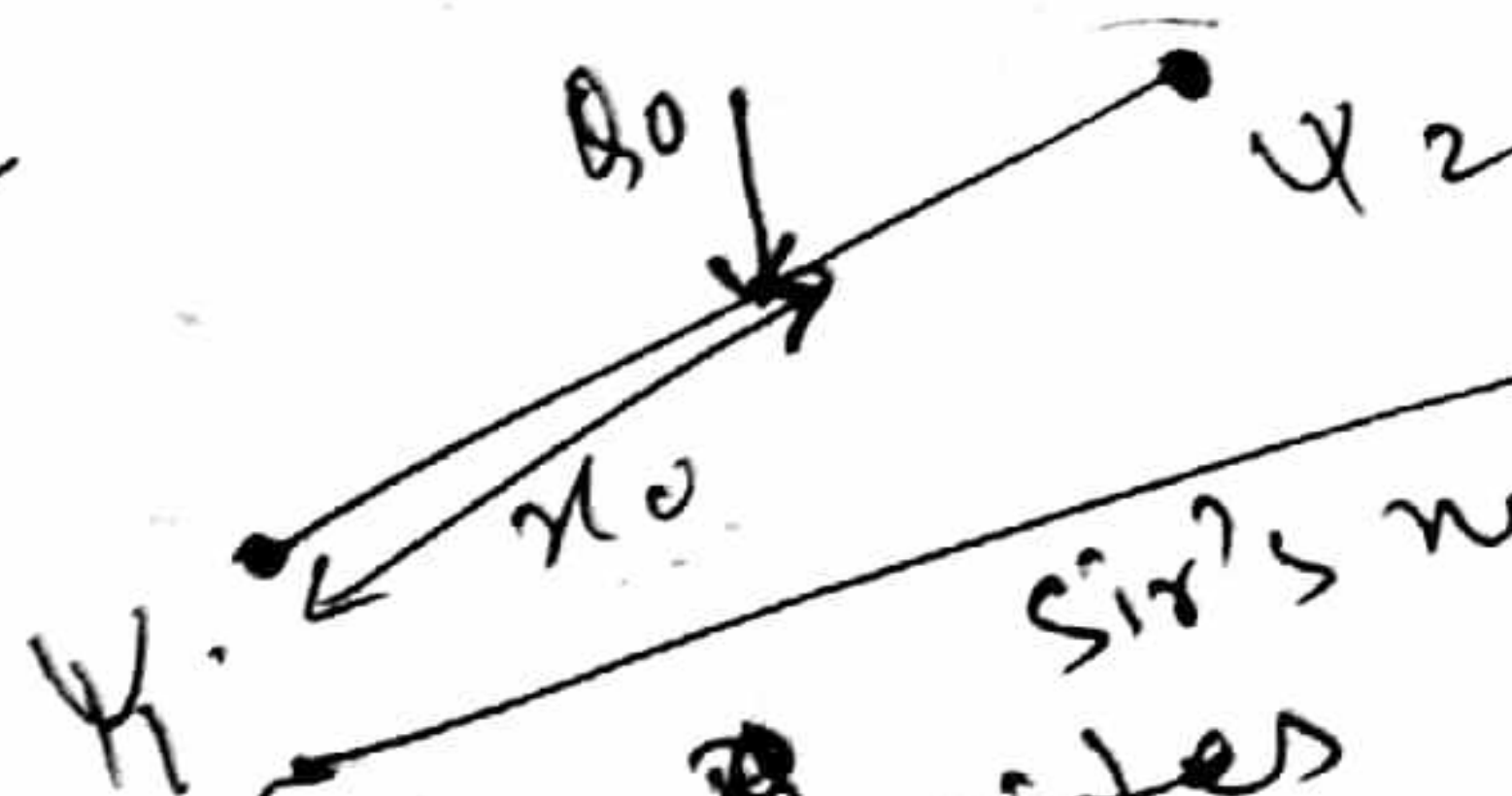
$q(x) = q_0 \delta(x - x_0)$   
 $\int_{-\infty}^{\infty} F(x) \delta(x - x_0) dx = f(x_0)$

$f_i = \int q(x) \psi_i dx$   
 $= \int q_0 \delta(x - x_0) \psi_i dx$   
 $= q_0 \psi_i(x_0)$

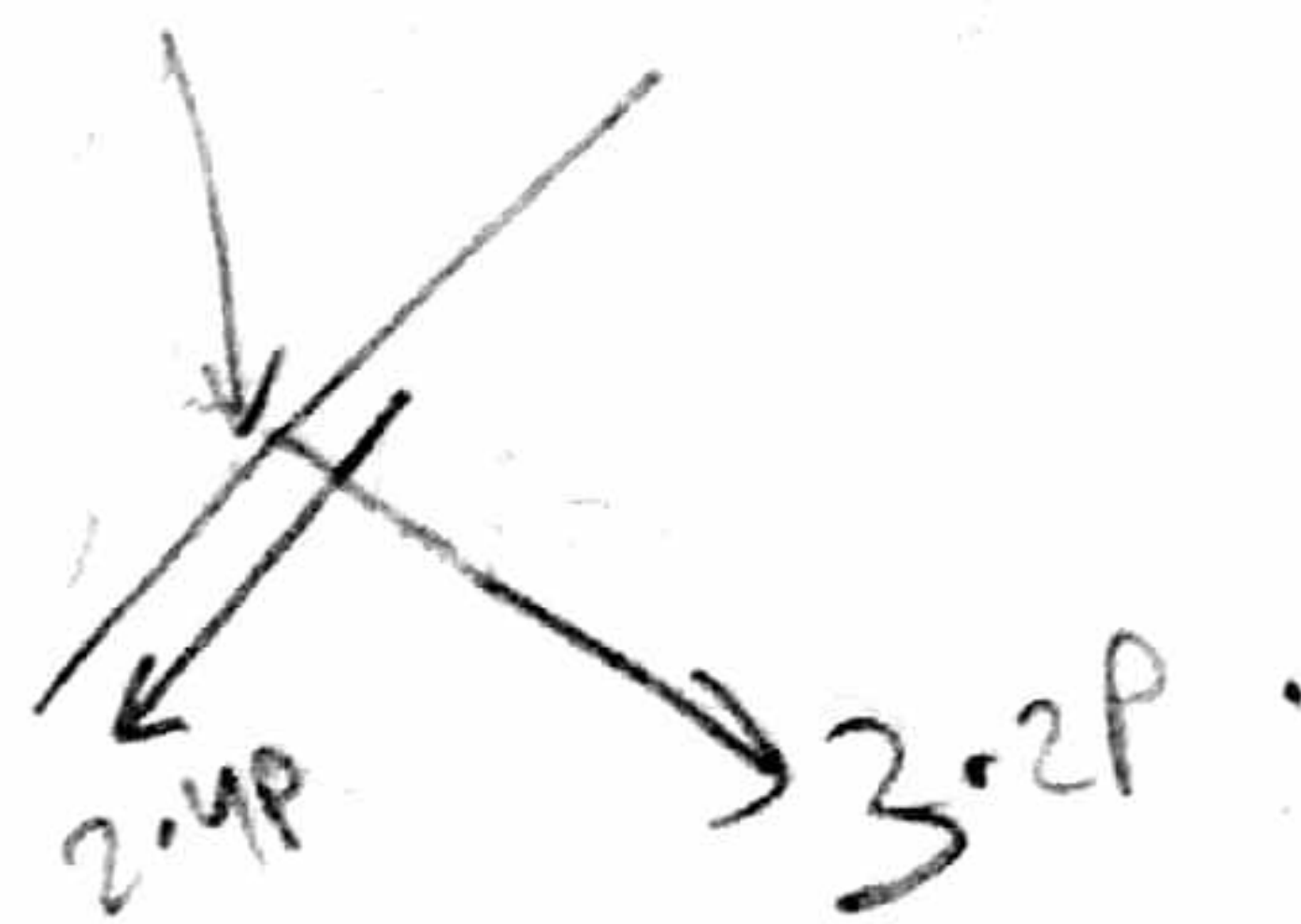
$f = \frac{q}{72}$

$E = 2 \times 10^6$

$A = 10$   
 $I = 10$



It He writes  $-1.2P$  it means load  $1.2P$  acting in left direction.



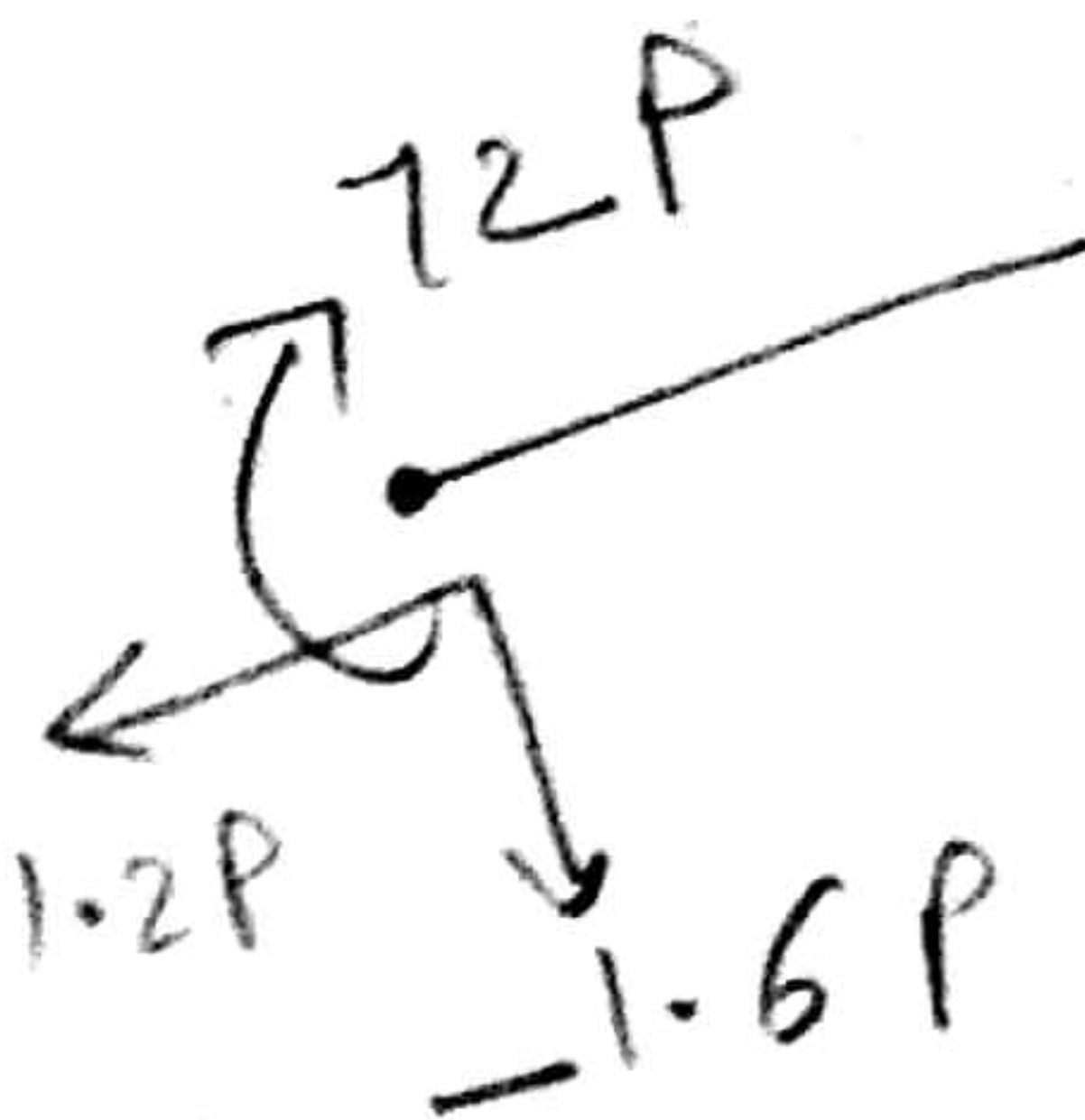
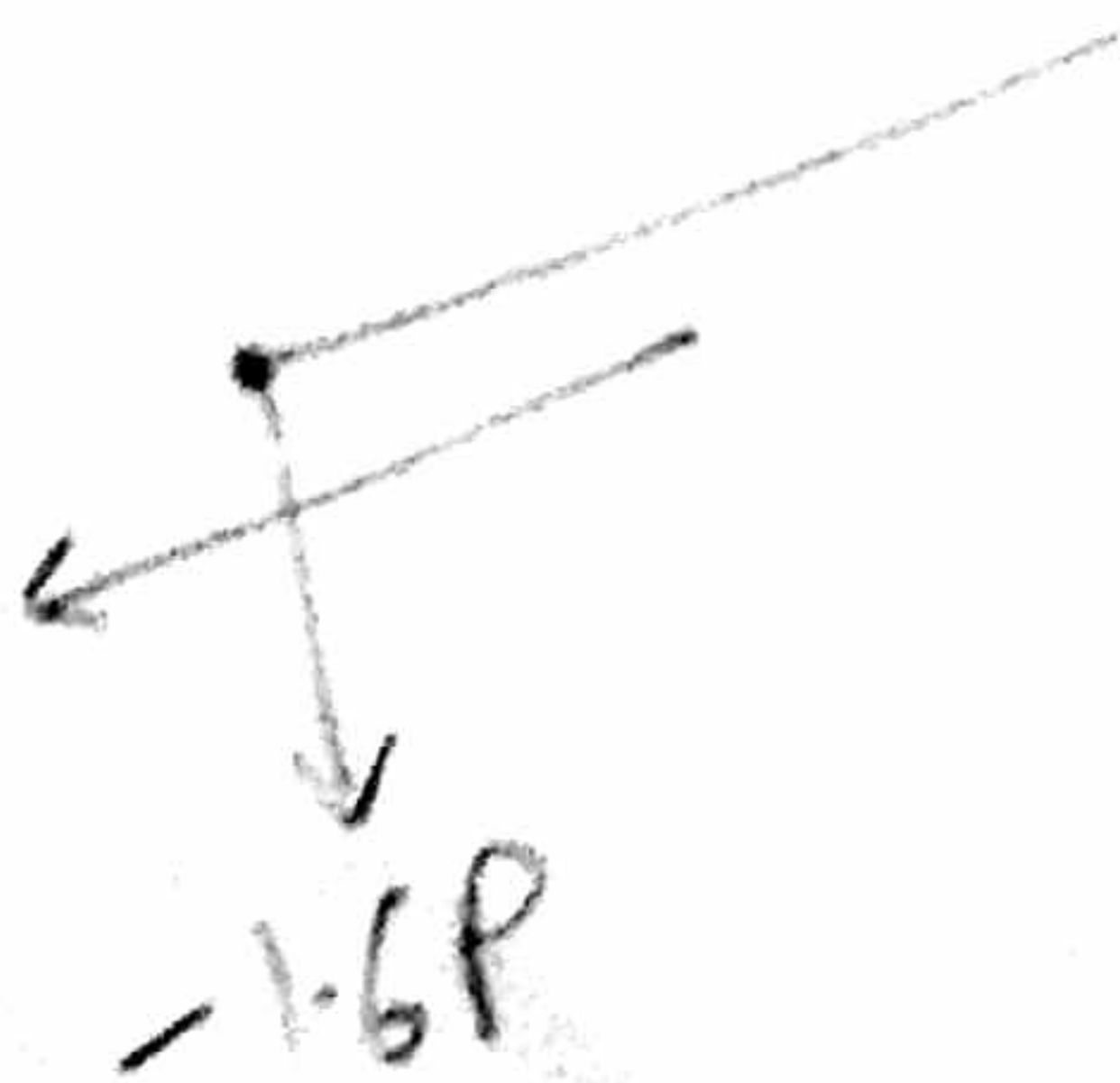
for axial load

$\psi_1 = (1 - \frac{x}{h})$

$\psi_1(x_0) = (1 - \frac{x_0}{h})$

for axial load

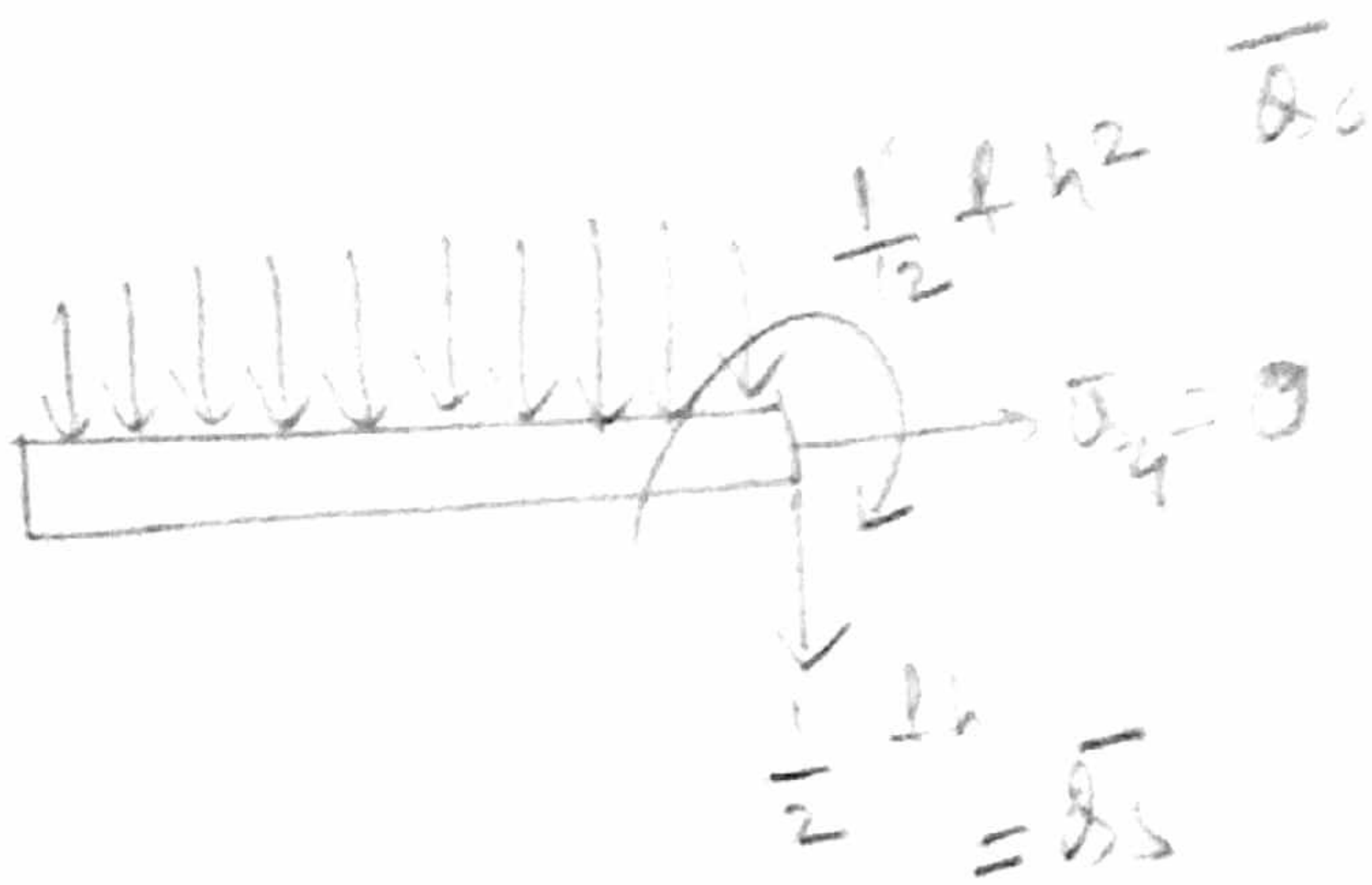
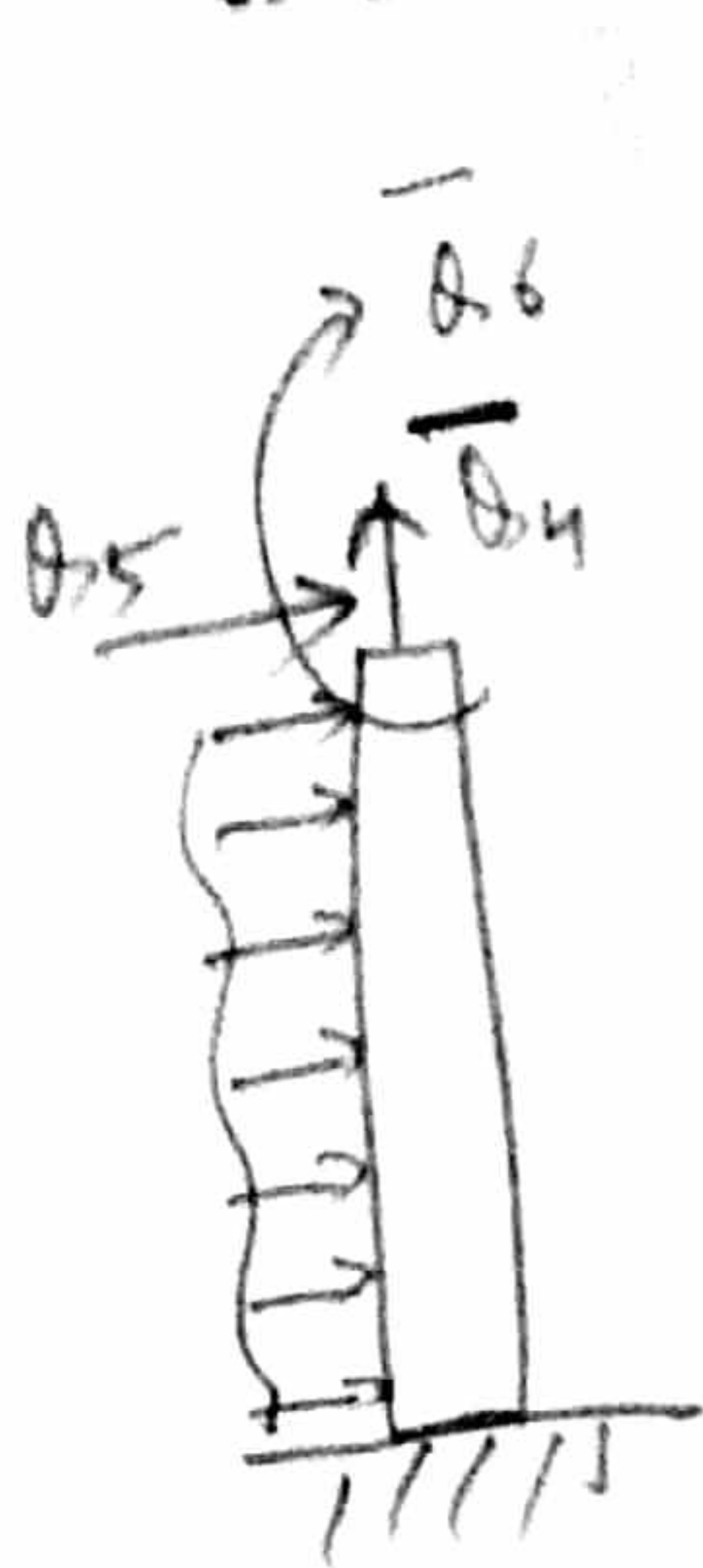
$\psi_2(x_0) = \frac{x_0}{h}$   
 $(q)_0 = q_0 (1 - \frac{x_0}{h})$   
 $q_0 = q_0 \frac{x_0}{h}$



$$\phi_1 = 1 - 3 \left( \frac{\bar{x}}{h} \right)^2 + 2 \left( \frac{\bar{x}}{h} \right)^3 = 0.5$$

$$\phi_2 = -\bar{x} \left( 1 - \frac{\bar{x}}{h} \right) = -90 * 0.5^2 = -22.5$$

$$Q_0 = Q_0 \phi_2 = -22.5 (3 \cdot 2P) = 72P$$



$$\bar{Q}_4 = 0$$

$$\bar{Q}_5 = \frac{1}{2} \frac{P}{72} \cdot 144 = -1P$$

$$\bar{Q}_6 = -\frac{1}{12} * \frac{P}{72} * (144)^2 = -24P$$

$$\begin{bmatrix} Q_4 \\ Q_5 \end{bmatrix} = \begin{bmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{bmatrix} \begin{bmatrix} \bar{Q}_4 \\ \bar{Q}_5 \end{bmatrix}$$

$$Q_4 = \cos \alpha \bar{Q}_4 - \sin \alpha \bar{Q}_5 = P$$

$$Q_5 = \sin \alpha \bar{Q}_4 + \cos \alpha \bar{Q}_5 = 0$$

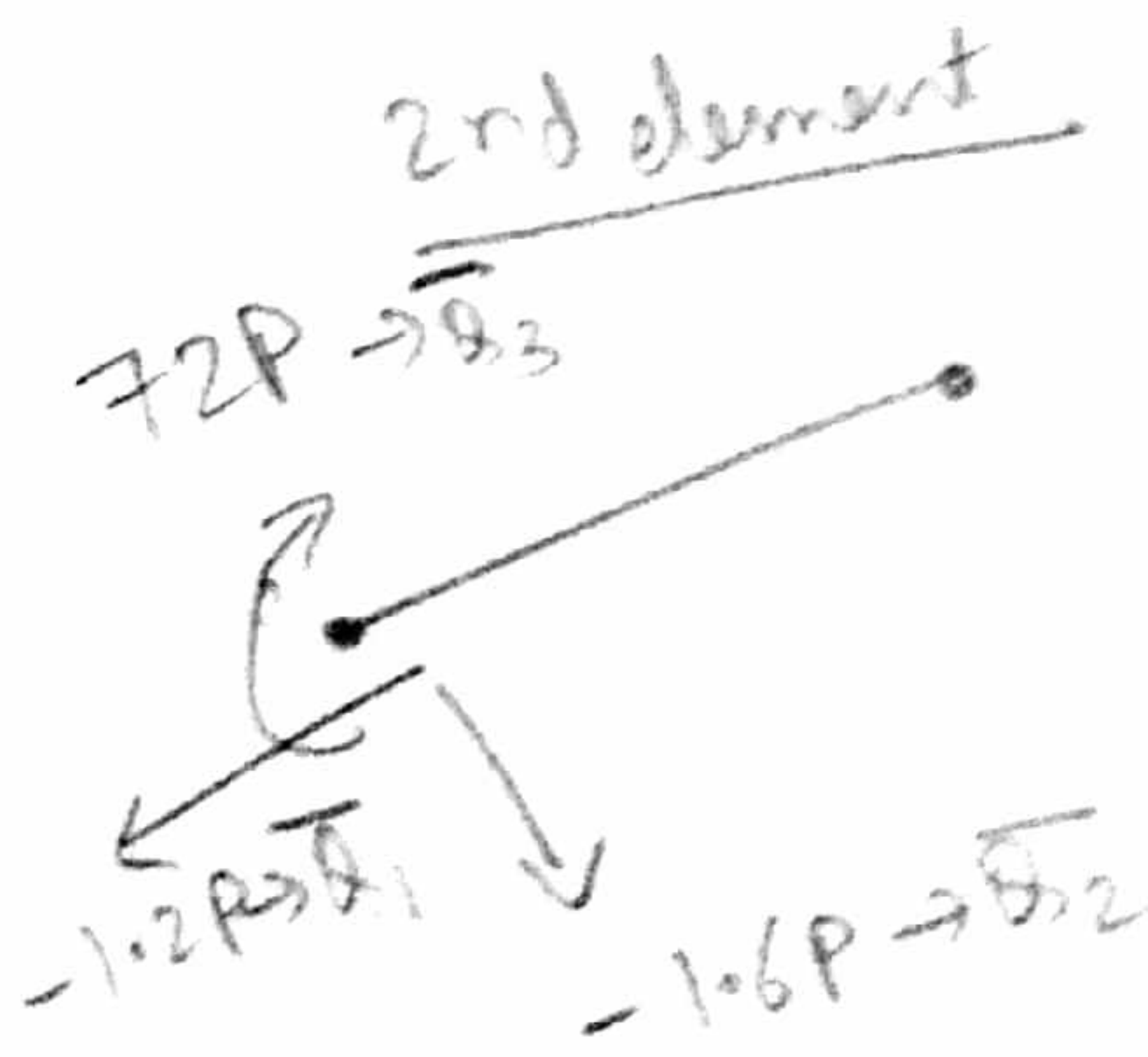
$$Q_6 = -24P$$

$$\begin{bmatrix} Q_1 \\ Q_2 \end{bmatrix} = \begin{bmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{bmatrix} \begin{bmatrix} \bar{Q}_1 \\ \bar{Q}_2 \end{bmatrix}$$

$$Q_1 = \bar{Q}_1 \cos \alpha - \sin \alpha \bar{Q}_2$$

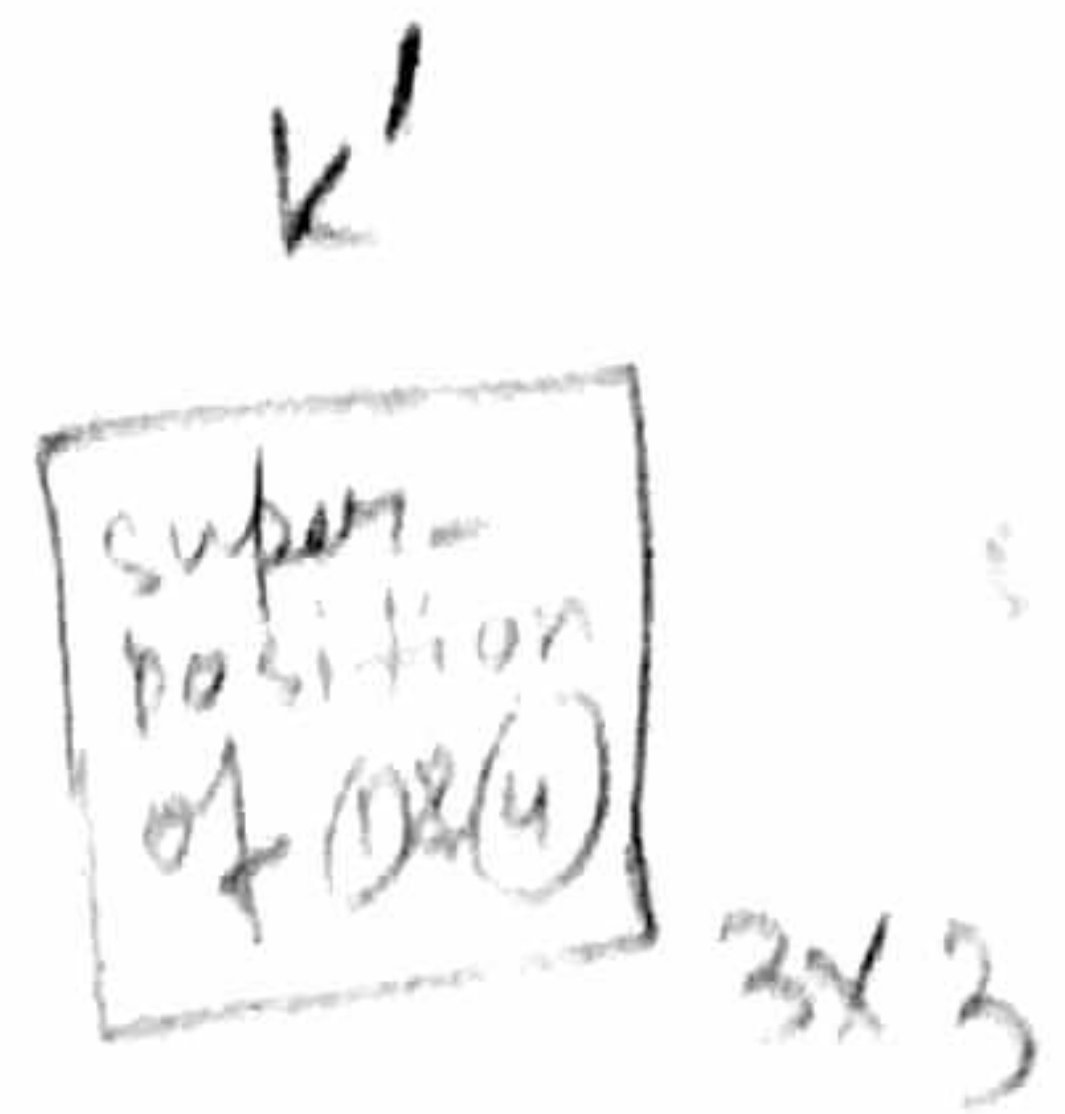
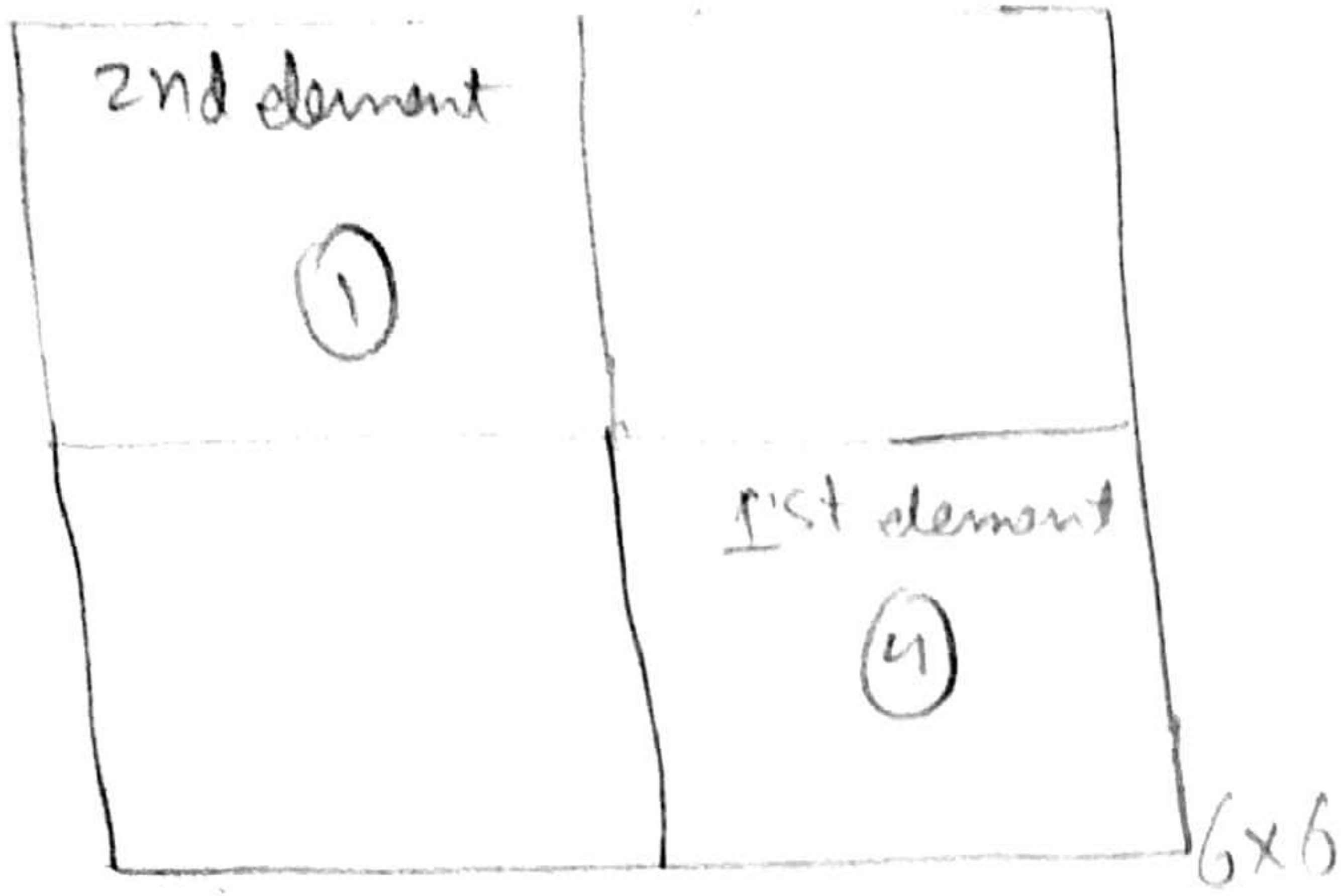
$$= 0.8(-1.2) - 0.6(-1.6)$$

$$Q_1 = 0$$



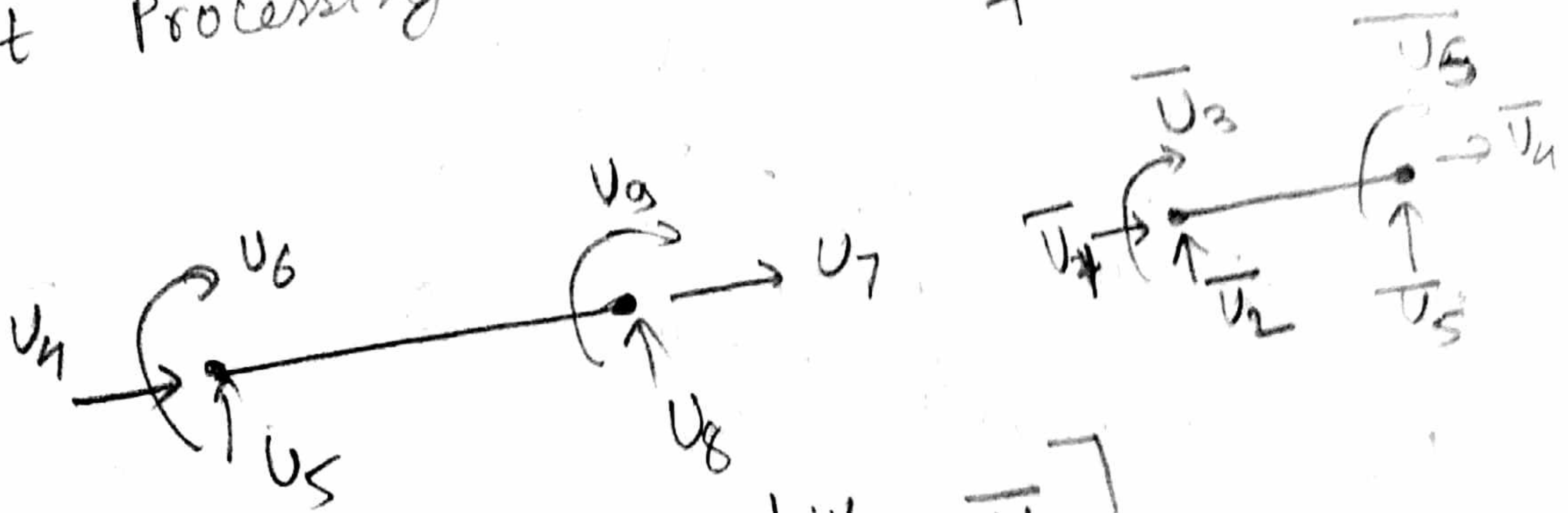
$$Q_2 = 0.6(-1.2)P + 0.8(-1.6)P$$

$$Q_2 = -2P$$



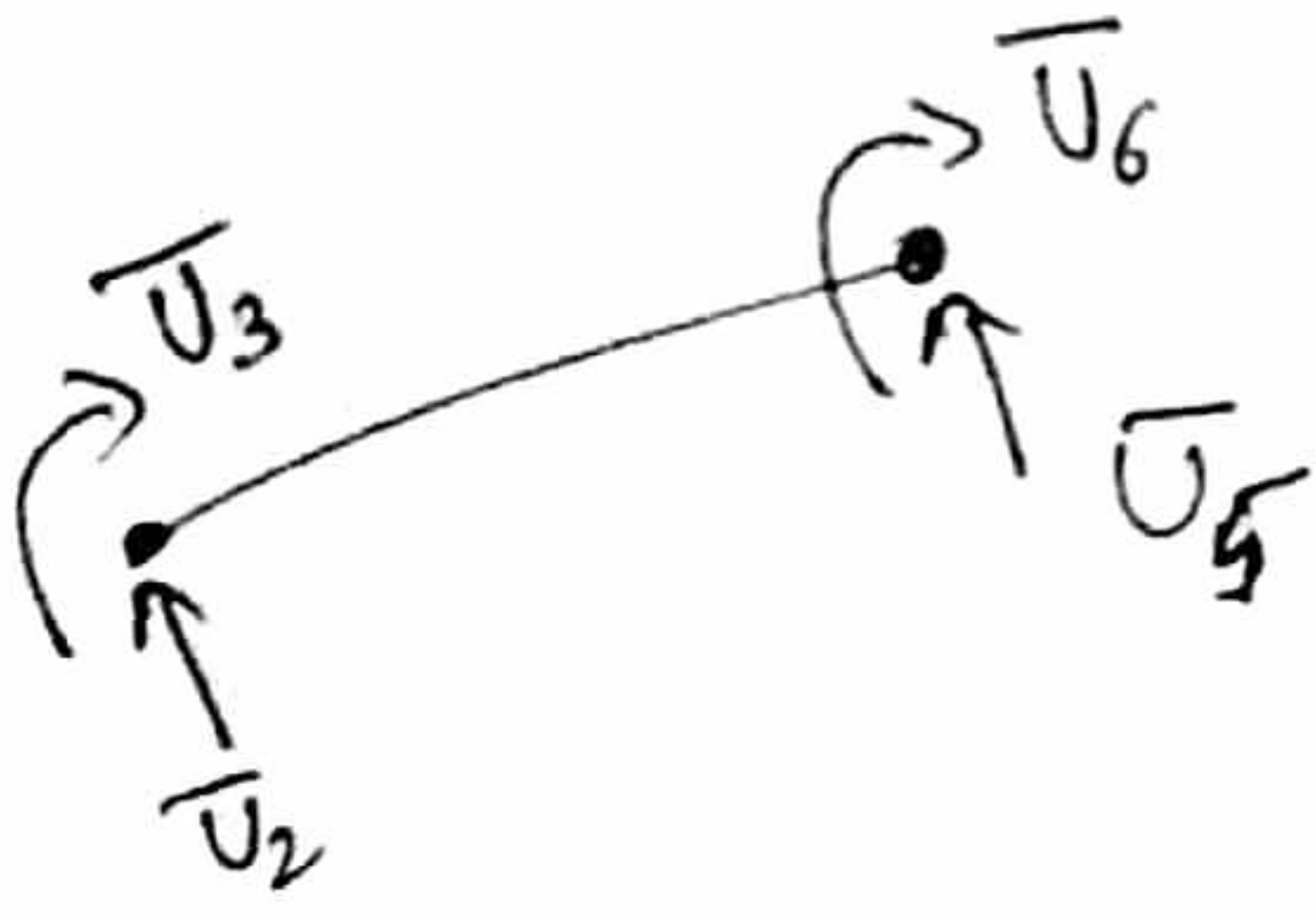
$$\begin{bmatrix} \phantom{0} \\ \phantom{0} \\ \phantom{0} \\ \phantom{0} \\ \phantom{0} \\ \phantom{0} \end{bmatrix}
 \begin{bmatrix} U_4 \\ U_5 \\ U_6 \end{bmatrix}
 =
 \begin{Bmatrix} P+0 \\ 0+(-2P) \\ -24P \\ +72P \end{Bmatrix}
 +
 \begin{Bmatrix} 0 \\ -2P \\ 0 \end{Bmatrix}
 =
 \begin{Bmatrix} P \\ -4P \\ 48P \end{Bmatrix}$$

Post Processing



$$\bar{F}_4 = EA \left[ \frac{d\psi_1}{dx} \bar{U}_1 + \frac{d\psi_2}{dx} \bar{U}_4 \right]$$

$$f_5 = \frac{d}{dx} EI \frac{d^2 w}{dx^2} \equiv EI \frac{d^3 w}{dx^3} \equiv EI \sum_{j=1}^4 \frac{d^3 \phi_j}{dx^3} \times \bar{U}_j$$



Bending moment

$$f_6 = EI \frac{d^2 w}{dx^2} = EI \sum_{j=1}^4 \left( \frac{d^2 \phi_j}{dx^2} \times U_j \right)$$

$$\sigma_{xx} = \frac{M_z y}{I_z}$$

- for  $j=1$   
 $U_3 = U_2$
- for  $j=2$   
 $U_3 = U_3$
- for  $j=3$   
 $U_3 = U_5$
- for  $j=4$   
 $U_3 = U_6$

\* Notes :-

- ① Anything axial  $\rightarrow$  Lagrangian interpolation  $f^n$
- ② Anything transverse  $\rightarrow$  Hermite interpolation  $f^n$

$$\begin{aligned}
 \Pi &= \int_0^h \frac{1}{2} \sigma E A dx - P \delta_2 - Q \delta_1 \\
 &= \int_0^h \frac{1}{2} (E \epsilon) E A dx - P \delta_2 + Q \delta_1 \\
 &= \frac{1}{2} EA \int_0^h \epsilon^2 dx - P \delta_2 + Q \delta_1 \\
 &= \frac{1}{2} EA \int_0^h \left( \frac{du}{dx} \right)^2 dx - P \delta_2 + Q \delta_1 \\
 &= \frac{1}{2} EA \int_0^h \left( \frac{d}{dx} (u_1 \psi_1 + u_2 \psi_2) \right)^2 dx - P \delta_2 + Q \delta_1 = 0
 \end{aligned}$$

$$\frac{\partial \Pi}{\partial u_1} = EA \int_0^h \frac{d\psi_1}{dx} \left( \frac{d}{dx} (\psi_1 u_1 + \psi_2 u_2) \right) dx + Q = 0$$

$$\frac{\partial \Pi}{\partial u_2} = EA \int_0^h \frac{d\psi_2}{dx} \left( \frac{d}{dx} (\psi_1 u_1 + \psi_2 u_2) \right) dx - P = 0$$

$$\frac{EA}{h} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \begin{Bmatrix} u_1 \\ u_2 \end{Bmatrix} = \begin{Bmatrix} -B \\ P \end{Bmatrix}$$

Internal strain energy

$$dU = \sigma_x dydz d(u + \frac{\partial u}{\partial x} dx)$$

$$= \sigma_x dydz du$$

$$dU = \sigma_x dydz d\left(\frac{\partial u}{\partial x}\right) dx$$

$$dU = \sigma_x d\left(\frac{\partial u}{\partial x}\right) dx dy dz$$

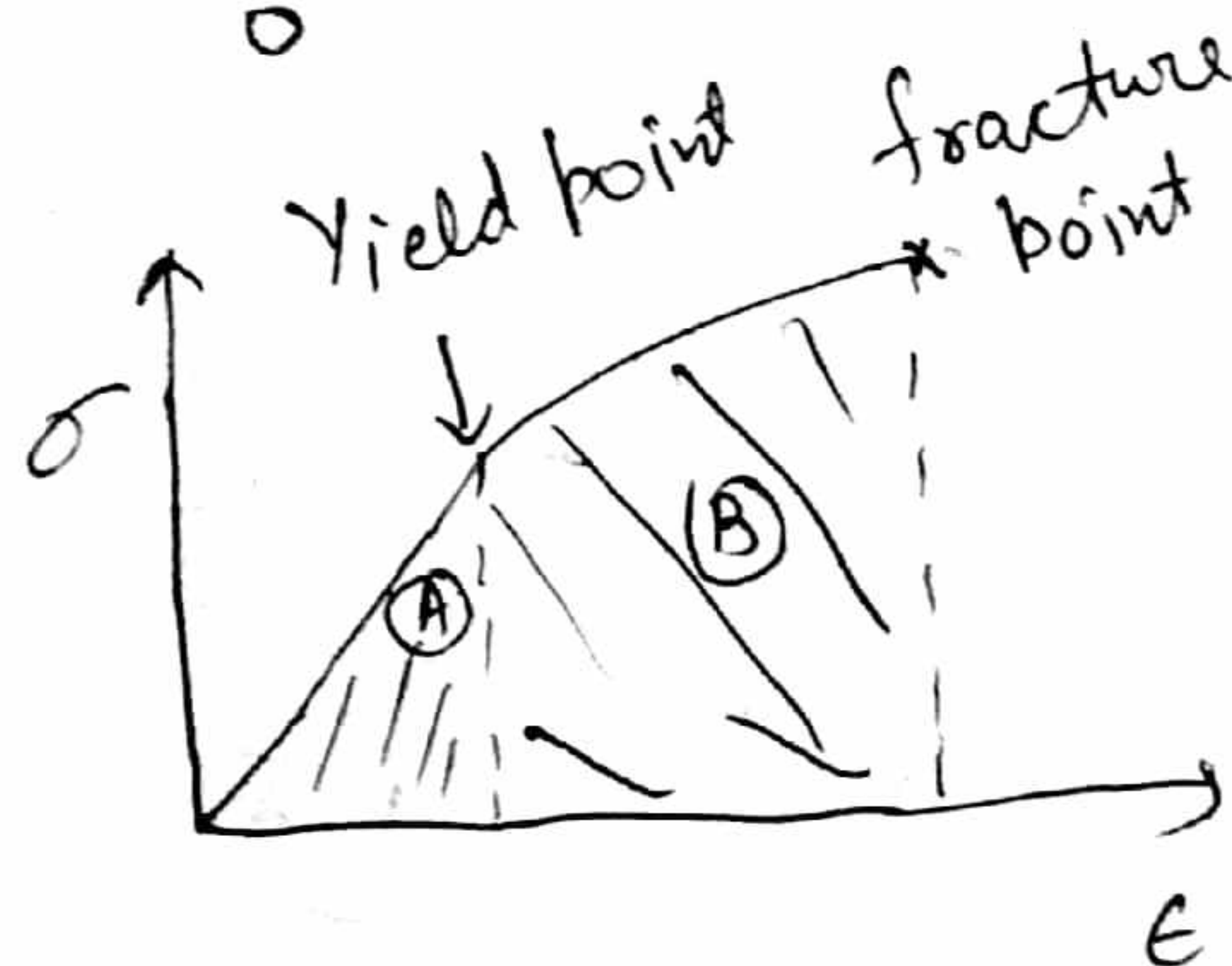
$$dU = \sigma_x (d\epsilon_x) \frac{dx dy dz}{\text{Volume}}$$

$$\text{Total strain energy} = \int_0^{\epsilon_x} \sigma_x d\epsilon_x dv + \int_0^{\epsilon_y} \sigma_y d\epsilon_y dv + \int_0^{\epsilon_z} \sigma_z d\epsilon_z dv + \int_0^{\gamma_{xy}} \tau_{xy} d\gamma_{xy} dv + \int_0^{\gamma_{xz}} \tau_{xz} d\gamma_{xz} dv + \int_0^{\gamma_{yz}} \tau_{yz} d\gamma_{yz} dv = U$$

$U_0 = \text{Strain energy / unit volume}$

(A) → Modulus of resilience

(A) + (B) → Modulus of toughness



$$dU_0 = \sigma_x d\epsilon_x + \sigma_y d\epsilon_y + \sigma_z d\epsilon_z + \tau_{xz} d\gamma_{xz} + \tau_{xy} d\gamma_{xy} + \tau_{yz} d\gamma_{yz}$$

$$dU_0 = \frac{\partial U_0}{\partial \epsilon_x} d\epsilon_x + \frac{\partial U_0}{\partial \epsilon_y} d\epsilon_y + \frac{\partial U_0}{\partial \epsilon_z} d\epsilon_z + \frac{\partial U_0}{\partial \gamma_{xz}} d\gamma_{xz}$$

$$+ \frac{\partial U_0}{\partial \gamma_{xy}} d\gamma_{xy} + \frac{\partial U_0}{\partial \gamma_{yz}} d\gamma_{yz}$$

On comparing above two equations strain energy density has the property that partial derivative of  $U_0$  w.r.t any strain stress component gives the corresponding

$$\sigma_x = \frac{\partial U_0}{\partial \epsilon_x}$$



$$\left\{ \frac{\partial U_0}{\partial \epsilon} \right\} = \{\sigma\} = [C] \{\epsilon\}$$

Integrating w.r.t  $d\epsilon$ ,  $U_0 = \frac{1}{2} \{\epsilon\}^T [C] \{\epsilon\}$

$$U = \frac{1}{2} \int_V \{\epsilon\}^T [C] \{\epsilon\} dV \rightarrow \text{strain energy}$$

External force effect

$$W = - \int_V \{u\}^T \{x\} dV - \int_S \{u\}^T \{p\} ds$$

$$\int_V (x_b u + y_b v + z_b w) dV \quad \int_S (x_s u + y_s v + z_s w) ds$$

⊙ Total strain energy ( $\Pi$ ) =  $\frac{1}{2} \int_V \{\epsilon\}^T [C] \{\epsilon\} dV - \int_V \{u\}^T \{x\} dV - \int_S \{u\}^T \{p\} ds$

$$\begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{bmatrix} \frac{\partial u}{\partial x} \\ \frac{\partial v}{\partial y} \\ \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \end{bmatrix}$$

$$\Rightarrow \begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} \end{bmatrix} \begin{Bmatrix} u \\ v \end{Bmatrix}$$

$$\begin{Bmatrix} \epsilon_x \\ \epsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{bmatrix} \frac{\partial \psi_1}{\partial x} & 0 & \frac{\partial \psi_2}{\partial x} & 0 & \frac{\partial \psi_3}{\partial x} & 0 \\ 0 & \frac{\partial \psi_1}{\partial y} & 0 & \frac{\partial \psi_2}{\partial y} & 0 & \frac{\partial \psi_3}{\partial y} \\ \frac{\partial \psi_1}{\partial x} & \frac{\partial \psi_1}{\partial y} & \frac{\partial \psi_2}{\partial x} & \frac{\partial \psi_2}{\partial y} & \frac{\partial \psi_3}{\partial x} & \frac{\partial \psi_3}{\partial y} \end{bmatrix} \begin{Bmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \end{Bmatrix}$$

$B \rightarrow$  Interpolation operator matrix  $\rightarrow$  Displacement matrix

$$\{\epsilon\} = [B] \{d\}$$

$\{d\} \rightarrow$  specifically nodal value whose values are constant

$$\{u\} = \{\psi\} \{d\}$$

$\{u\} \rightarrow$  generalised displacement matrix

$$\Pi = \frac{1}{2} \int_V \{d\}^T [B]^T [C] [B] \{d\} dV$$

$$- \int_V \{d\}^T \{\psi\}^T \{x\} dV - \int_V \{d\}^T \{\psi\}^T \{P\} ds$$

The first variation of  $\Pi$  must be zero for all equilibrium conditions.

→ Principle of stationary potential energy.

$$\delta \Pi = \{ \delta d \}^T \left[ \int_V [B]^T [C] [B] dV \{d\} - \int_V [ \psi ]^T \{x\} dV - \int_V [ \psi ]^T \{P\} ds \right] = 0$$

$$\left[ \int_V [B]^T [C] [B] dV \right] \{d\} = \left[ \int_V [ \psi ]^T \{x\} dV + \int_V [ \psi ]^T \{P\} ds \right] + \{Q\} + \{f\}$$

K matrix

$$[K] \{u\} = \{f\} + \{Q\}$$

This term will come, if at any node external loads are given

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Plane stress problem

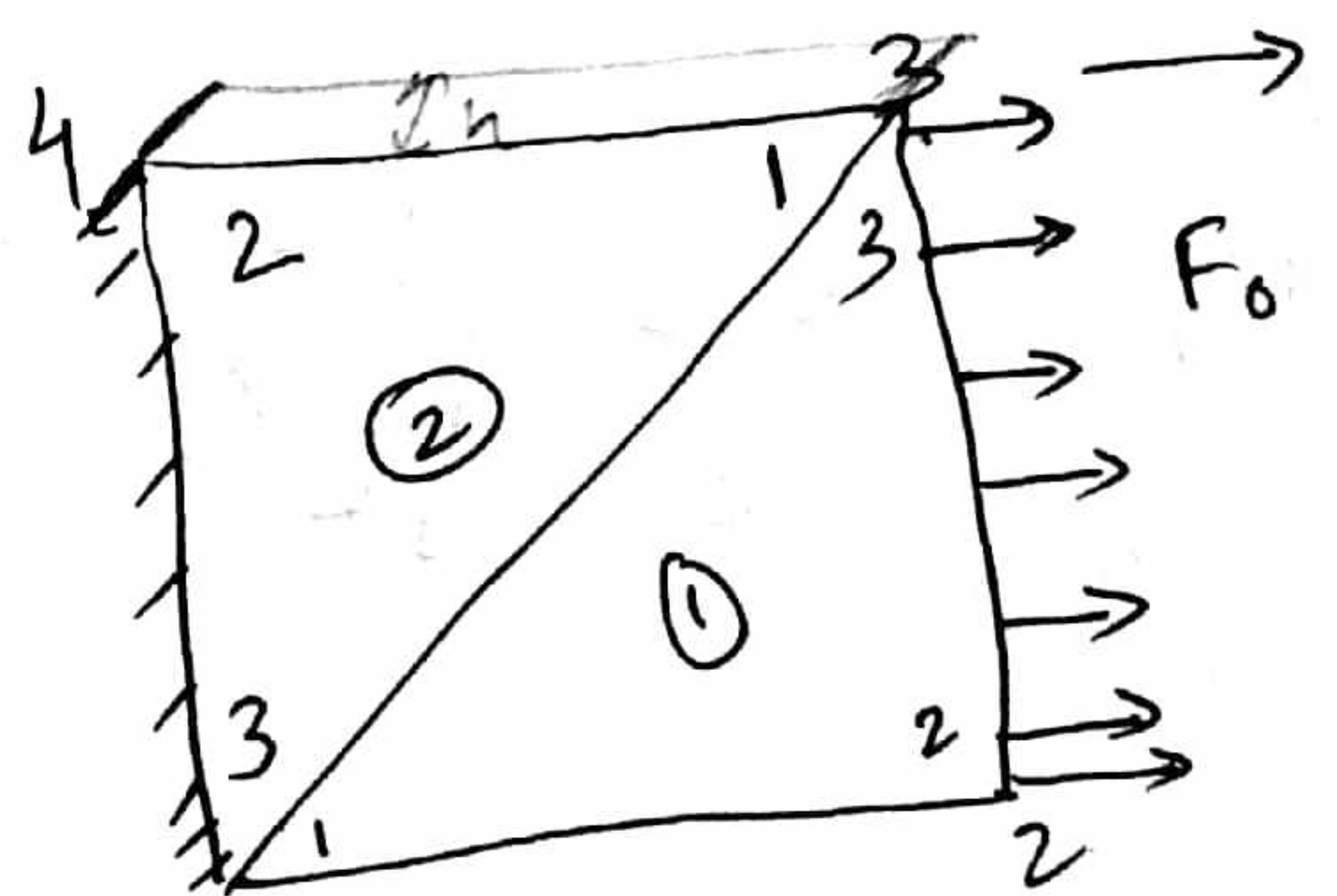
$$k_{ij} = h \int_{\Omega} B^T C B \frac{dx dy}{\text{Area}}$$

↑  
thickness

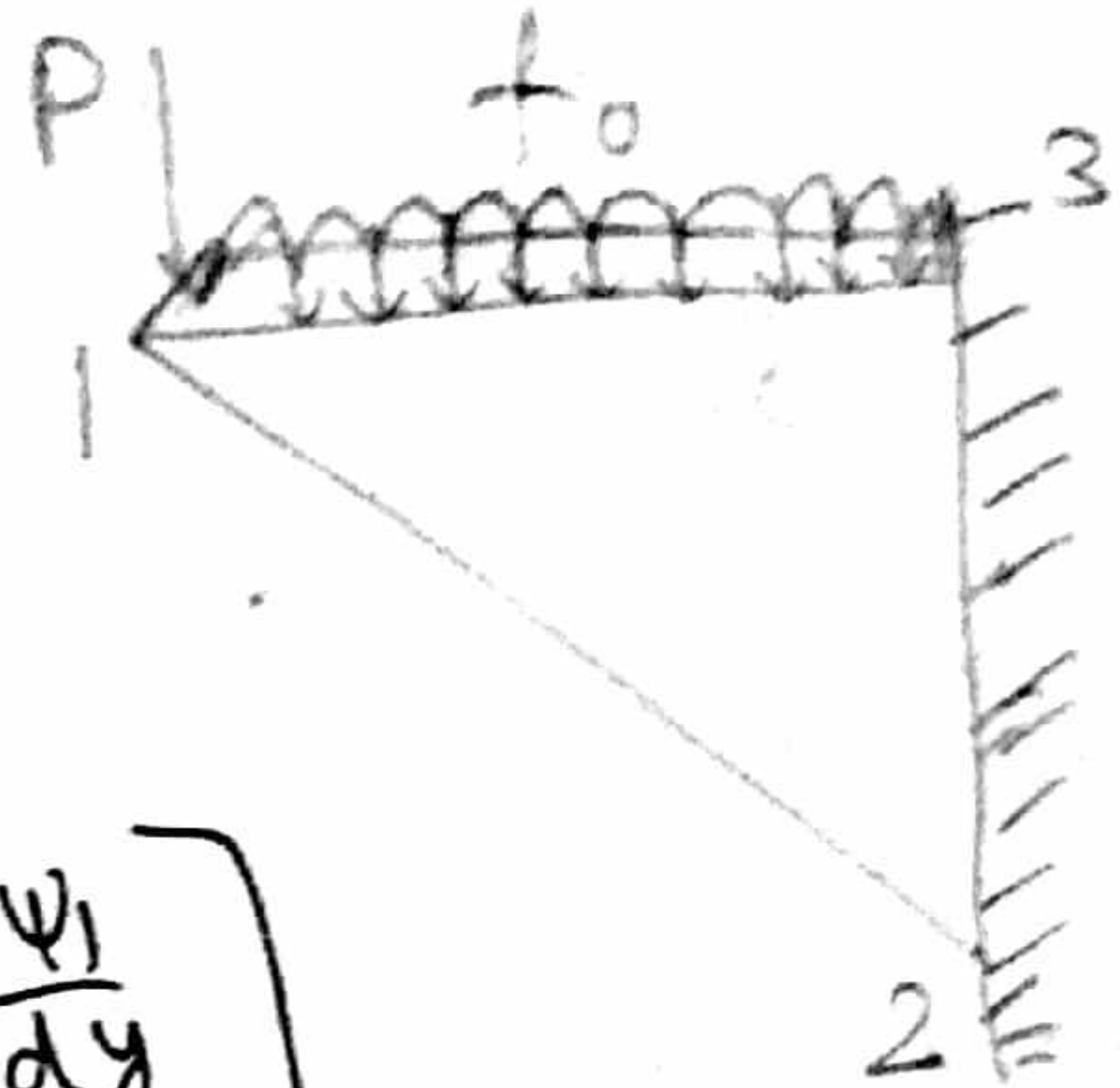
$$[C] = \begin{bmatrix} C_{11} & C_{12} & 0 \\ C_{21} & C_{22} & 0 \\ 0 & 0 & C_{66} \end{bmatrix}$$

$$[B] = \begin{bmatrix} \frac{d\phi_j}{dx} & 0 \\ 0 & \frac{d\phi_j}{dy} \\ \frac{d\phi_j}{dy} & \frac{d\phi_j}{dx} \end{bmatrix}$$

→ repeats with  $j=1,2,3$



$$\begin{bmatrix} \frac{d\psi_1}{dx} & 0 & \frac{d\psi_1}{dy} \\ 0 & \frac{d\psi_1}{dy} & \frac{d\psi_1}{dx} \end{bmatrix} \begin{bmatrix} c_{11} & c_{12} & 0 \\ c_{21} & c_{22} & 0 \\ 0 & 0 & c_{66} \end{bmatrix} \begin{bmatrix} \frac{d\psi_1}{dx} & 0 \\ 0 & \frac{d\psi_1}{dy} \\ \frac{d\psi_1}{dy} & \frac{d\psi_1}{dx} \end{bmatrix}$$



$$\begin{bmatrix} \frac{d\psi_1}{dx} & 0 & \frac{d\psi_1}{dy} \\ 0 & \frac{d\psi_1}{dy} & \frac{d\psi_1}{dx} \end{bmatrix} \begin{bmatrix} c_{11} \frac{d\psi_1}{dx} & c_{12} \frac{d\psi_1}{dy} \\ c_{21} \frac{d\psi_1}{dx} & c_{22} \frac{d\psi_1}{dy} \\ c_{66} \frac{d\psi_1}{dy} & c_{66} \frac{d\psi_1}{dx} \end{bmatrix}$$

$$c_{11} \left( \frac{d\psi_1}{dx} \right)^2 + c_{66} \left( \frac{d\psi_1}{dy} \right)^2 - c_{12} \frac{d\psi_1}{dy}$$

$$h \int_A \left[ c_{11} \frac{d\psi_1}{dx} \cdot \frac{d\psi_1}{dx} + c_{66} \frac{d\psi_1}{dy} \cdot \frac{d\psi_1}{dy} + c_{21} \frac{d\psi_1}{dx} \frac{d\psi_1}{dy} + c_{66} \frac{d\psi_1}{dy} \frac{d\psi_1}{dx} \right] dx dy$$

$$c_{12} \frac{d\psi_1}{dy} \frac{d\psi_1}{dx} + c_{66} \frac{d\psi_1}{dx} \frac{d\psi_1}{dy} + c_{22} \frac{d\psi_1}{dy} \frac{d\psi_1}{dy} + c_{66} \frac{d\psi_1}{dx} \frac{d\psi_1}{dx}$$

$dx dy = \frac{1}{2} h dx$

$$\begin{bmatrix} k_{11} & k_{12} \\ k_{21} & k_{22} \end{bmatrix} \begin{bmatrix} u_1 \\ v_1 \end{bmatrix} = \begin{Bmatrix} -\frac{f_0 h}{2} \\ -\frac{f_0 h}{2} \end{Bmatrix} + \begin{Bmatrix} 0 \\ -P \end{Bmatrix}$$

For triangular element

$$\psi_i = \frac{1}{2A} (\alpha_i + \beta_i x + \gamma_i y)$$

$$\frac{\partial \psi_i}{\partial x} = \frac{\beta_i}{2A}$$

$$\frac{\partial \psi_i}{\partial y} = \frac{\gamma_i}{2A}$$

$$h \int_A c_{11} \frac{d\psi_1}{dx} \frac{d\psi_1}{dx} dx dy = c_{11} \frac{\beta_i^2 h}{4A}$$

$$K = \begin{bmatrix} \frac{h C_{11} \beta_1^2}{4A} + \frac{C_{66} h \gamma_1^2}{4A} & \\ \frac{C_{21} h \beta_1 \gamma_1}{4A} + \frac{C_{66} \beta_1 \gamma_1 h}{4A} & \end{bmatrix}$$

$$\begin{bmatrix} \frac{C_{12} h \beta_1 \gamma_1}{4A} + \frac{C_{66} \beta_1 \gamma_1 h}{4A} & \\ \frac{h C_{22} \gamma_1^2}{4A} + \frac{C_{66} \beta_1^2 h}{4A} & \end{bmatrix}$$

$$K = \frac{h}{4A} \begin{bmatrix} (C_{11} \beta_1^2 + C_{66} \gamma_1^2) & \\ (C_{12} + C_{66}) \beta_1 \gamma_1 & \end{bmatrix}$$

$$\begin{bmatrix} (C_{12} + C_{66}) \beta_1 \gamma_1 & \\ (C_{22} \gamma_1^2 + C_{66} \beta_1^2) & \end{bmatrix}$$